



Radiative Transfer and Star Formation

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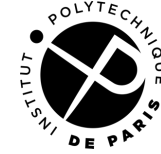
Under the guidance of Matthias González and Patrick Hennebelle

Baccalauréat S

Abu Dhabi, Emirats Arabes Unis



Programme PhD Track



2019

B.Sc.

2023

M1

2024

M2

2025

NEW YORK
UNIVERSITY



ABU DHABI

Bachelor of Science

- ❖ Spécialisation en Physique
- ❖ Mineure en mathématiques appliquées et en philosophie
- ❖ Stage en astrophysique théorique (1.5 an) a NYUAD – Andrea MACCIO



M1: Parcours Physique

- ❖ Cours du master High Energy Physics (HEP)
- ❖ Stage 4 mois AIM/LMPA – Patrick HENNEBELLE



M2: Physique des Plasmas et de la Fusion

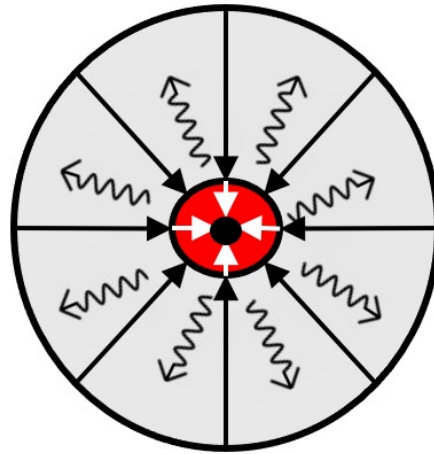
- ❖ Spécialisation Plasmas Astrophysiques
- ❖ Stage 5.5 mois AIM/LMPA – Matthias GONZÁLEZ

The collapse of dense molecular clouds

RT is important for:

- Realistic temperature gradients
- Dust enrichment
- Formation of planets

Second Collapse



Optically Thin Fluid

- H_2 dissociates
- Quasi-isothermal contraction

Modeling complicated by:

- Scale dynamics
- Variety of processes

Need approximations!

Pillars of Creation — JWST. Credit : NASA / ESA / CSA / STScI – 2022

RAMSES branches

	RAMSES ISM	RAMSES TINE
Photon Wavelengths:	...	...
Moment Model:	...	...
Solver Type:	...	...
Flux Limiting Scheme:	...	...

- Let's quickly walk through the differences in RT treatment:

The radiative transfer equation

$$\frac{1}{c} \frac{\partial I_\nu}{\partial t} + \mathbf{n} \cdot \nabla I_\nu = -\chi_\nu I_\nu + \eta_\nu,$$

where $I_\nu(\mathbf{x}, \mathbf{n}, t)$ is the radiation specific intensity, $\chi_\nu(\mathbf{x}, \mathbf{n}, t)$ is the absorption coefficient and $\eta_\nu(\mathbf{x}, \mathbf{n}, t)$ is the emission coefficient.

ISM – Gray photons



Difference 1: Photon wavelength

TINE – Multigroup photons



Moment models

$$E_r(\mathbf{x}, t) = \frac{1}{c} \int_0^\infty \int_{S^2} I_\nu(\mathbf{x}, \mathbf{n}, t) d\Omega d\nu,$$

$$\mathbf{F}_r(\mathbf{x}, t) = \int_0^\infty \int_{S^2} I_\nu(\mathbf{x}, \mathbf{n}, t) \hat{\mathbf{n}} d\Omega d\nu,$$

$$\mathbb{P}_r(\mathbf{x}, t) = \frac{1}{c} \int_0^\infty \int_{S^2} I_\nu(\mathbf{x}, \mathbf{n}, t) \hat{\mathbf{n}} \otimes \hat{\mathbf{n}} d\Omega d\nu.$$

$$\frac{\partial E_\nu}{\partial t} + \nabla \cdot \mathbf{F}_\nu = -\rho \kappa_\nu c E_\nu + S_\nu,$$

$$\frac{\partial \mathbf{F}_\nu}{\partial t} + c^2 \nabla \cdot \mathbb{P}_\nu = -\rho \kappa_\nu c \mathbf{F}_\nu.$$

Difference 2: Moment model

Algorithm

Timestepping with the Courant–Friedrichs–Lewy condition:

$$C = \frac{u\Delta t}{\Delta x} < C_{max},$$


Implicit solver + large number of iterations.

