

Abe et al. (2025)

A Novel Mechanism to Drive Turbulence in Massive and Magnetized Star-Forming Filaments

Daisei ABE¹,

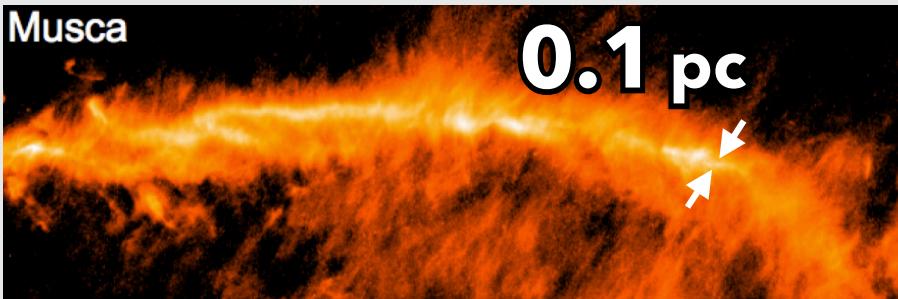
Tsuyoshi INOUE^{2,3}, Shu-ichiro INUTSUKA³, & Doris ARZOUMANIAN⁴

¹Tohoku University, ²Konan University, ³Nagoya University, ⁴Kyushu University

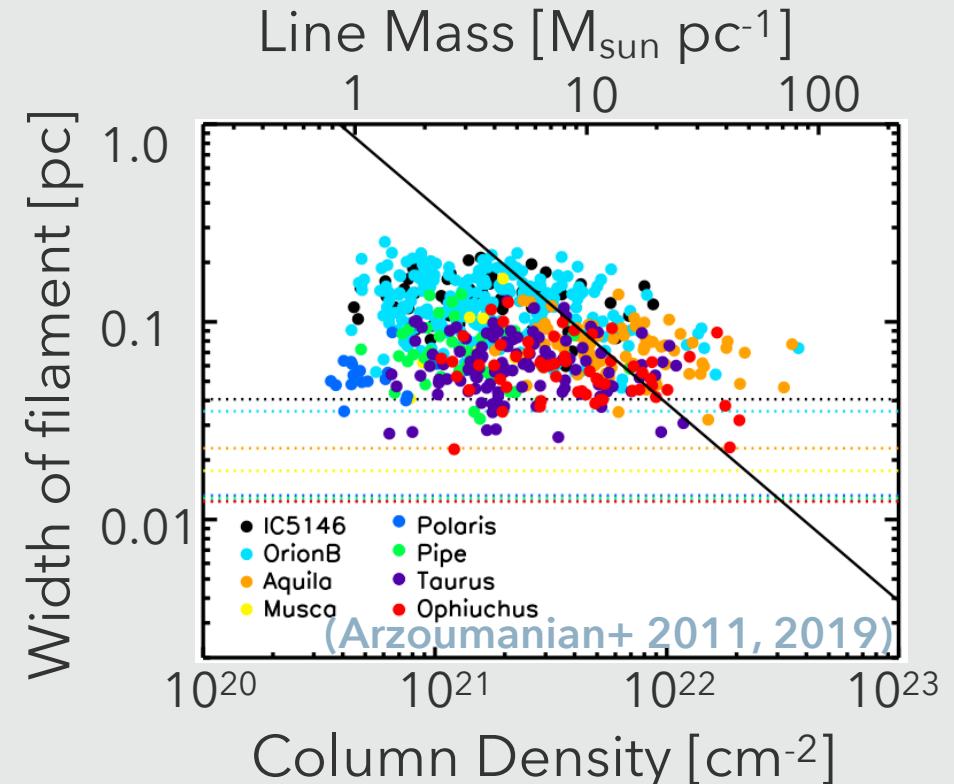
JSPS Research Fellow-PD [Tohoku University (Japan) / CEA Paris-Saclay

Questions: What controls the initial condition of star formation?

Star-forming Filaments



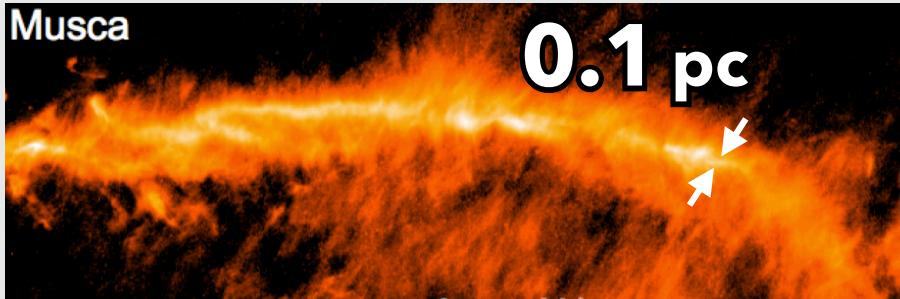
Filaments play an important role in setting the initial conditions.



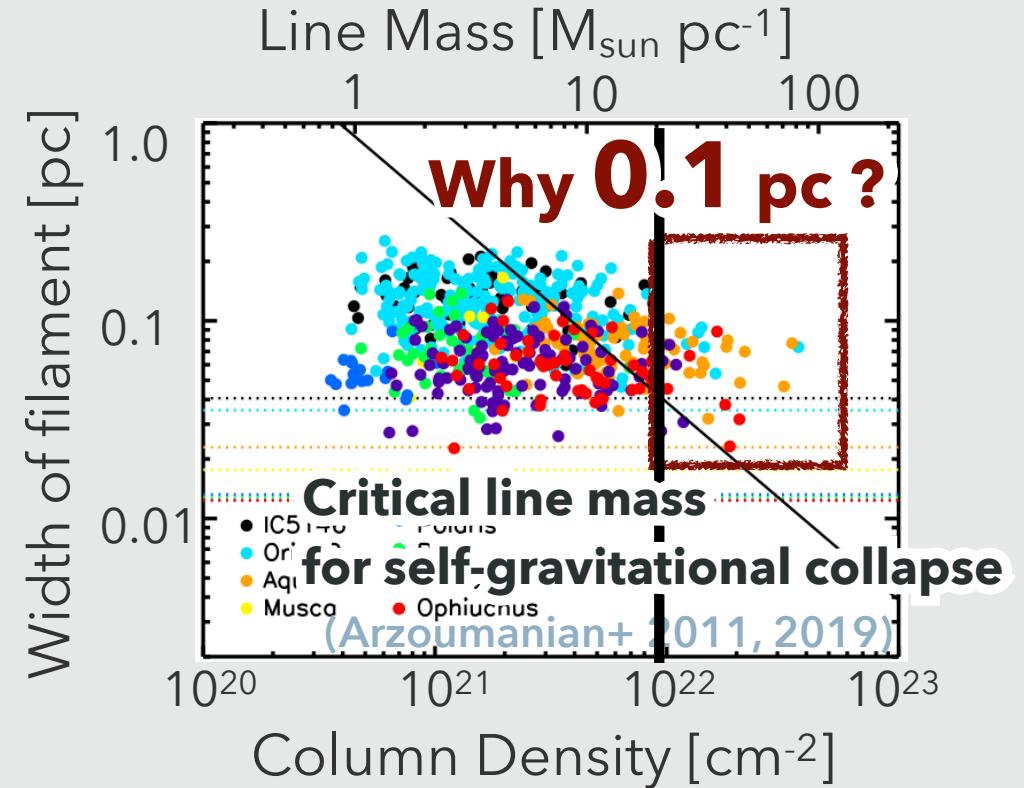
Observed filaments typically have a width of $\sim 0.1 \text{ pc}$

Questions: What controls the initial condition of star formation?

Star-forming Filaments



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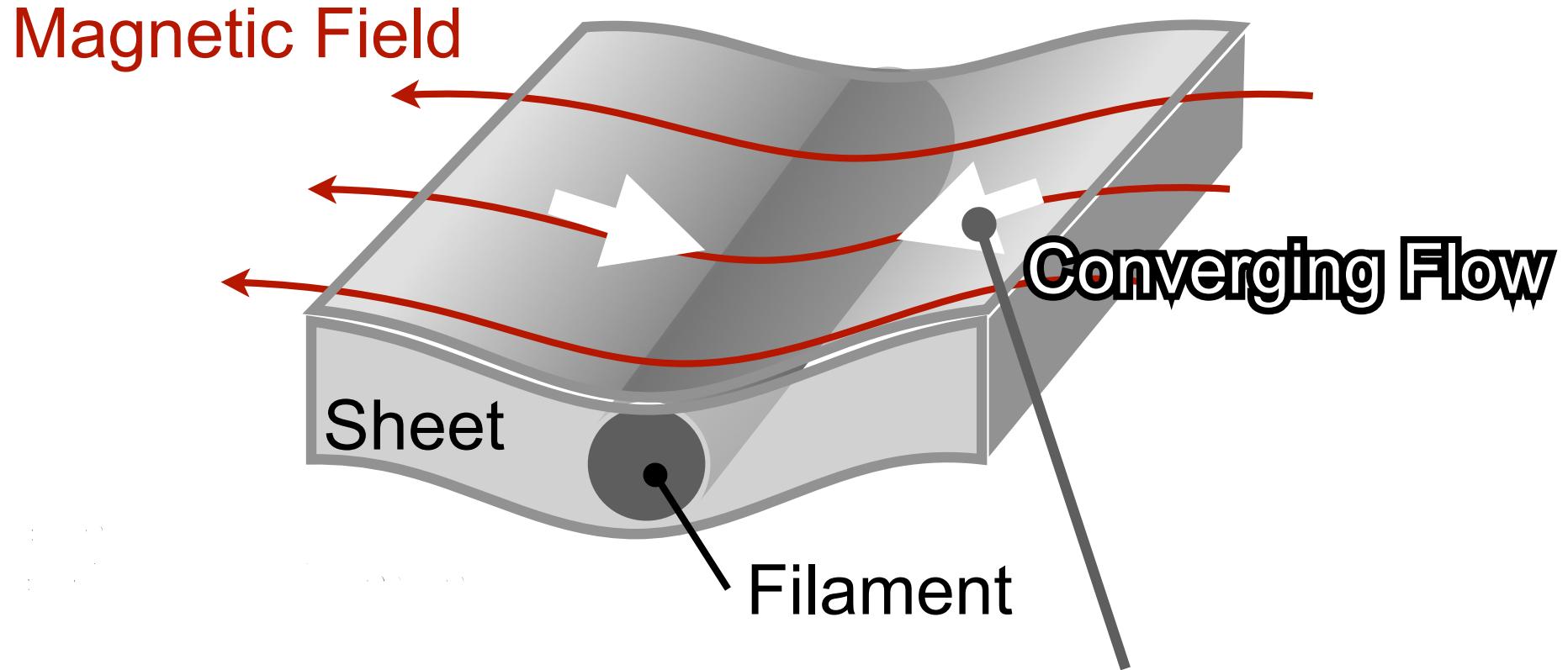
Observed filaments typically have a width of $\sim 0.1 \text{ pc}$

However, most simulations show much narrower width due to strong gravity.

→ **We are missing something.**

Theoretical Picture of Filament Evolution

e.g., Tomisaka & Ikeuchi (1983); Kitsionas & Whitworth (2007); Balfour et al. (2015); Inoue & Fukui (2013); Vaidya et al. (2013); Inoue et al. (2018); Chen & Ostriker (2014); Padoan & Nordlund (1999), Abe+ (2021), Pineda et al. (PPVII 2023)



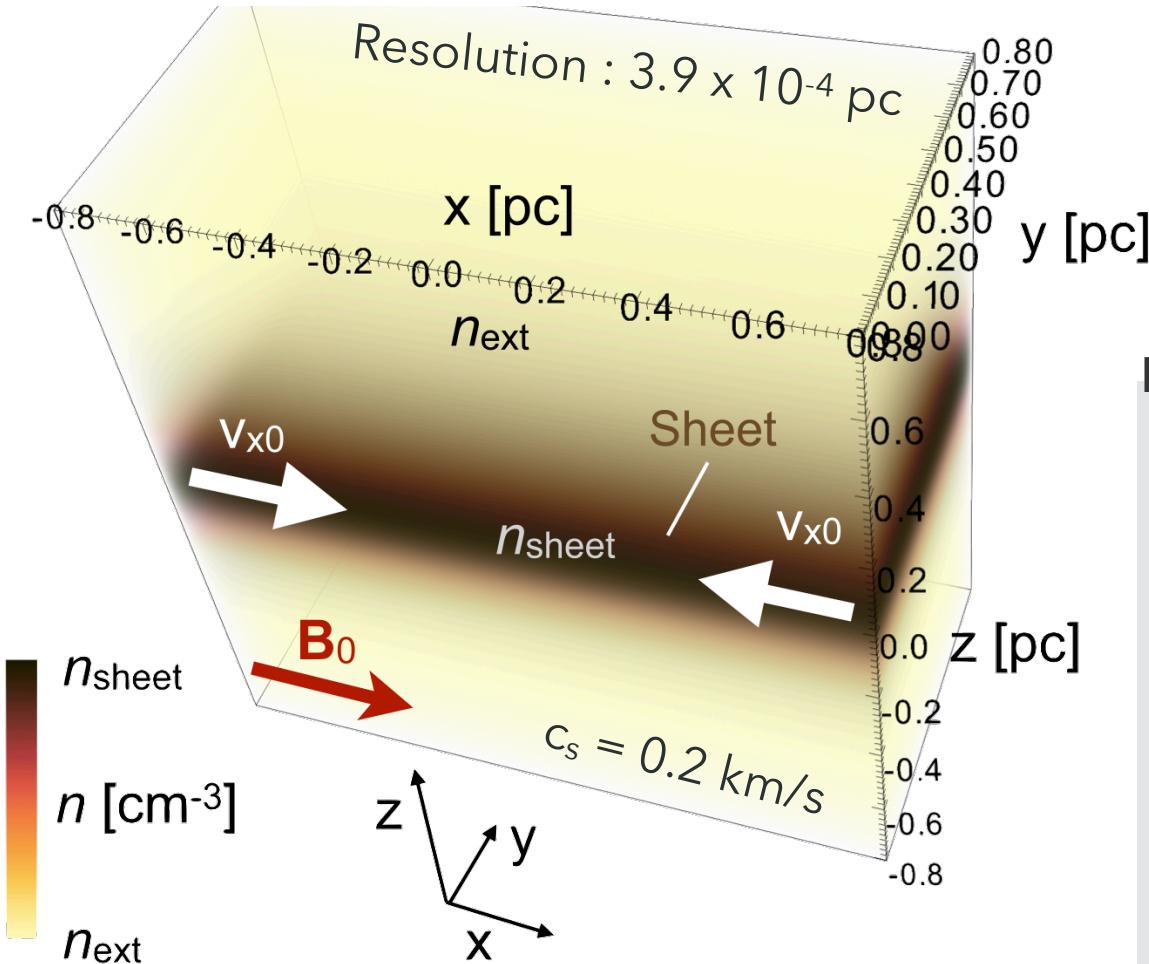
Filaments evolve via gas inflow along magnetic field lines.

We performed a simulation of this situation (Abe+ 2025).