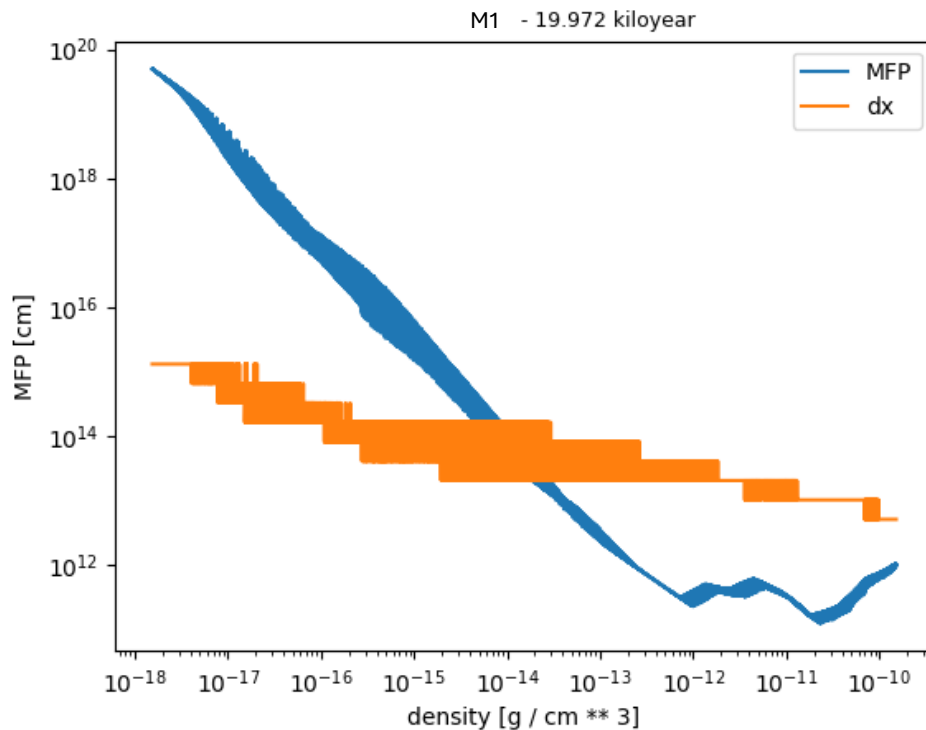
Explicit solver + reduced speed of light.

Difference 3: Algorithm

$$\tilde{c} = 10^{-4}c.$$

Flux limiting scheme

Difference 4: Flux Limiting Scheme



Problem - Overdiffusivity
Bouchut et al. (2004)



Trapped photons
Rosdhal et al. (2015)
RAMSES Tine

$$E_s = \frac{2}{2 + 3\tau_c} E_r, \quad \mathbf{F}_s = \mathbf{F}_r,$$

$$E_t = \frac{3\tau_c}{2 + 3\tau_c} E_r, \quad \mathbf{F}_t = \mathbf{0},$$



Flux limiter
Minerbo (1978)
RAMSES ISM

$$\mathbf{F}_r = -\frac{c\lambda}{\sigma_R} \nabla E_r, \quad R = |\nabla E_r| / (\sigma_R E_r)$$

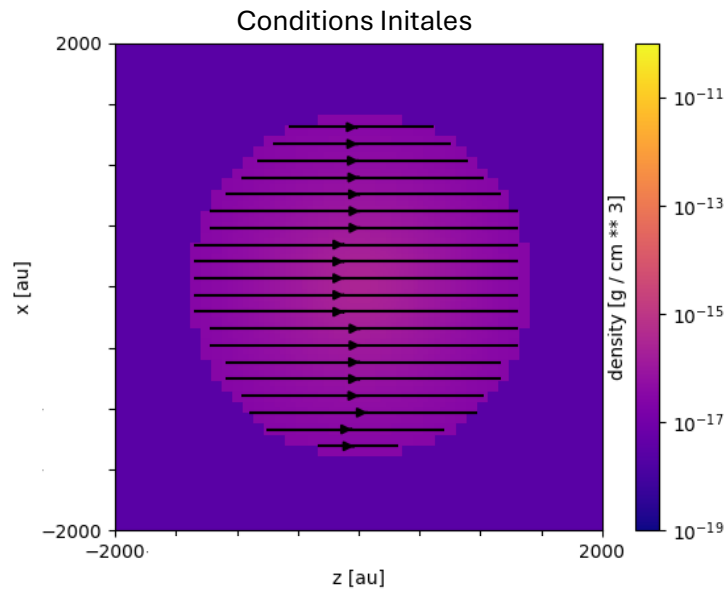
$$\lambda = \begin{cases} 2/(3 + \sqrt{9 + 12R^2}) & \text{if } 0 \leq R \leq 3/2, \\ (1 + R + \sqrt{1 + 2R})^{-1} & \text{if } 3/2 < R \leq \infty. \end{cases}$$

RAMSES branches: Summary

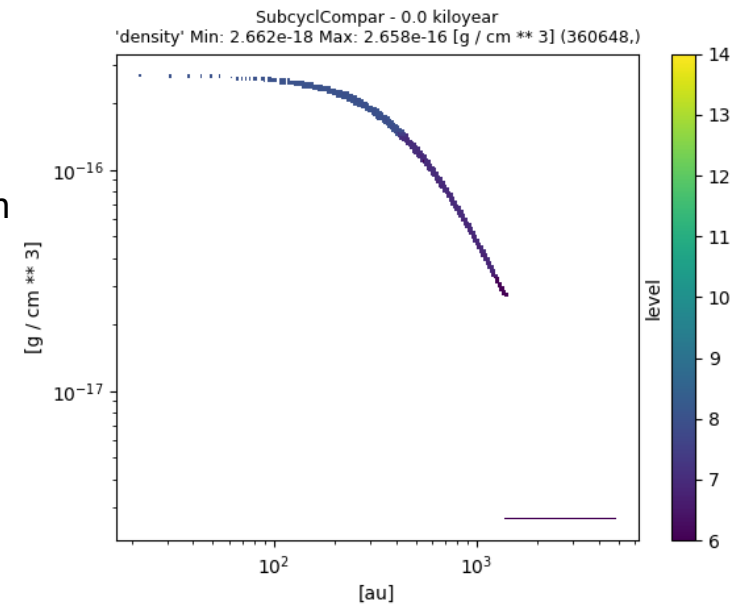
	RAMSES ISM (FLD)	RAMSES TINE (M1)
Photon Wavelengths:	Gray	Multigroup
Moment Model:	Flux Limited Diffusion	M1
Solver Type:	Implicit	Explicit
Flux Limiting Scheme:	Minerbo Limiter	Trapped Photons

- Protostellar collapses have historically been studied in RAMSES ISM
- RAMSES TINE has an advantageous treatment of RT, but is not made with smaller scales in mind
- **My main contribution so far:** Adapting RAMSES Tine to the scales of the ism, and implementing matching opacity computations from Vaytet et al. (2013)