Initial Condition for the Simulation

Simulations using Athena++ code (Stone+ 2020, Tomida & Stone 2023)

Initial Condition: Gas inflows along the B field \rightarrow filament formation



Boundary Condition

- x \rightarrow gas continue to flow in
- y,z \rightarrow Periodic

MHD including ambipolar diffusion

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + \mathbf{P} + \mathbf{B}^2/2) = -\rho \nabla \Phi$$

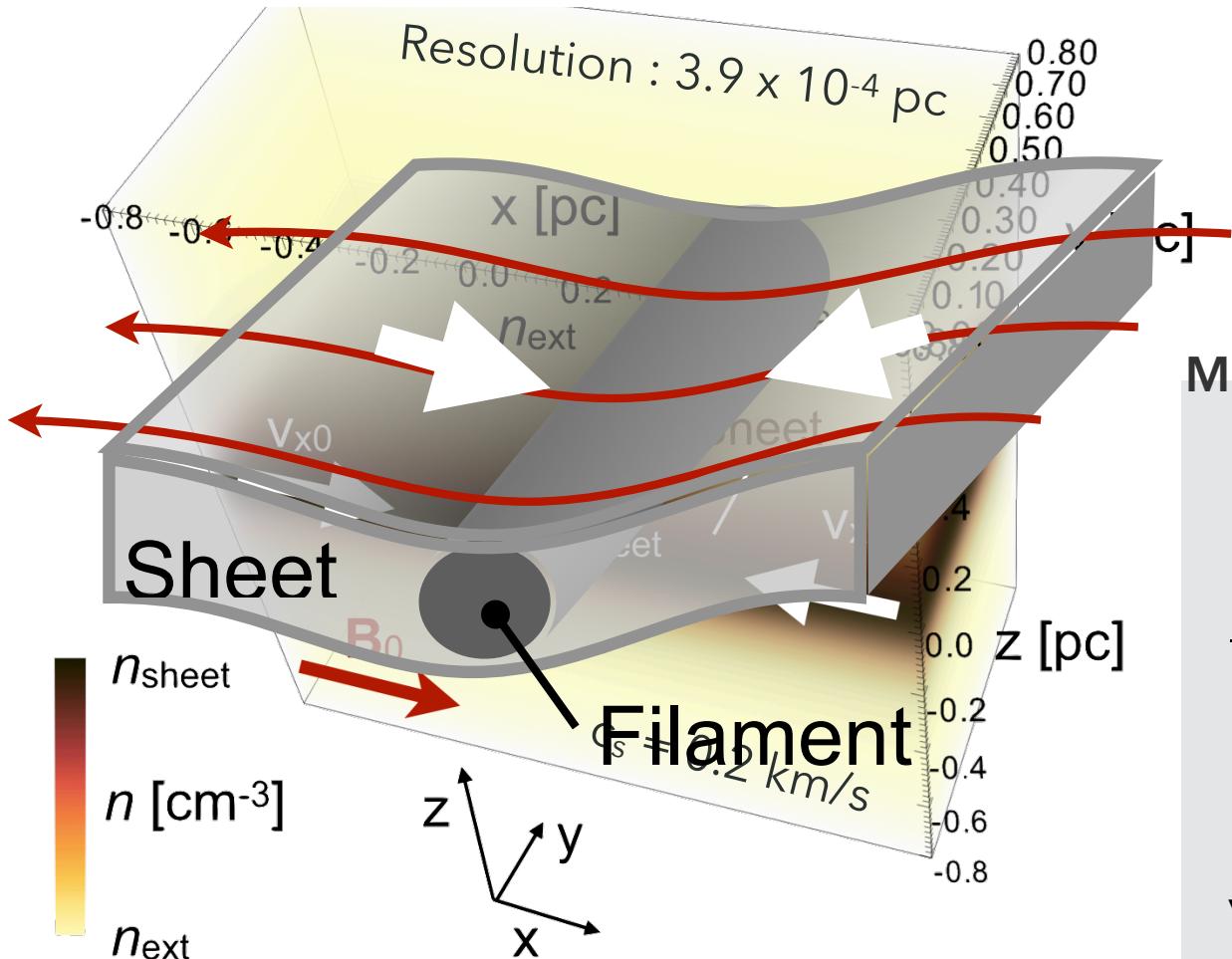
$$\frac{\partial \mathbf{B}}{\partial t} - \nabla \times \left[(\mathbf{v} \times \mathbf{B}) - \frac{\eta_{\text{AD}}}{|\mathbf{B}|^2} \mathbf{B} \times ((\nabla \times \mathbf{B}) \times \mathbf{B}) \right] = 0$$

$$\nabla^2 \Phi = 4\pi G \rho \quad \mathbf{P} = \rho c_s^2 \quad \eta_{\text{AD}} = \frac{B^2}{4\pi \gamma_{\text{in}} \rho_n \rho_i}$$

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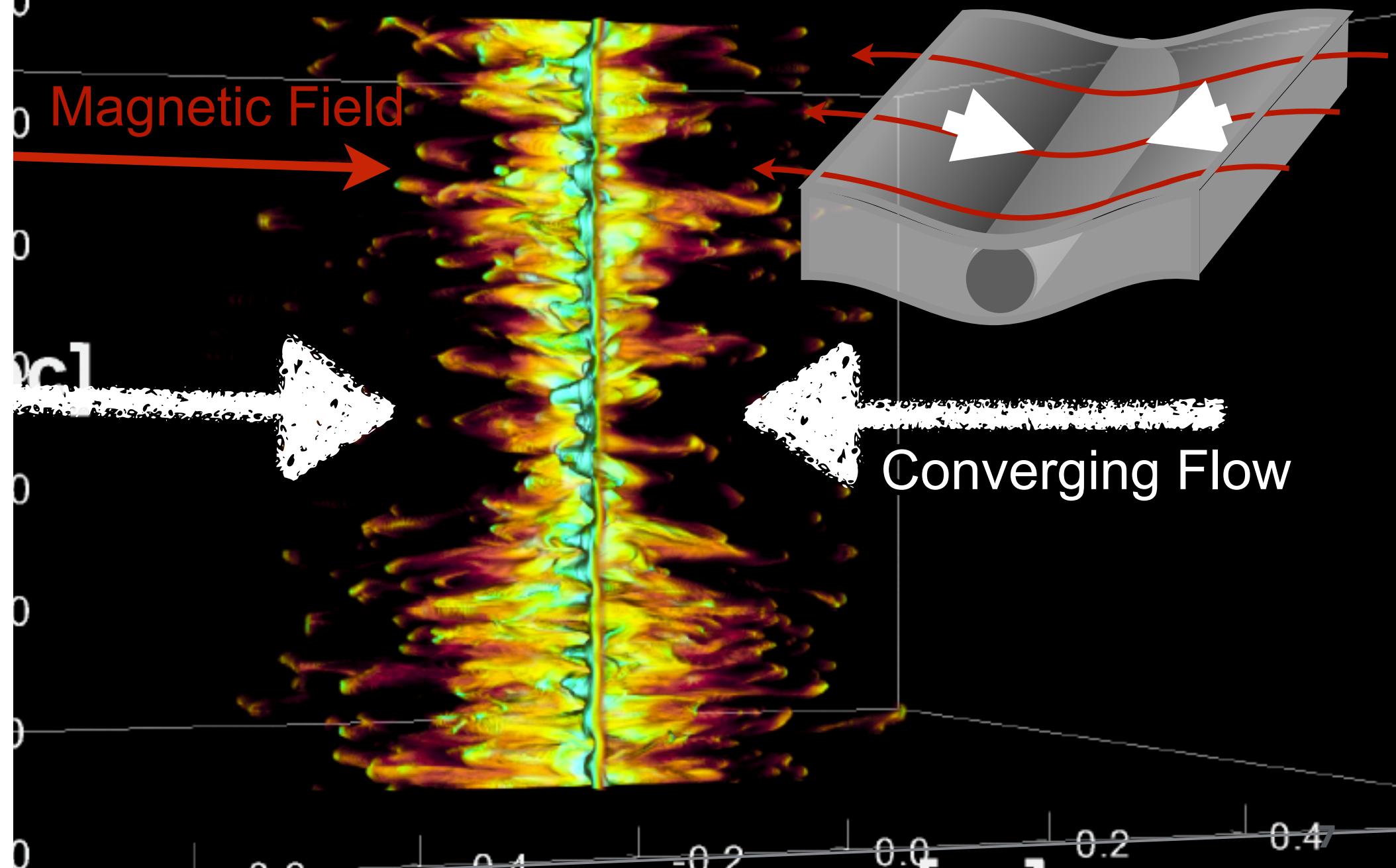
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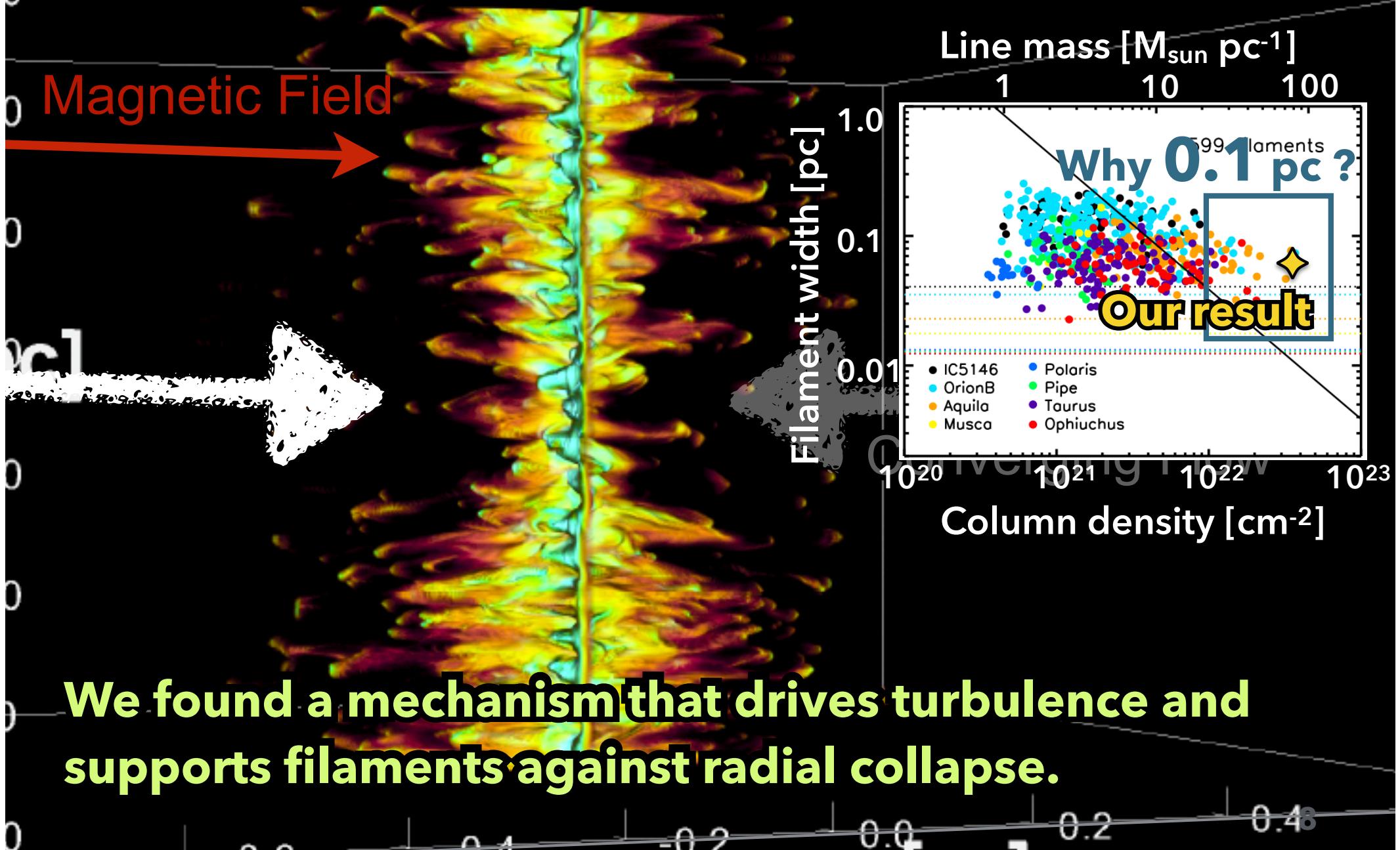
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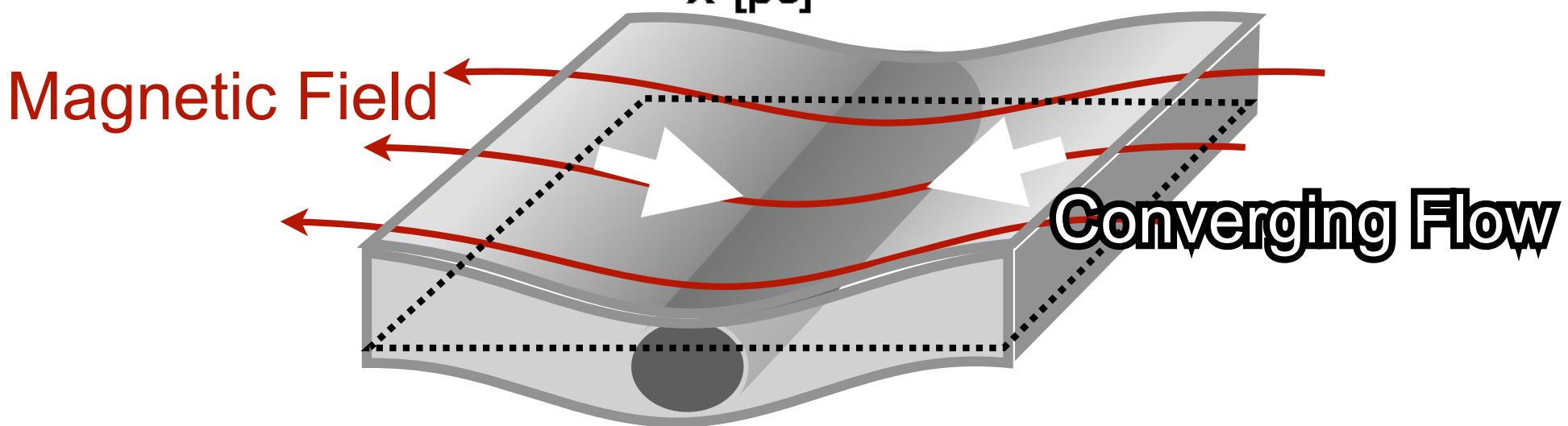
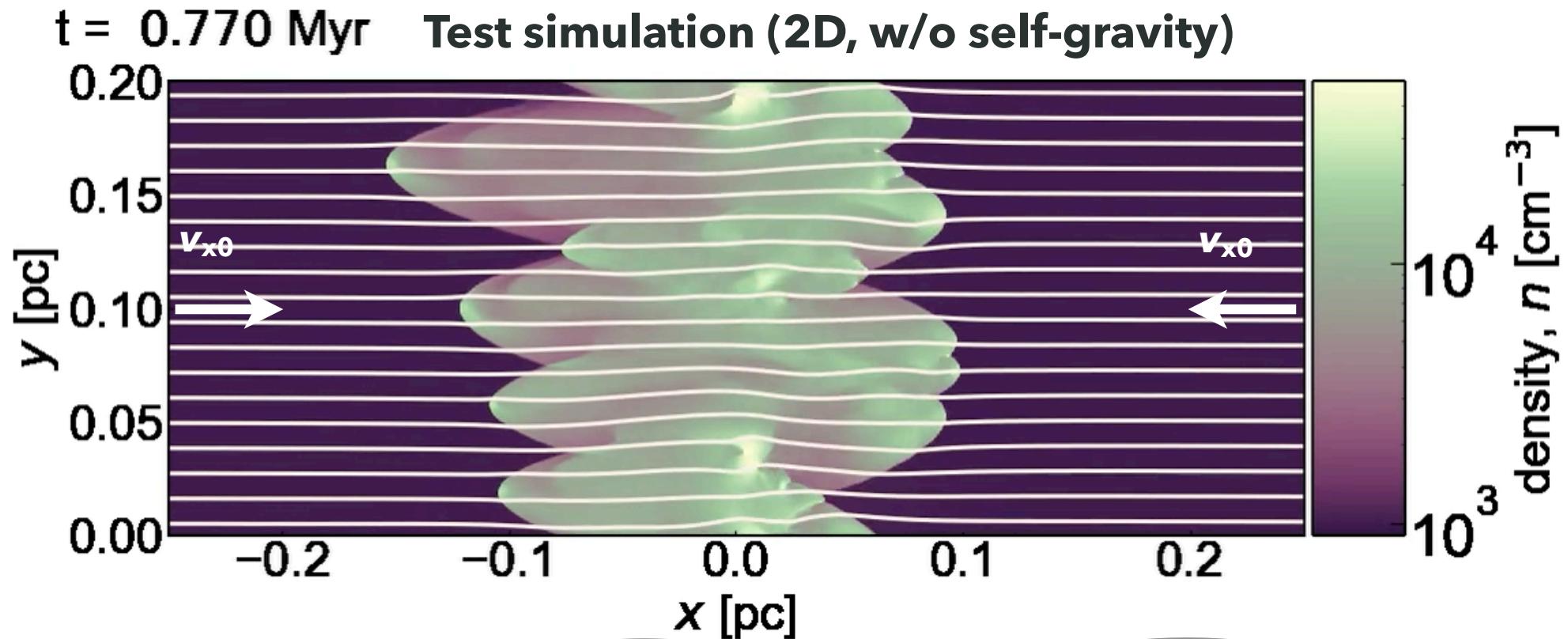
3D Simulation of Filament Evolution (Abe et al. 2025)