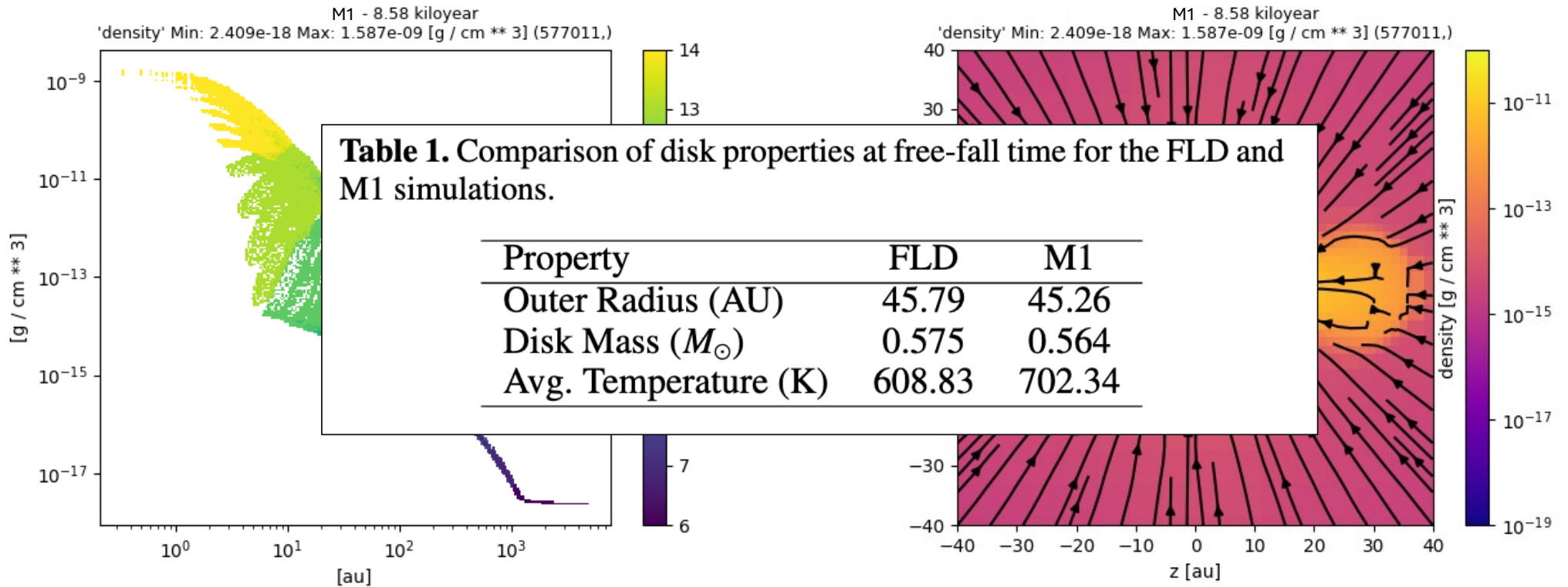
Initial conditions



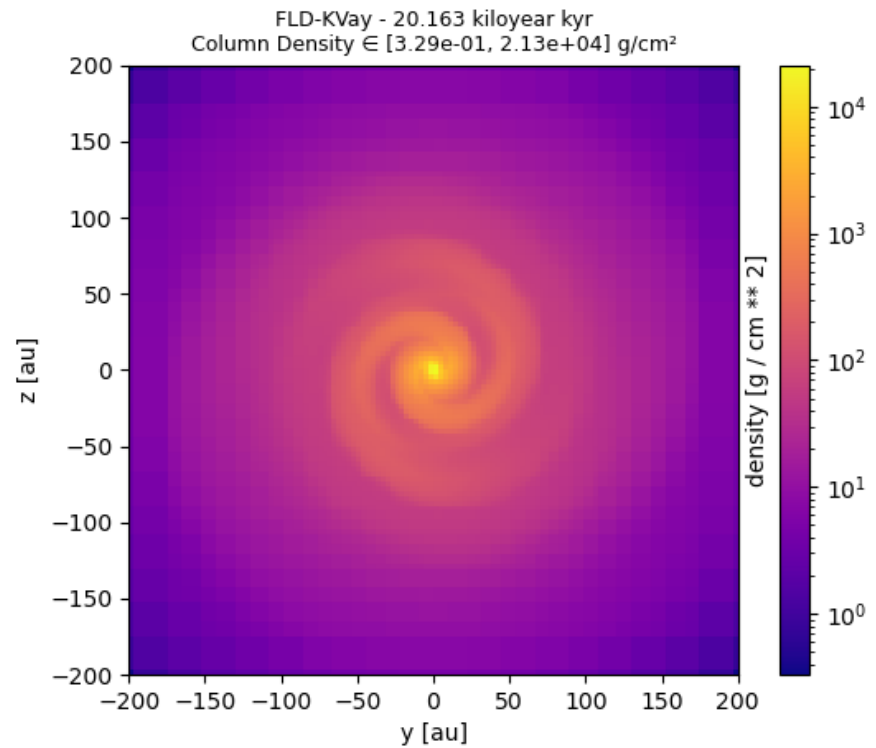
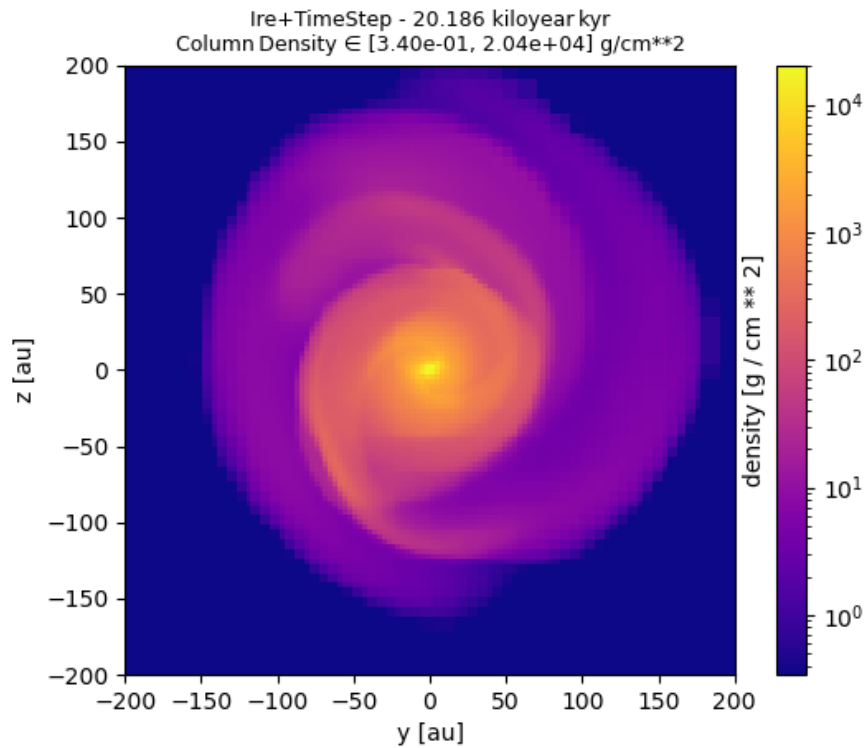
$M_{\text{core}} = 1 M_{\odot}$
 $R_{\text{core}} \approx 1400 \text{ AU}$
 H_2 and He in solar proportion
 $Z = Z_{\odot}$
 $\tau_{\text{ff}} \approx 8,5 \text{ Kiloyears}$
 $T = 10 \text{ K}$
 $L_{\text{box}} \approx 5600 \text{ AU}$
Dust to gas ratio = 1/100