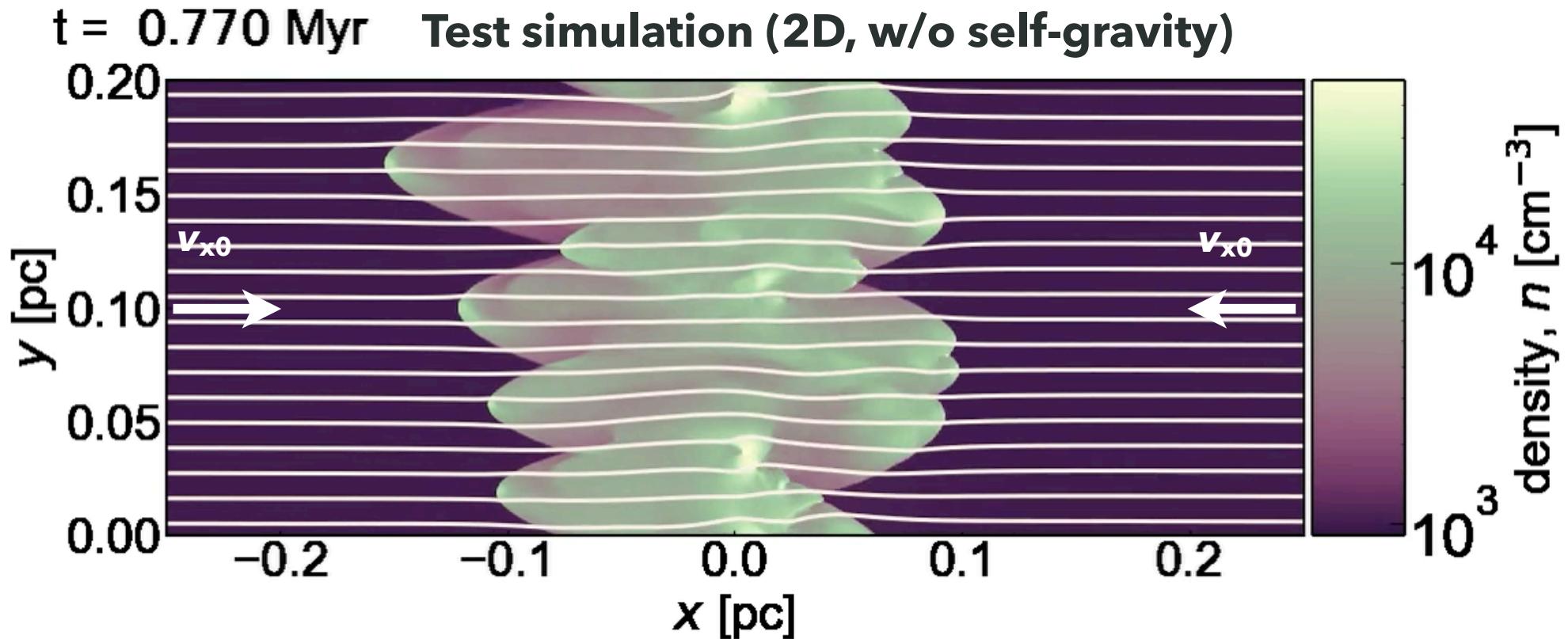
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What is the Mechanism?: Finding of “STORM”



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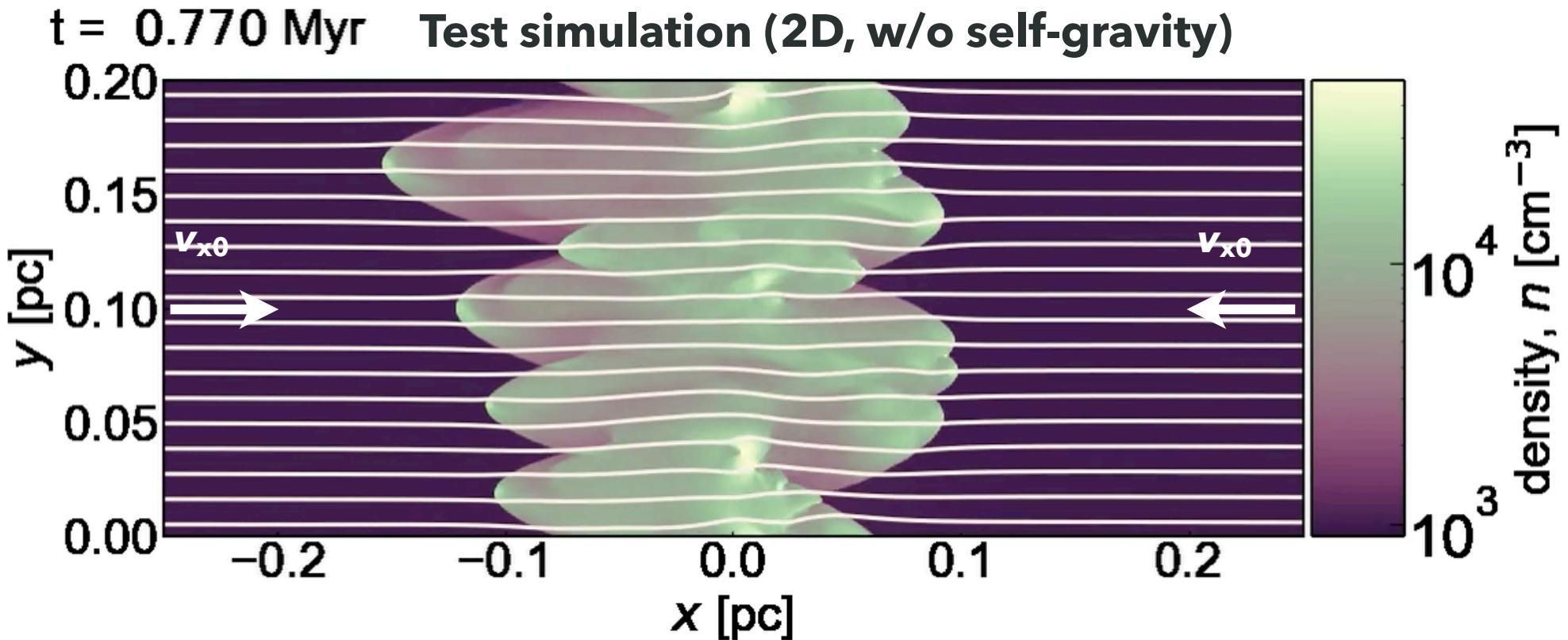


✓ Slow-shock Instability

Magnetic fields play an important role in this instability.

Perturbation of the shock front grows.

What is the Mechanism?: Finding of “STORM”



✓ Slow-shock Instability

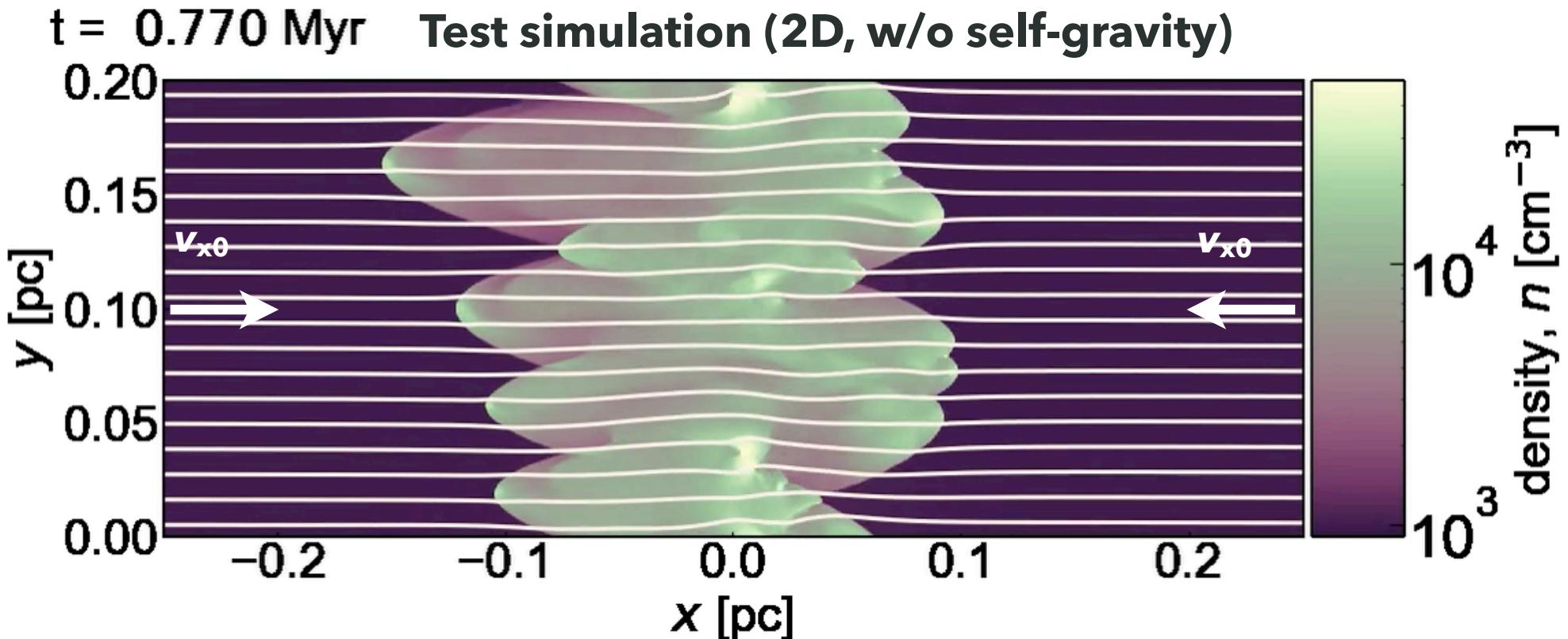
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✓ Partially ionized plasma (Ambipolar Diffusion)

Neutral gas can move freely without being frozen into the magnetic field.

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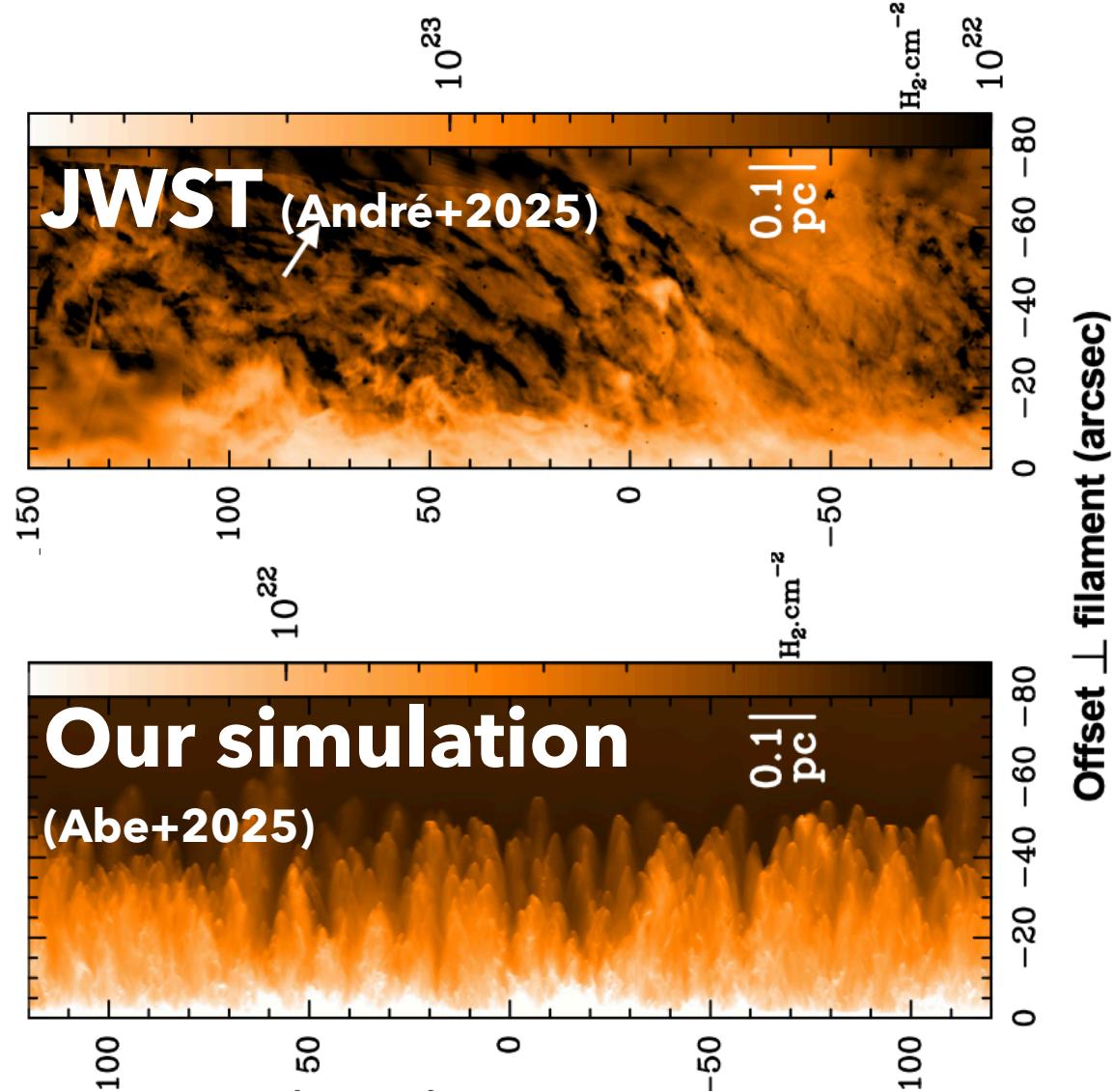
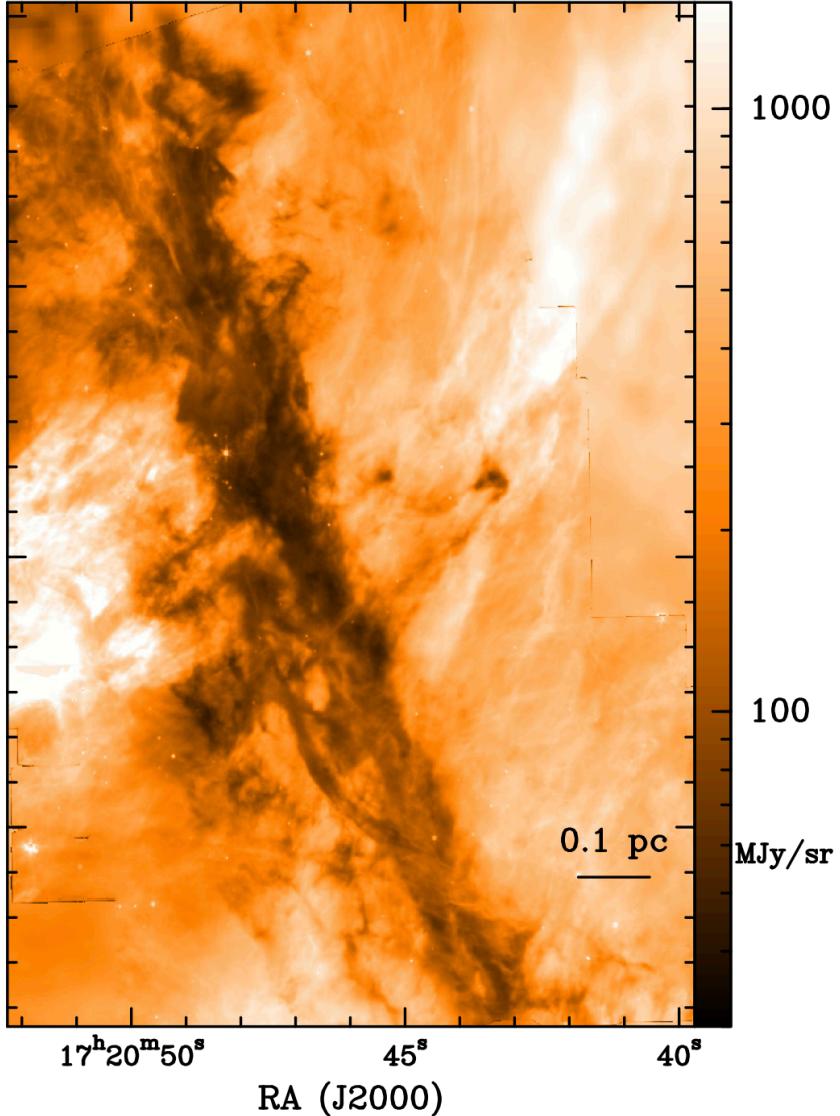
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“**STORM** (Slow-shock-mediated Turbulent flow Reinforced by Magnetic diffusion)”

Comparison with Observation



The spacing of the “fingers” & the column density power spectrum are consistent with the JWST observation.