Simulation properties at free fall time

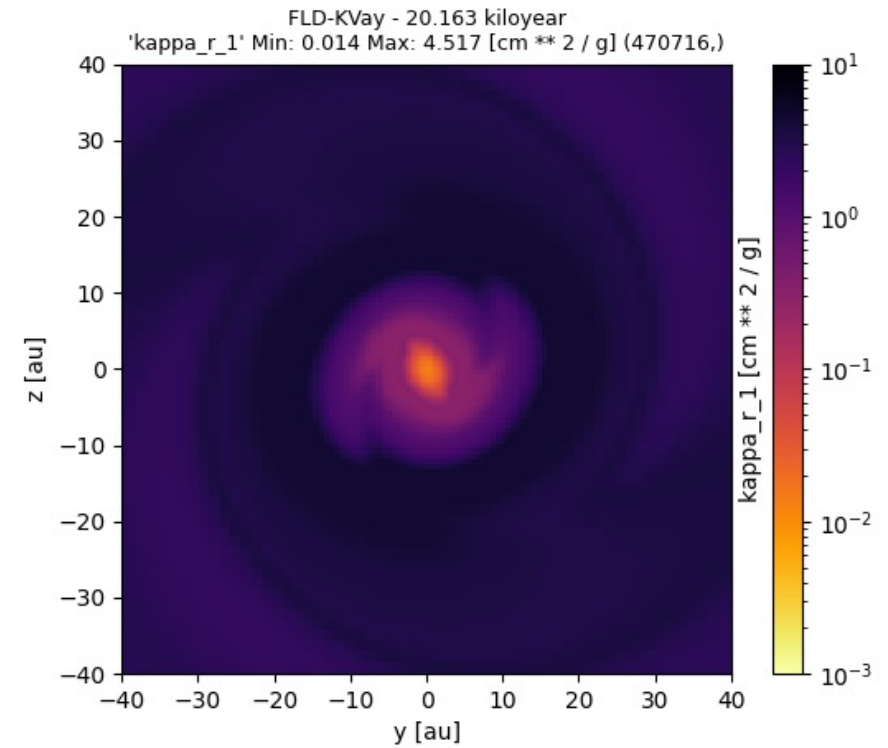
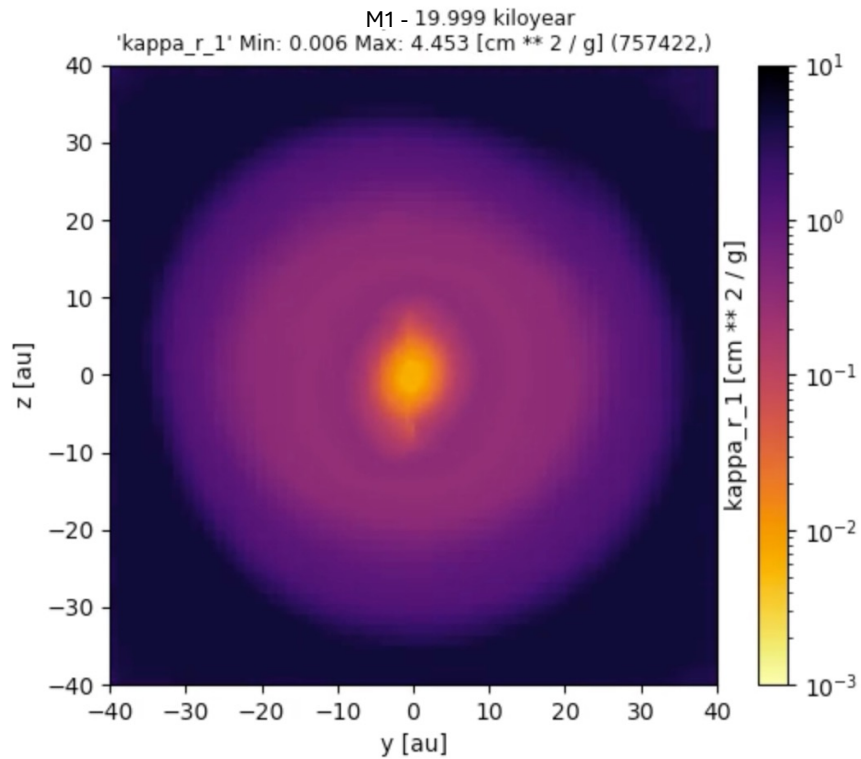