Summary

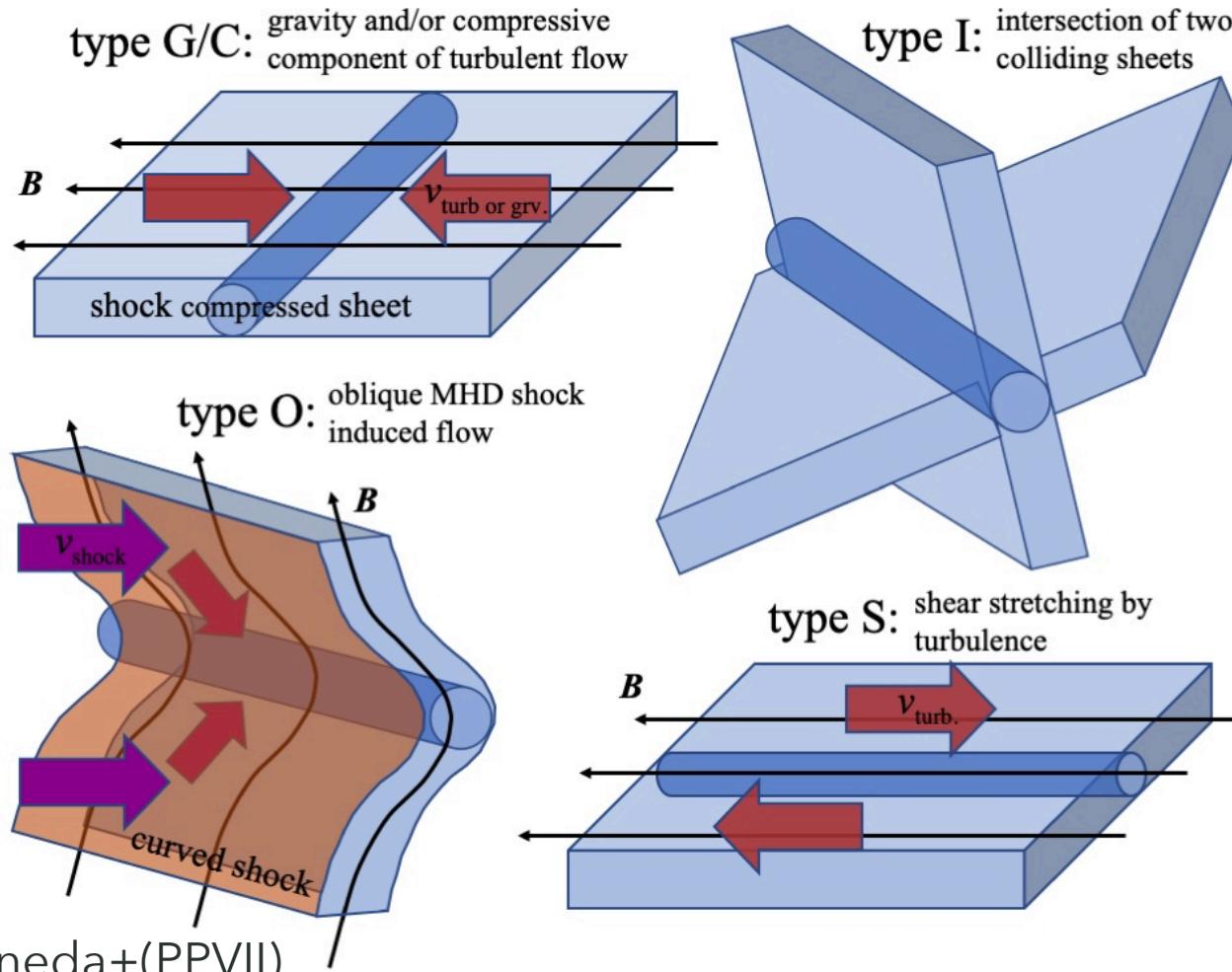
STORM

Slow-shock-mediated Turbulent flow Reinforced by Magnetic diffusion

can be important in determining
the initial conditions for star formation.

Filament Formation

There are 5 types of formation mechanisms

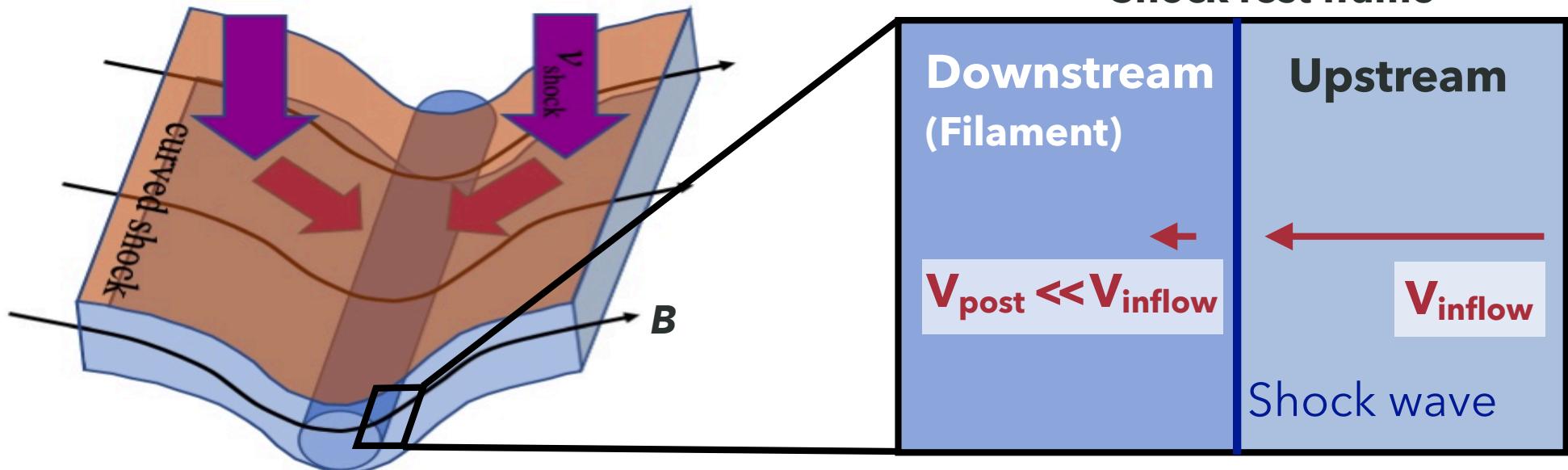


Abe+(2021), Pineda+(PPVII)

Formation process is well understood!

To Solve the 0.1 pc Problem

Pineda+(PPVII, 2023)



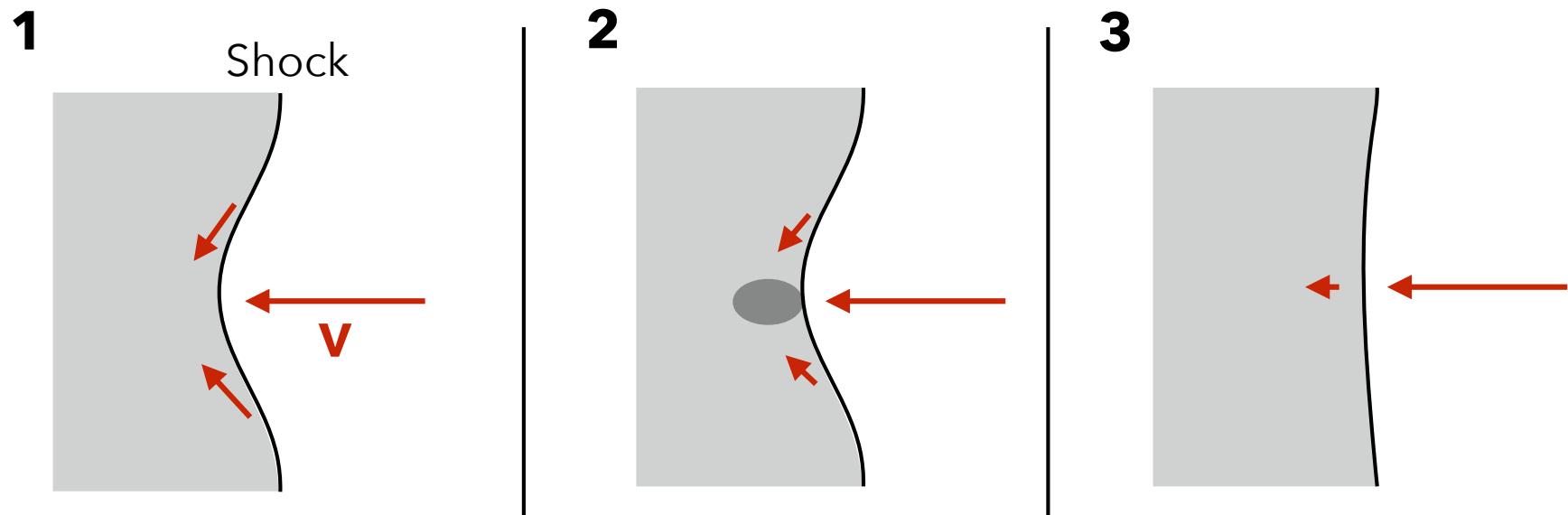
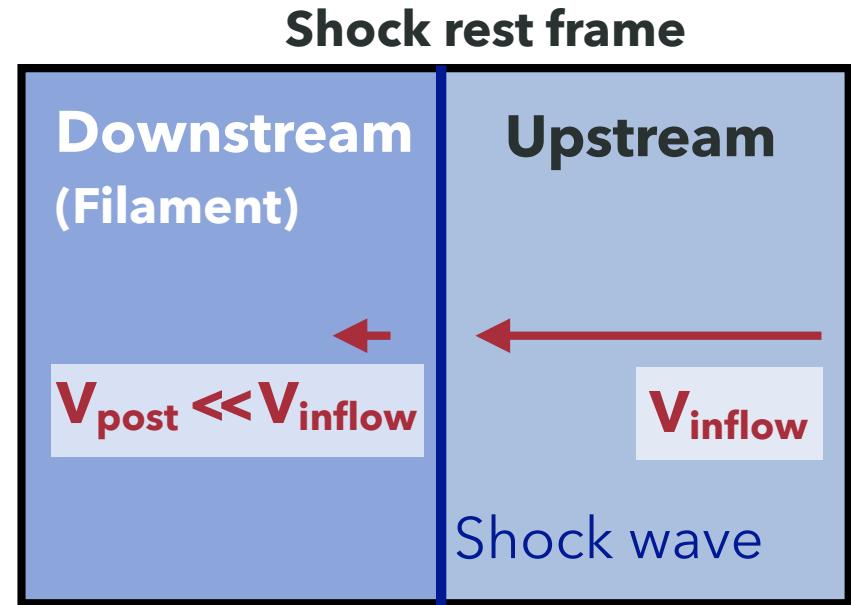
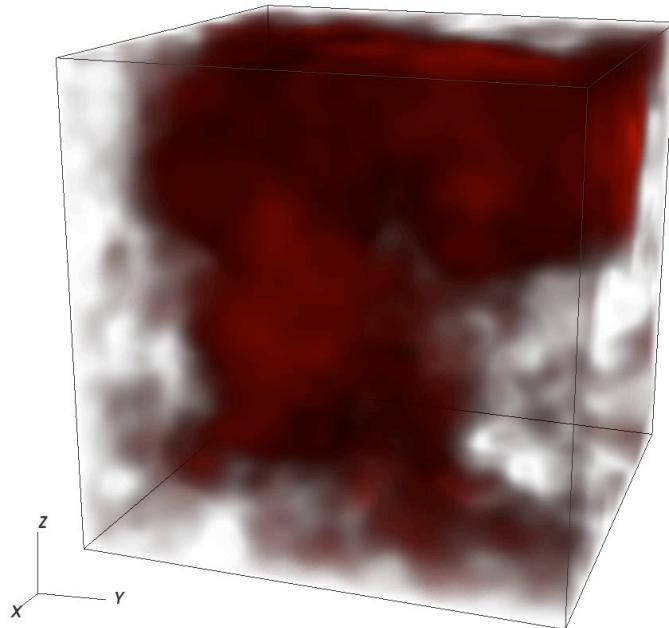
Shock waves bound the filament \rightarrow **dissipation** of the kinetic energy

Non-thermal pressure is necessary to keep the filament width.

$$M_{\text{line,crit}} = 2c_s^2/G \rightarrow 2(c_s^2 + \sigma_{\text{NT}}^2)/G$$

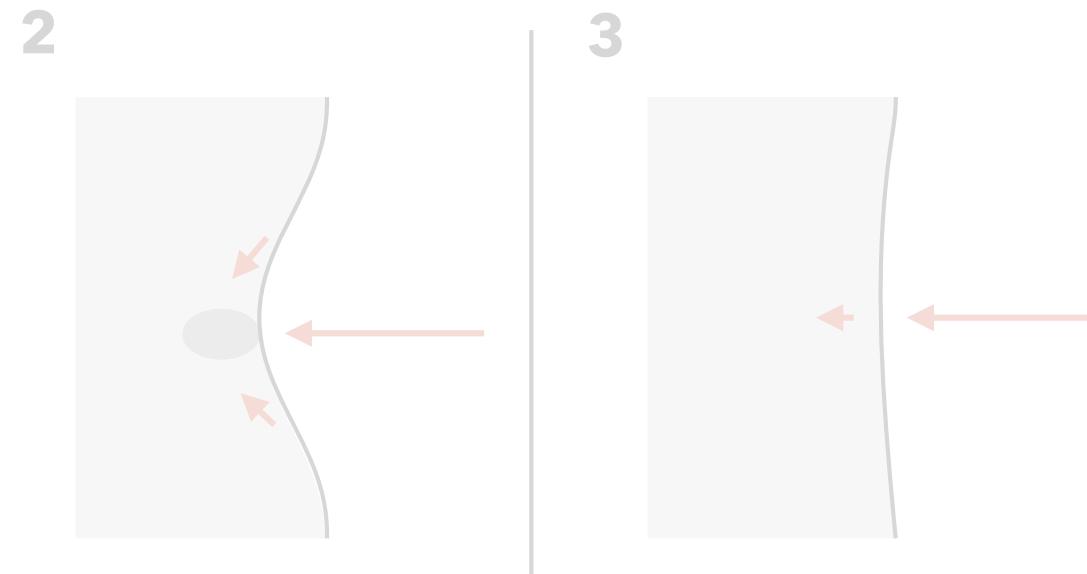
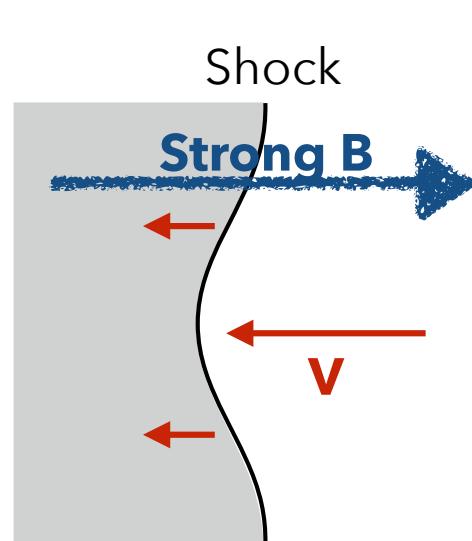
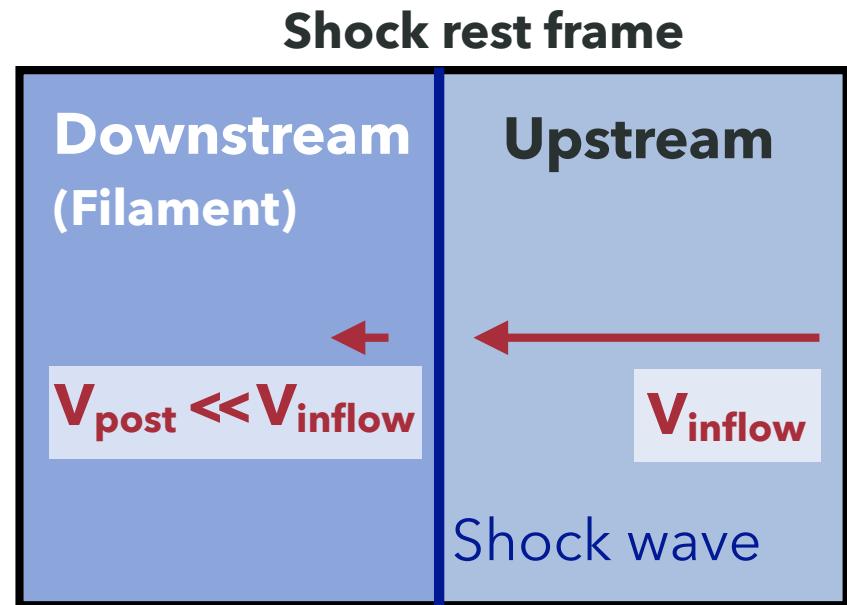
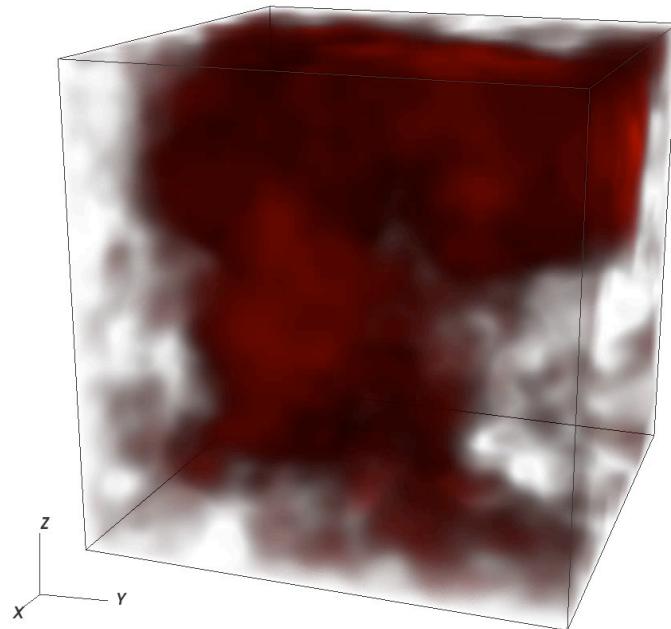
Non-thermal velocity dispersion

Remarks: Stability of Shock Fronts



(Usual) Shock fronts are stable to perturbations in their shape.

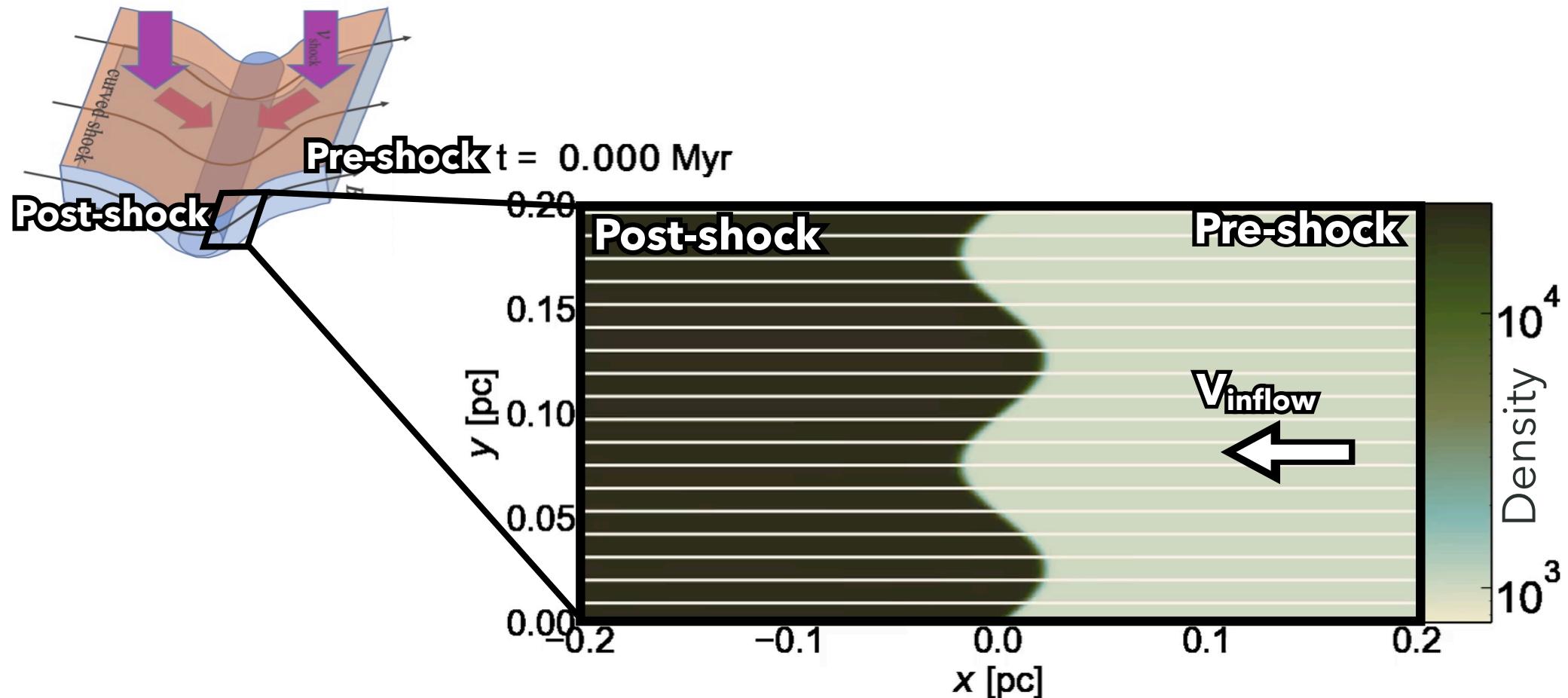
Brief Explanation of Slow-shock Instability



Accurate mechanism is different but brief understanding is OK with this.

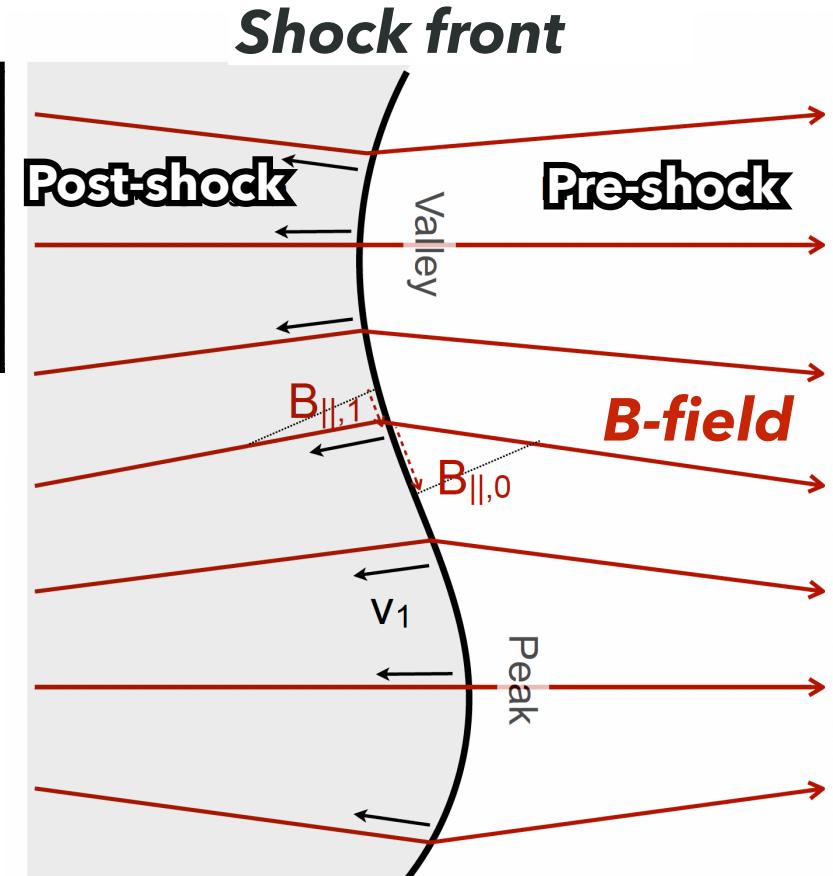
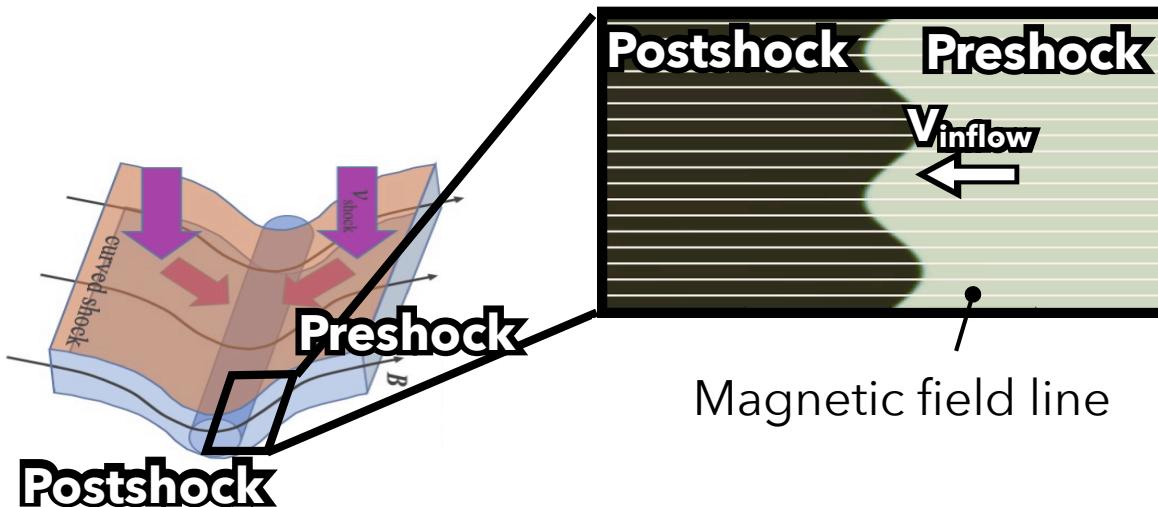
Key Process: Instability at Filament Boundaries

Slow-shock Instability (SSI) (e.g., Lesson & Deshpande 1967)



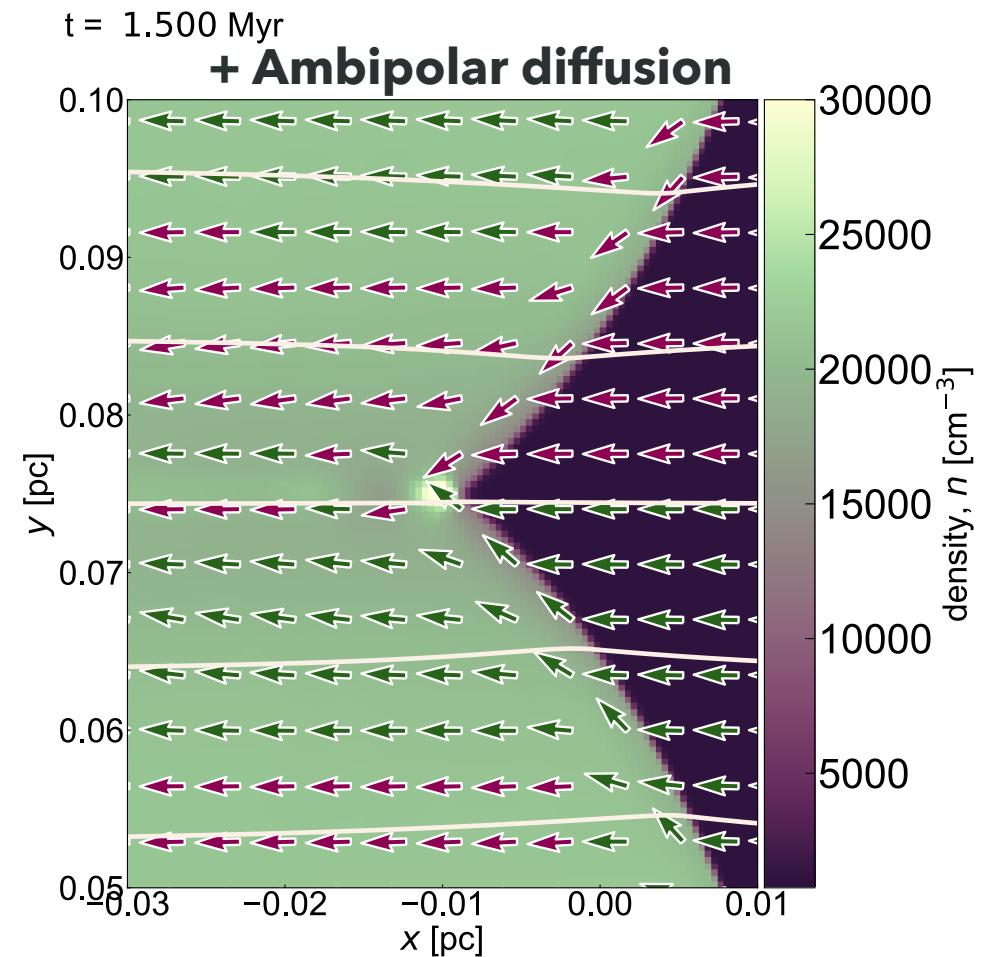
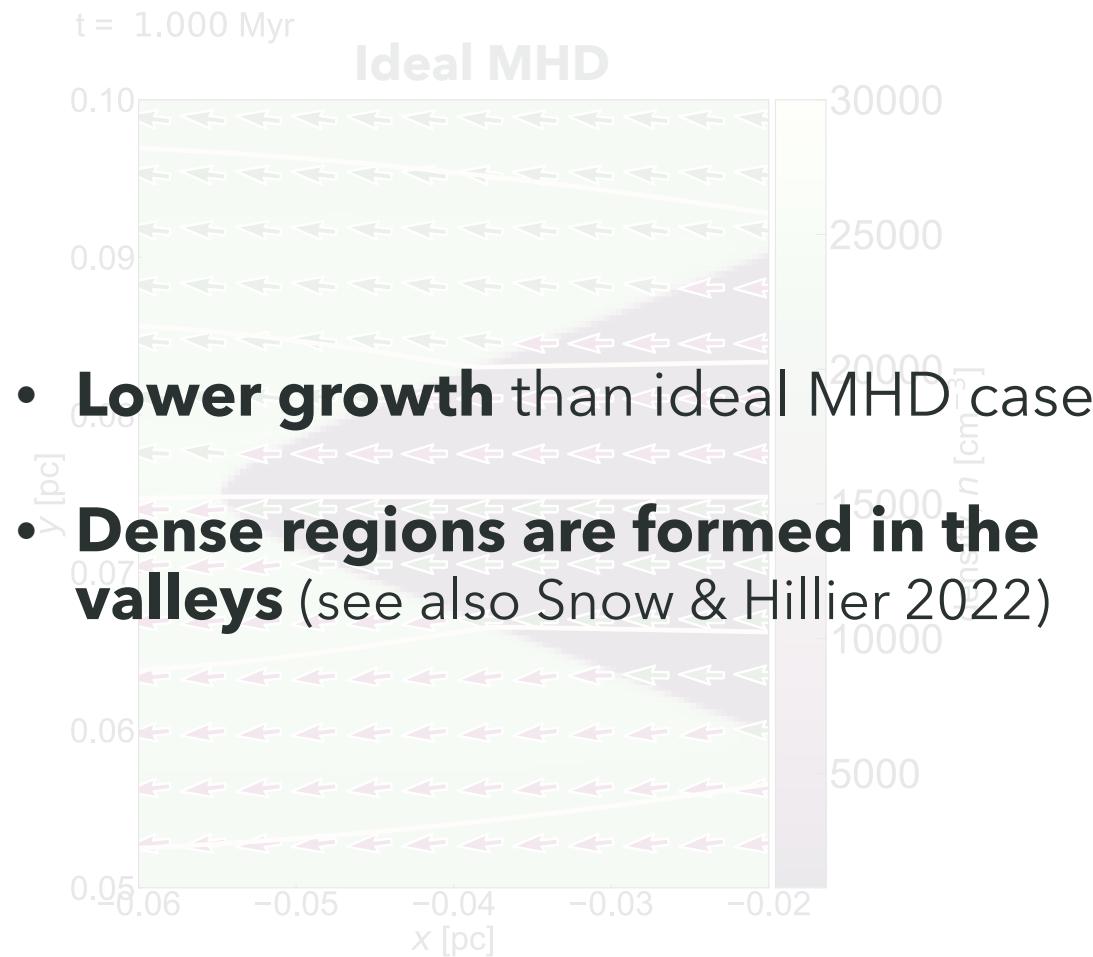
Linear Growth of the SSI