Later times – disk evolution

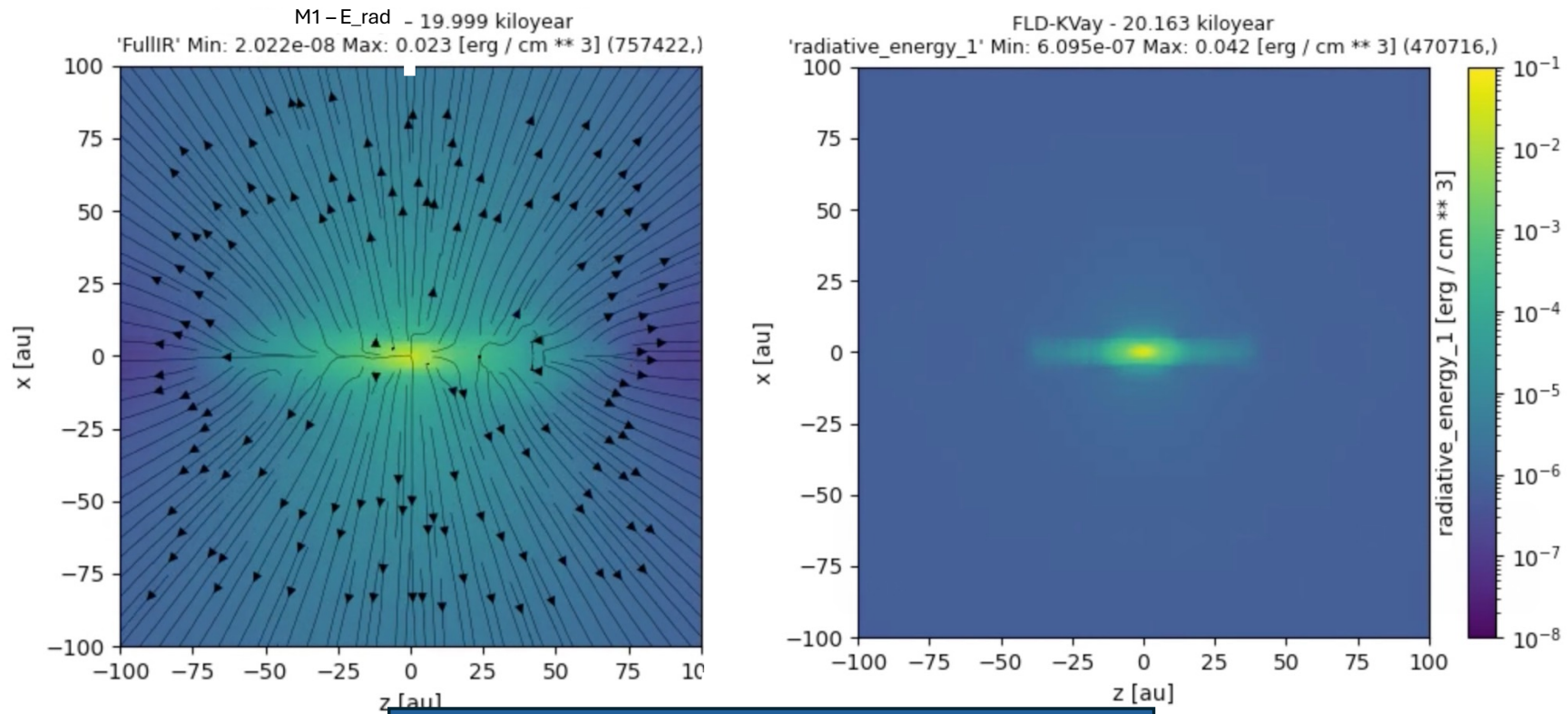


Disks look different: investigate optical properties

Opacity Maps

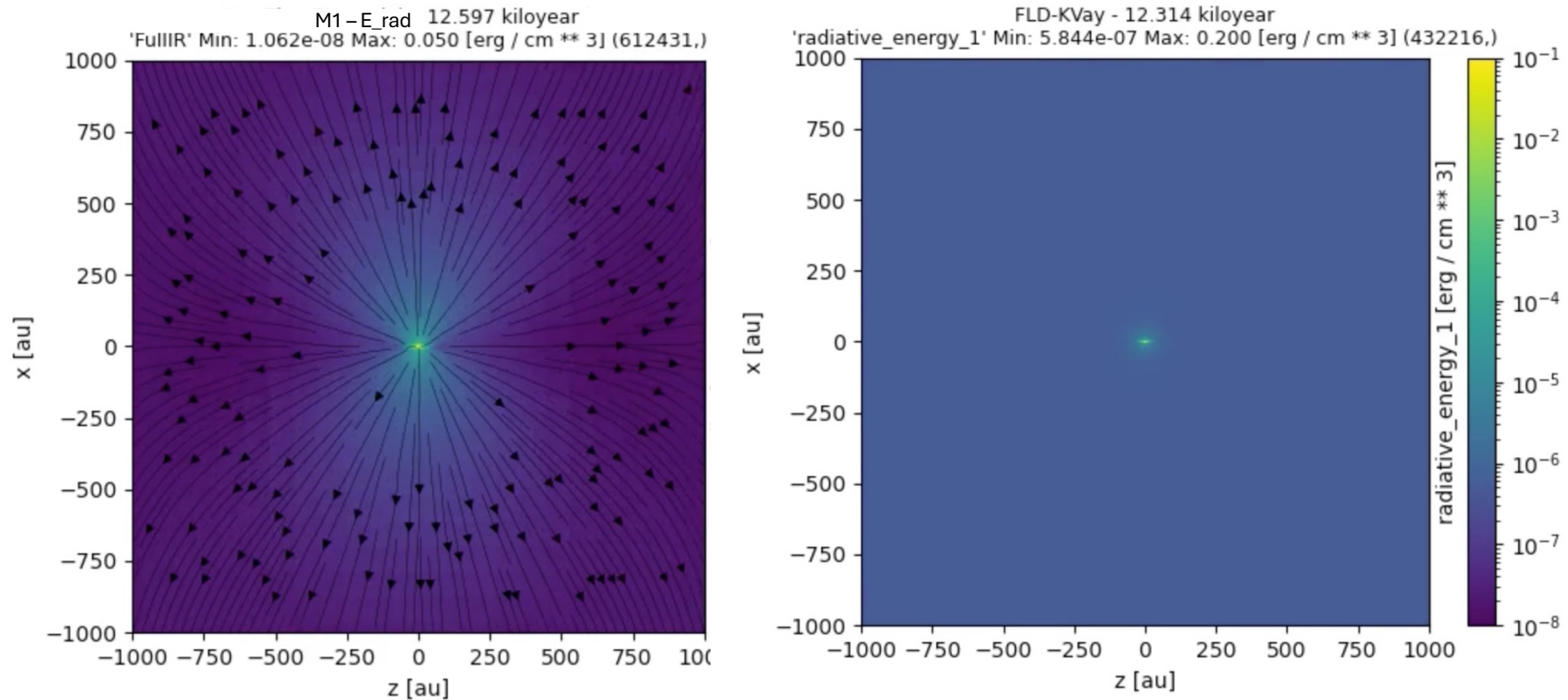


Radiative Energy Density



Main point: presence of shadows

Radiative Energy Density – Full box



Conclusion

- **At free fall time** FLD and M1 give similar disk radii (~ 45 AU) and masses ($\sim 0.57 M_{\odot}$); M1 disks are **hotter by $\sim 15\%$** .
- **After a few more kyrs** – disks present **differences** in size, temperature, susceptibility to fragmentation.

Next step – figure out if these results are physical:

- 1) Test the opacities: semi-analytic opacity-dependent polytrope (based on Matthew and Federath 2020)
- 2) Experiment with implementing a flux limiter in M1

Future project: realistic dust/gas temperature coupling $\rightarrow$ retroaction of dust on opacity $\rightarrow$ populations of disks $\rightarrow$ solve planet formation ;)