Slow-shock Instability (SSI) (e.g., Lesson & Deshpande 1967)



$B_{||}$ becomes small

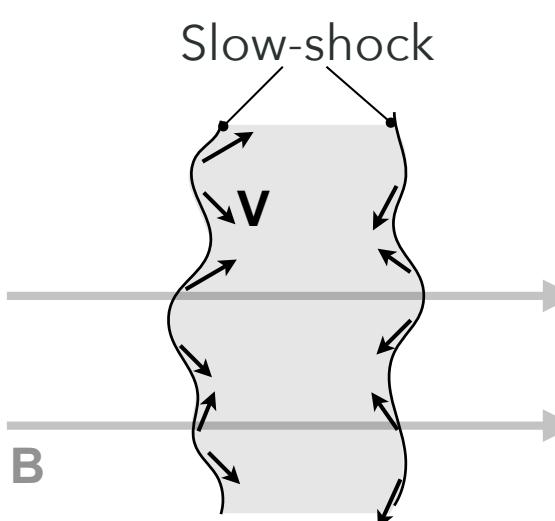
Nonlinear evolution of SSI + ambipolar diffusion



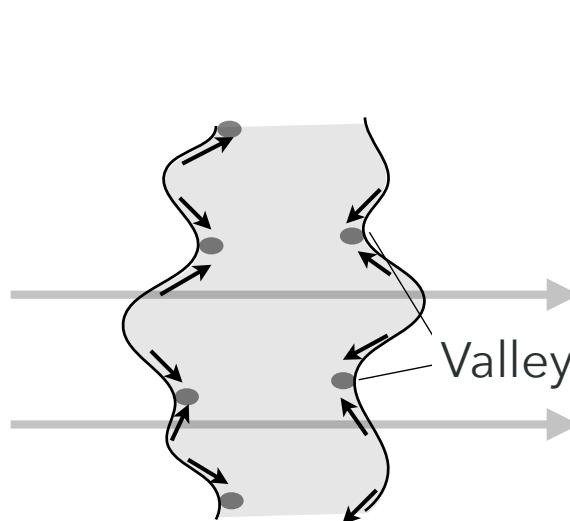
The Origin of Additional Pressure: “STORM”

STORM

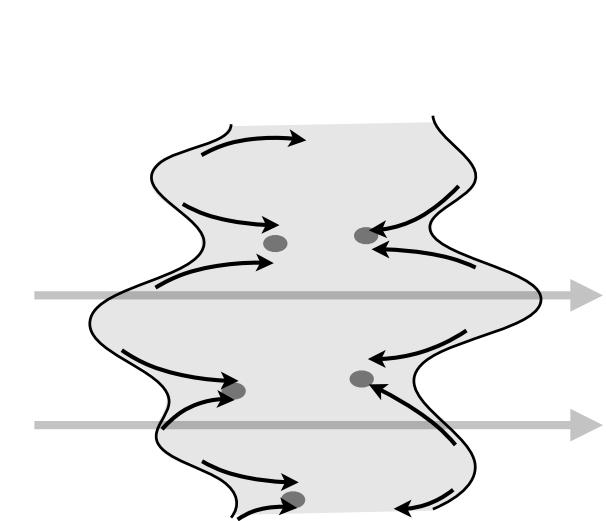
(Slow-shock-mediated **T**urbulent **f**lOw **R**einforced by **M**agnetic diffusion)



1. AD permits **flow across the magnetic field** lines



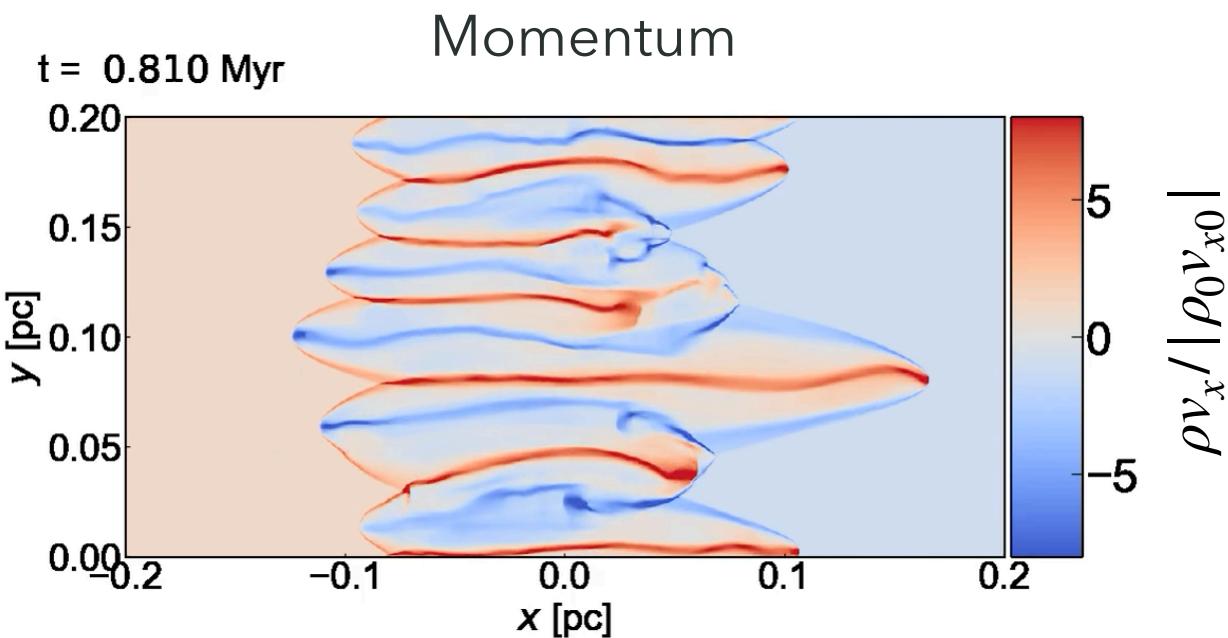
2. Gas accumulates in the valleys



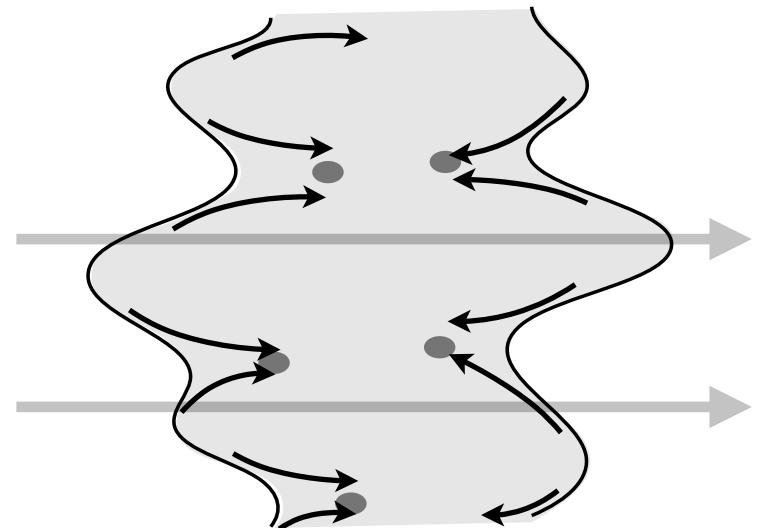
3. **Momentum injection** induces expansion of the shock-bound region

The corrugation of shock fronts continues to grow under the effect SSI.

The Origin of Additional Pressure: “STORM”



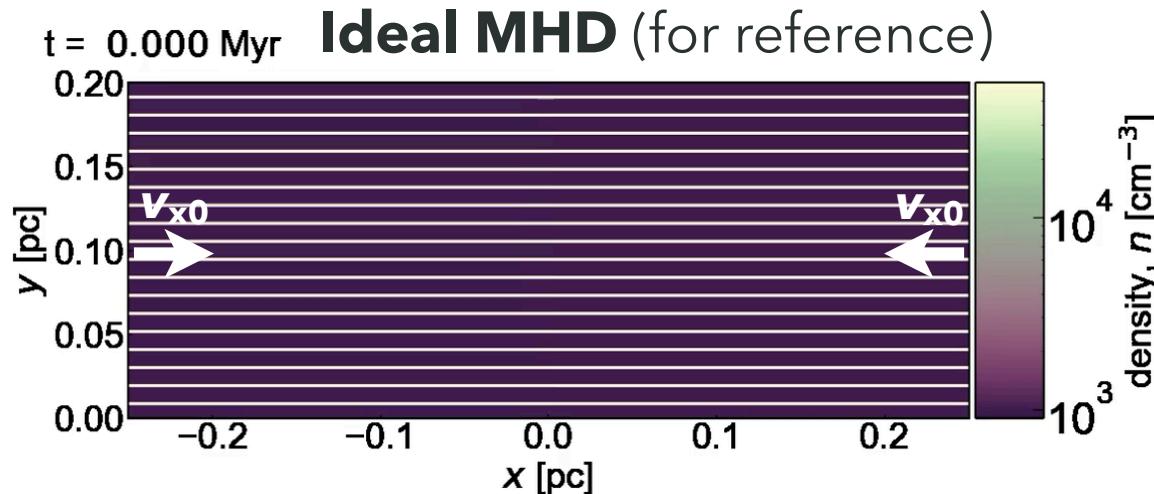
The blobs stir the filament gas like “hailstorm”.



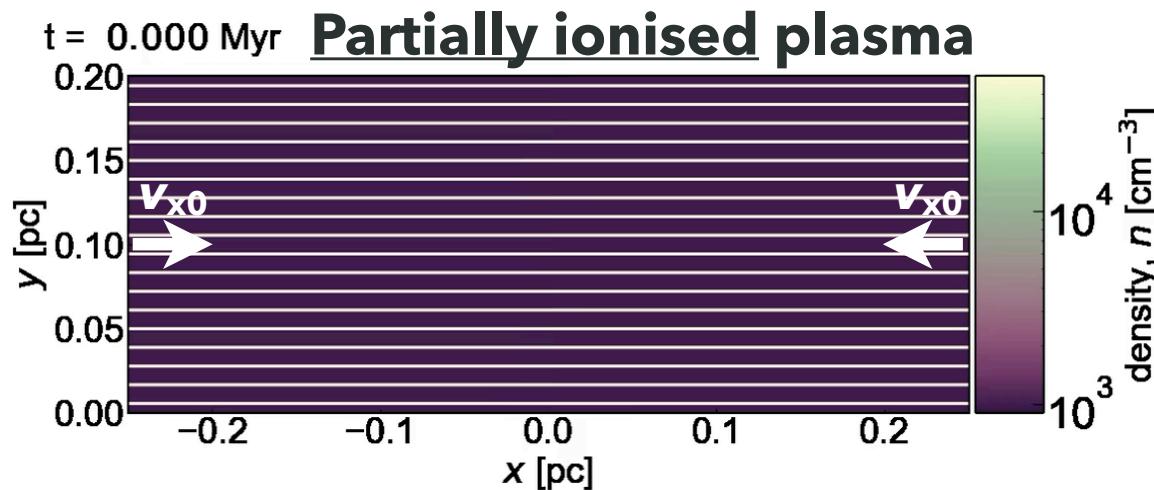
“STORM” provides ram pressure to increase filament width

Result: Simulation of Filament Evolution with Ambipolar Diffusion

Evolution of density field



- SSI causes shock front corrugation.
- **No turbulence**



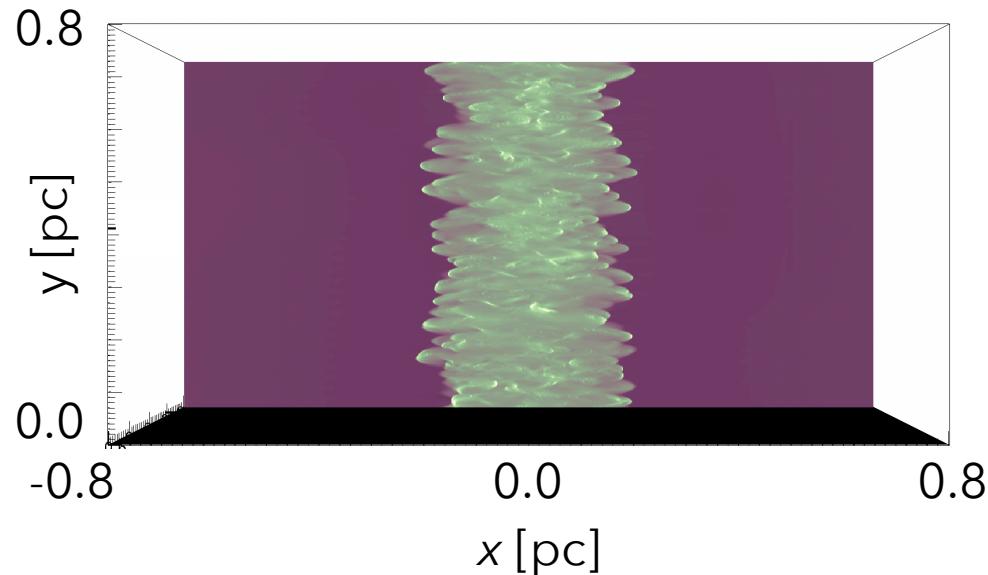
More turbulent than in ideal MHD

Importance of Resolution

Density Slice

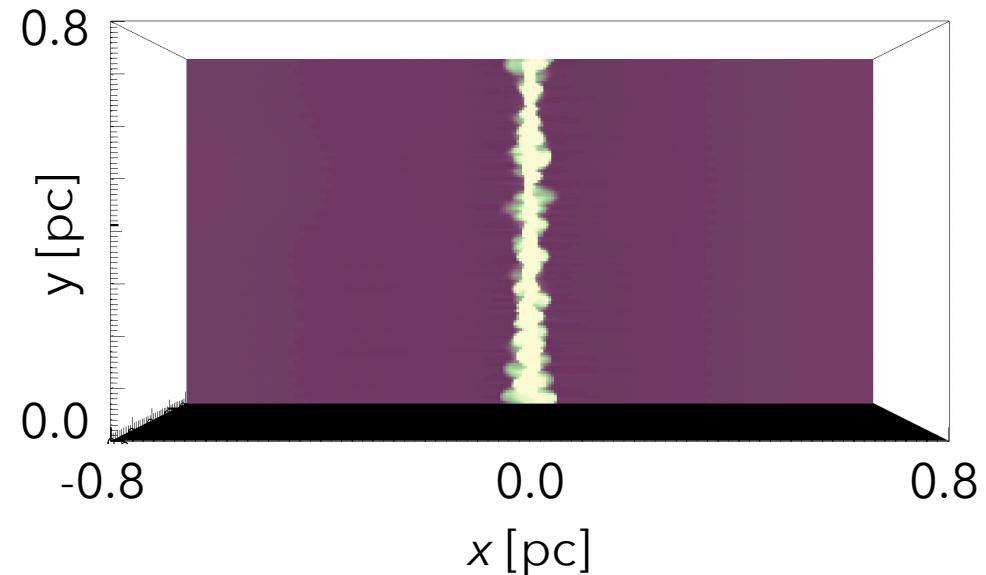
High resolution

Finest resolution: $\Delta x \sim 1.5 \times 10^{-3}$ pc



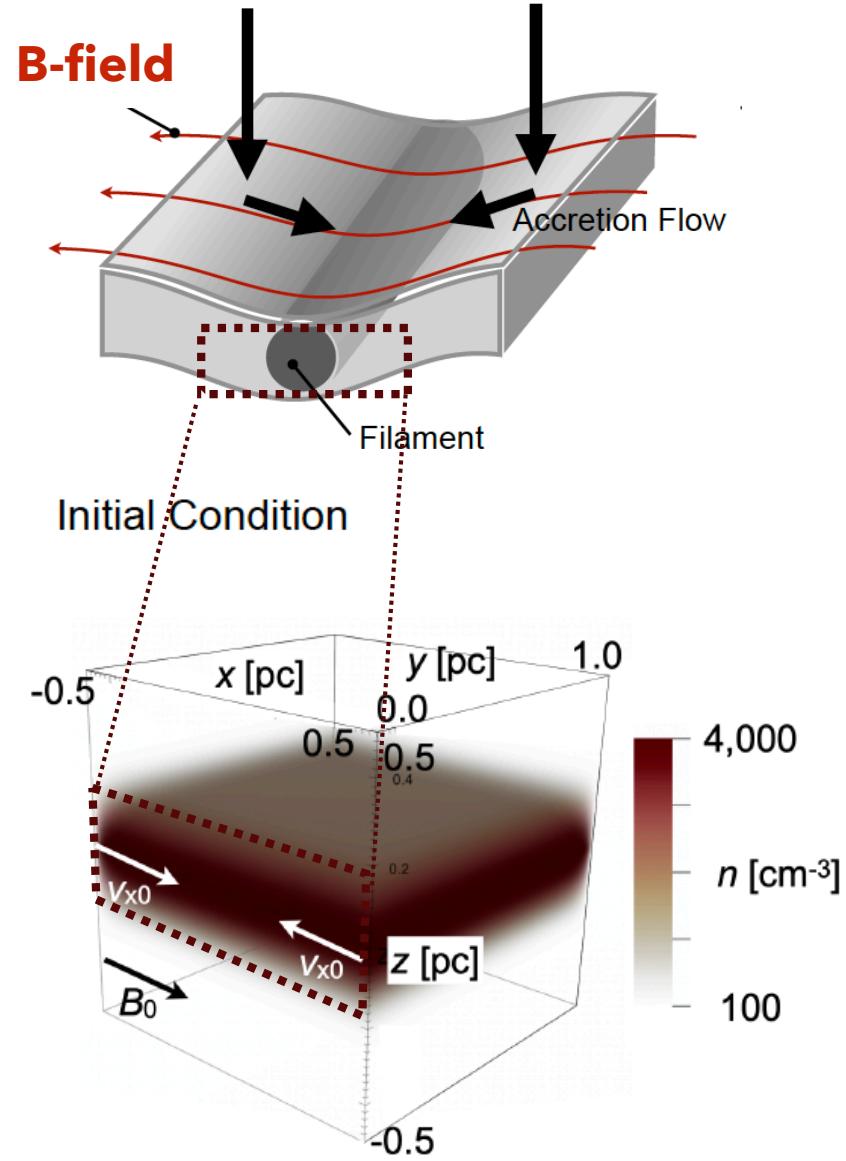
Low resolution

$\Delta x \sim 6.2 \times 10^{-3}$ pc

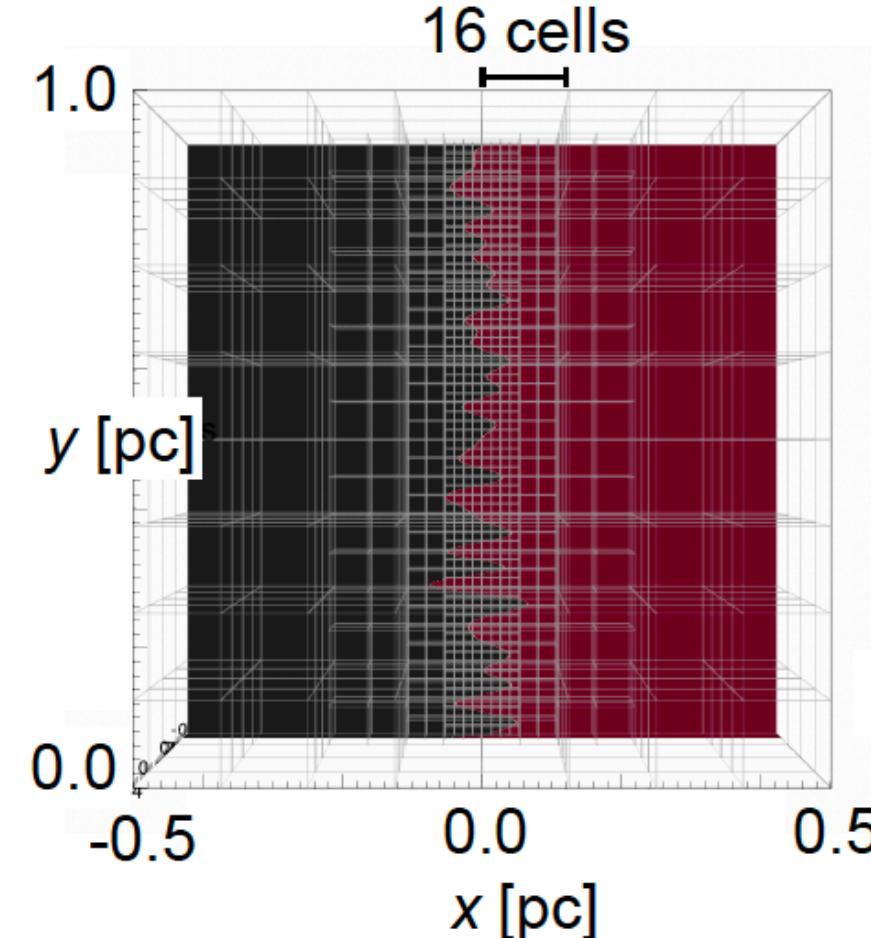


We need **high resolution**
to appropriately solve the filament evolution

Setup for 3D Simulation

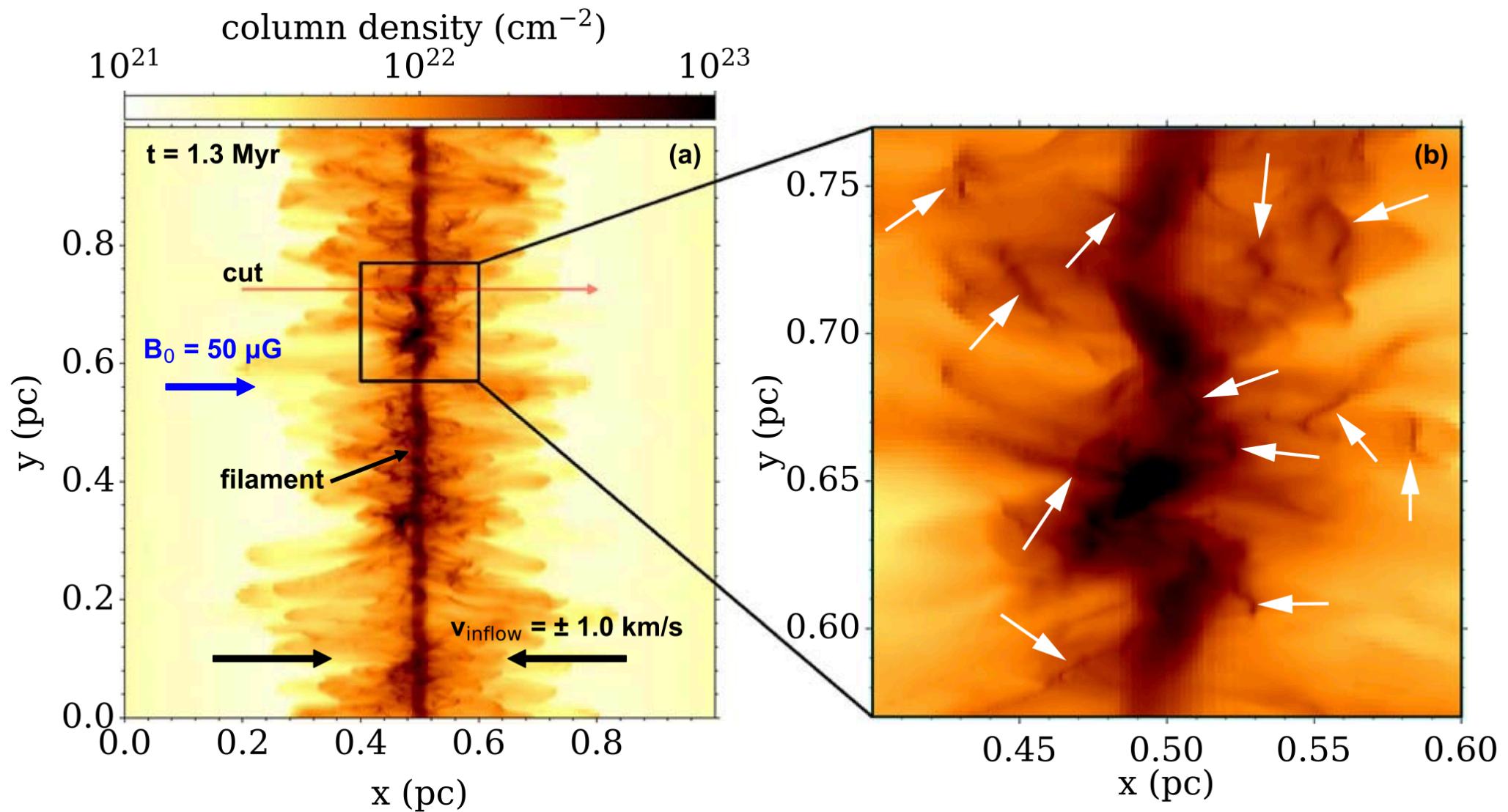


Velocity and grid structure

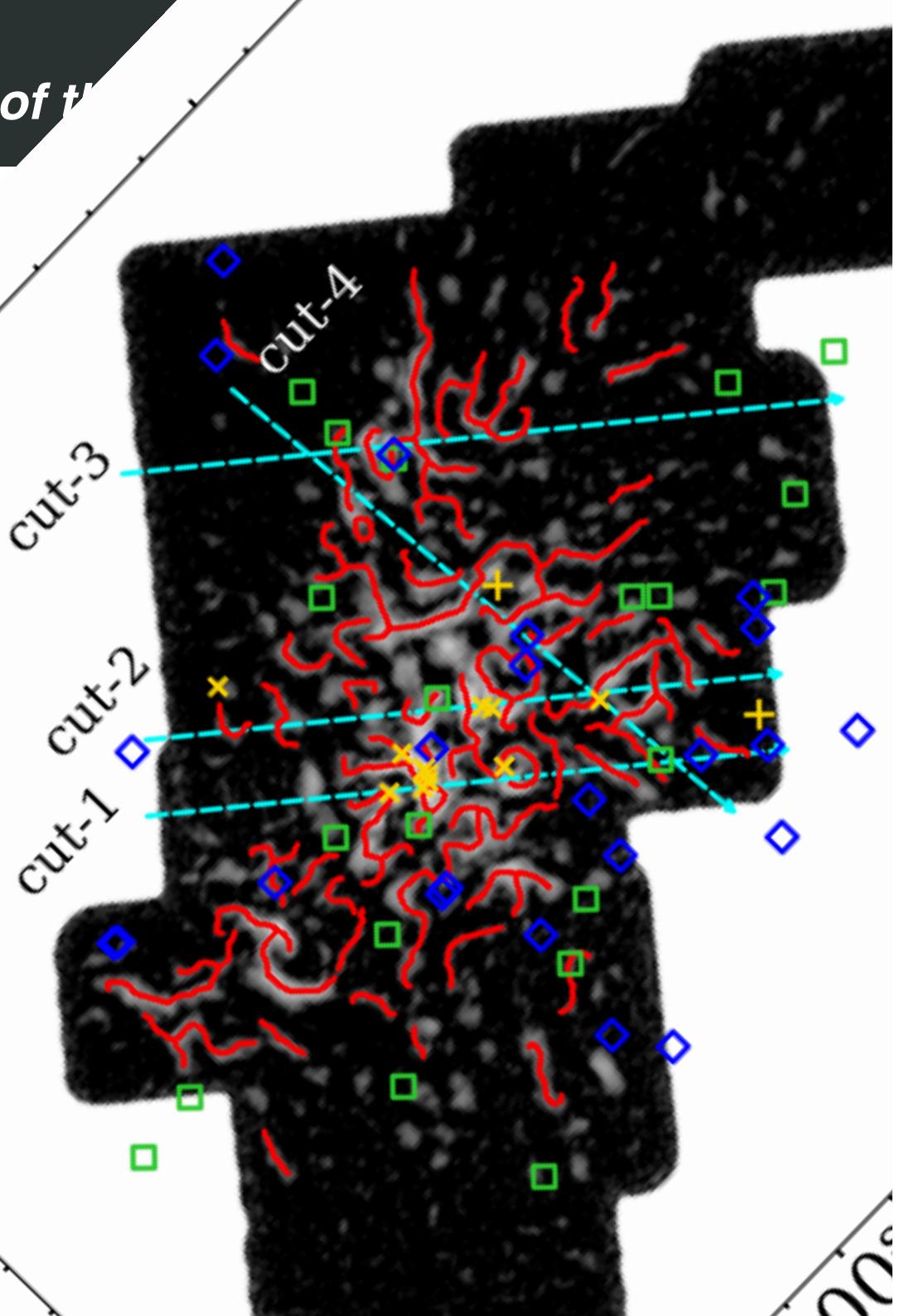
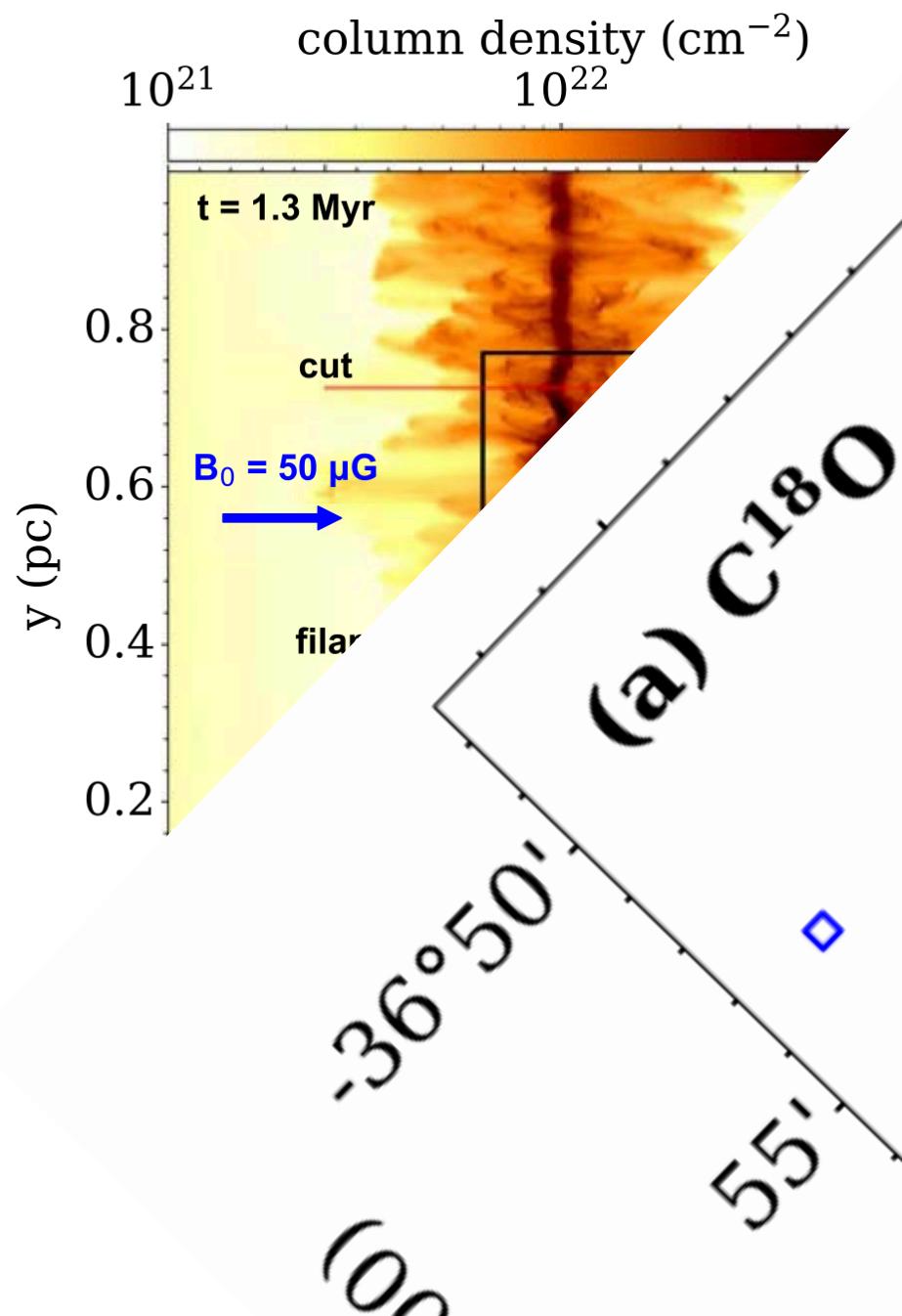


Use hierarchical grids
Finest resolution: ~ 0.001 pc

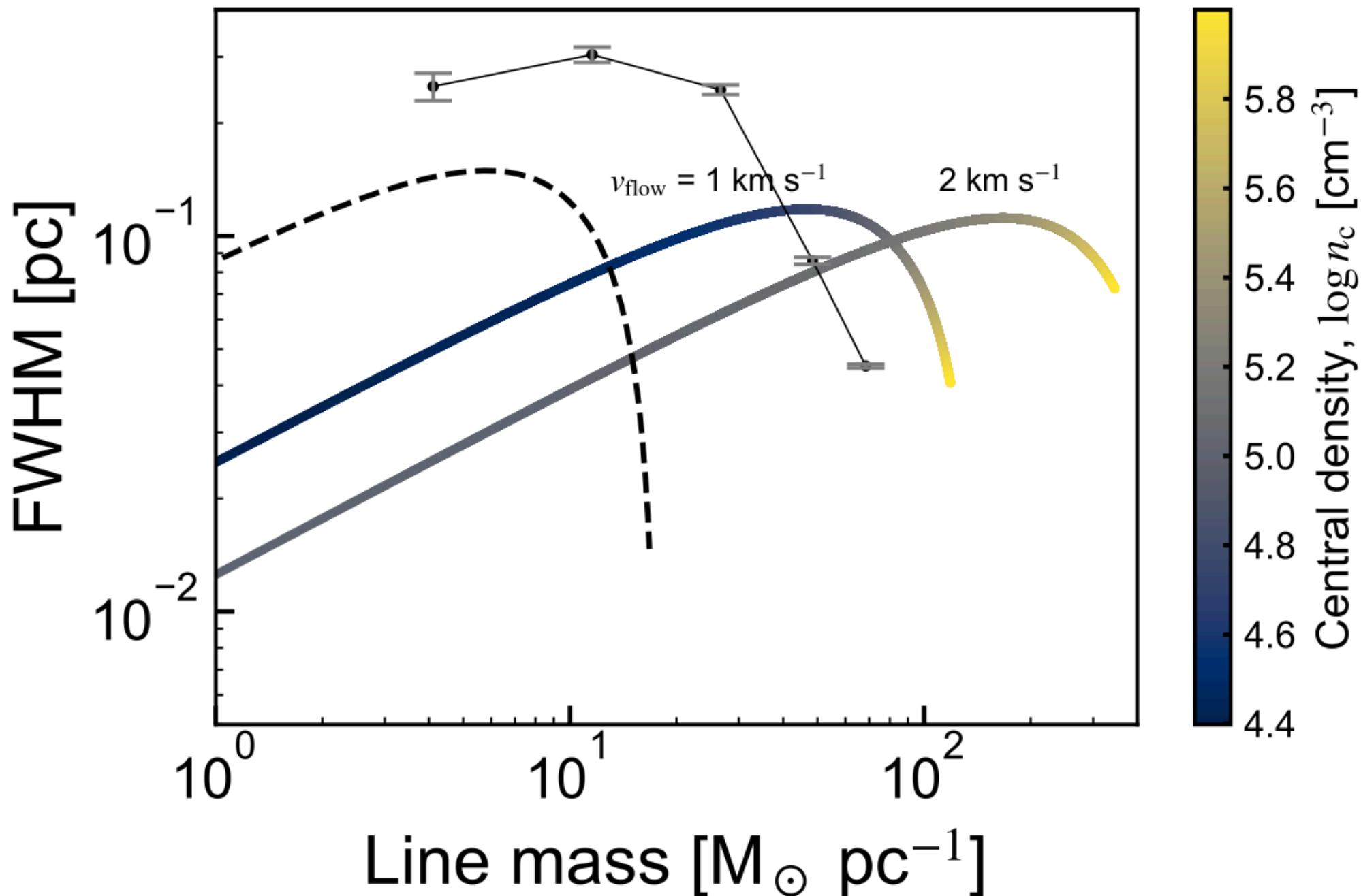
Internal 1000 au Scale Structures of the R CrA Cluster-forming Cloud



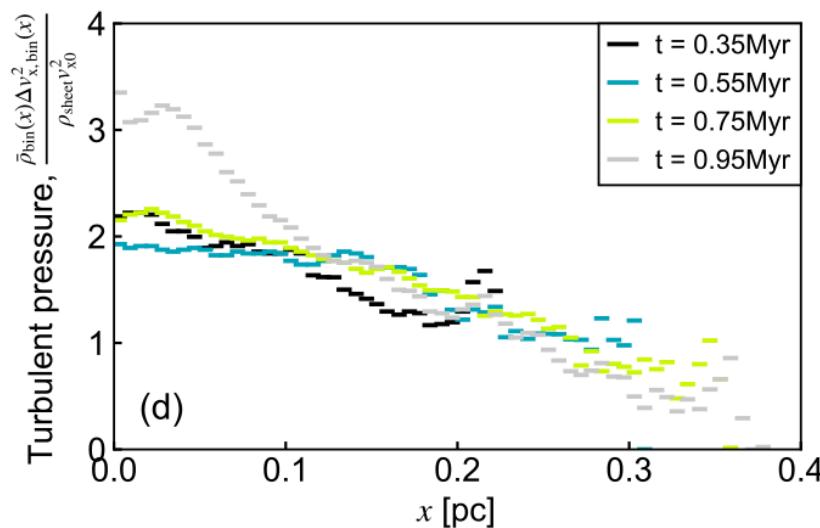
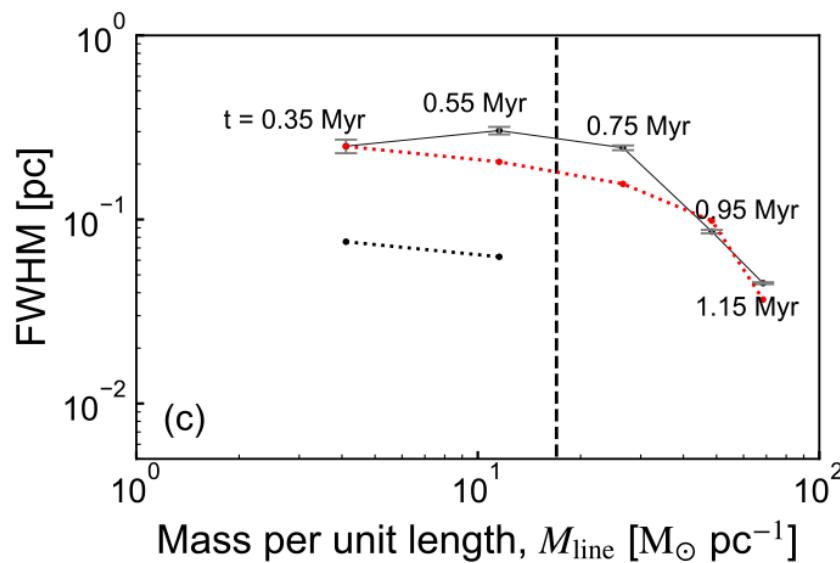
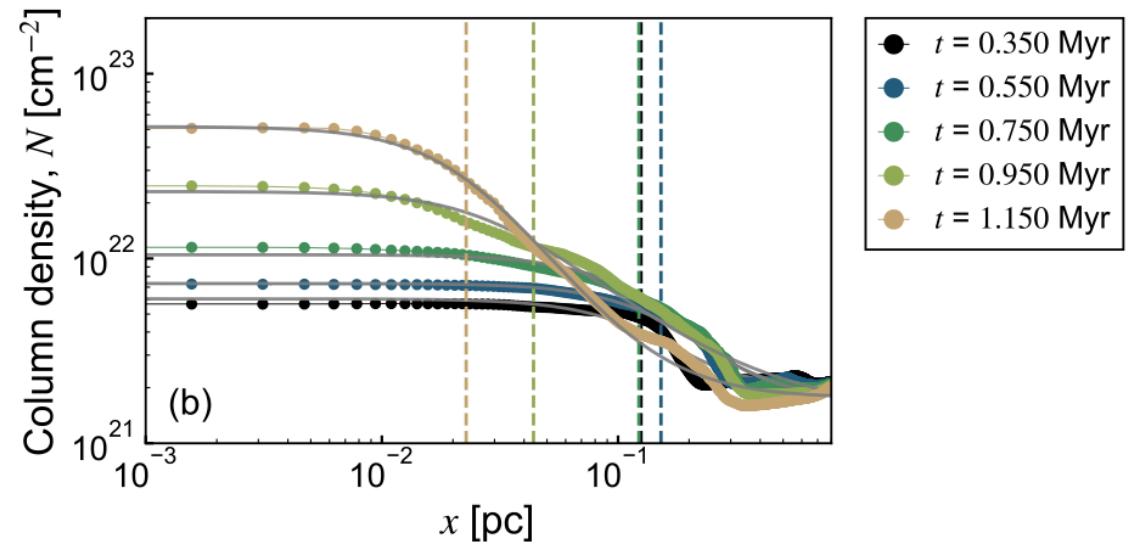
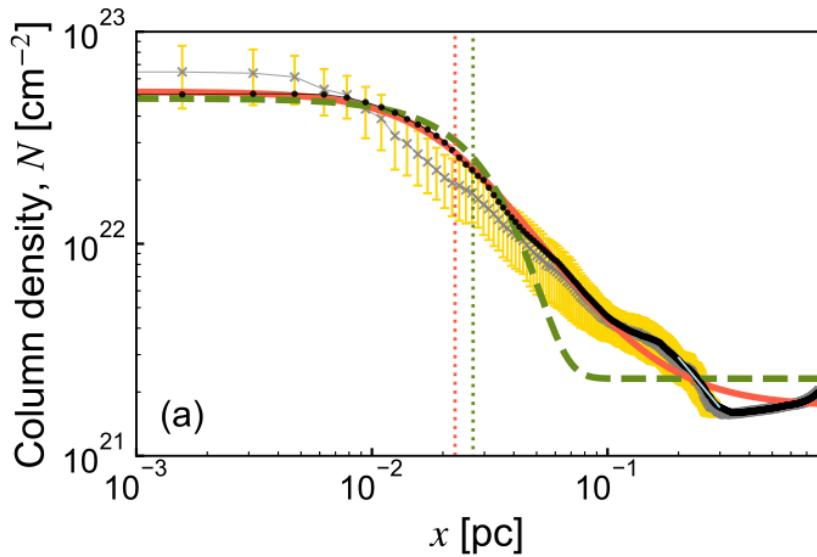
Internal 1000 au Scale Structures of the Taurus Molecular Cloud



Analytic Model considering the *STORM* mechanism

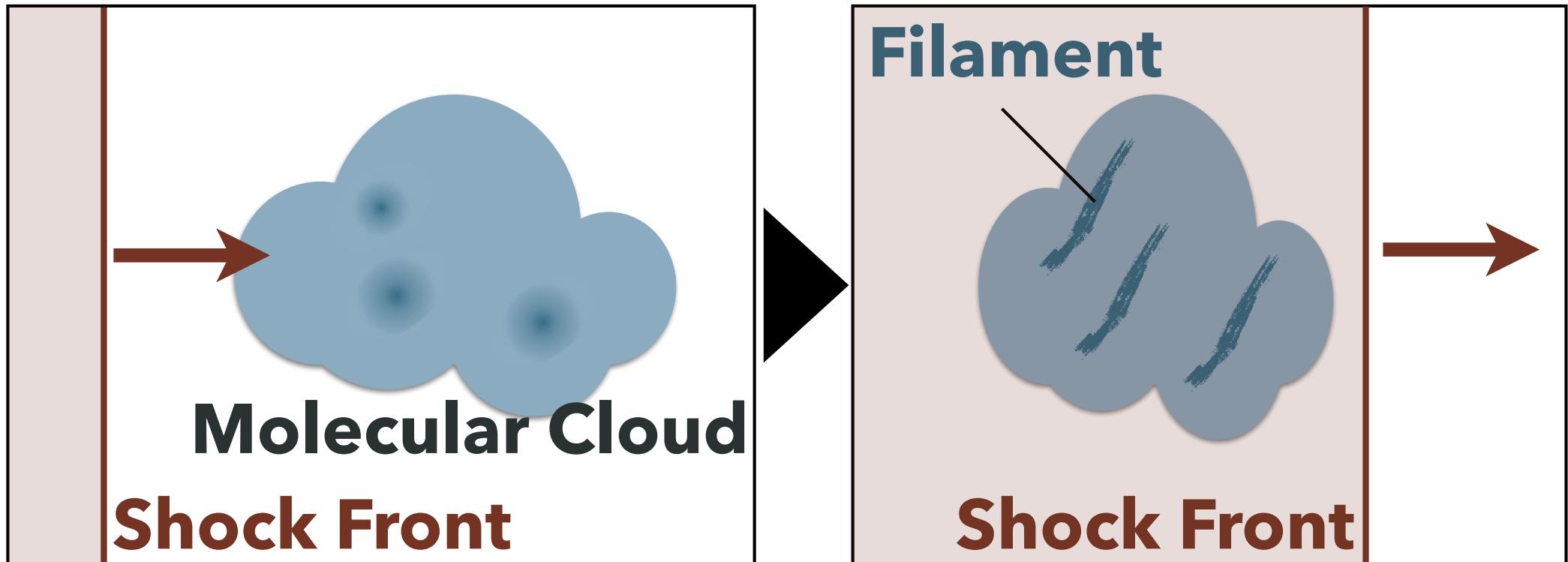


Properties of the STORM filament



Trigger of Filament Formation

Molecular clouds are **frequently** ($\sim 1 \text{ Myr}^{-1}$) swept by **shock waves**.
(ex. Super Novae, Cloud-cloud Collision, Expansion of HII regions etc.)

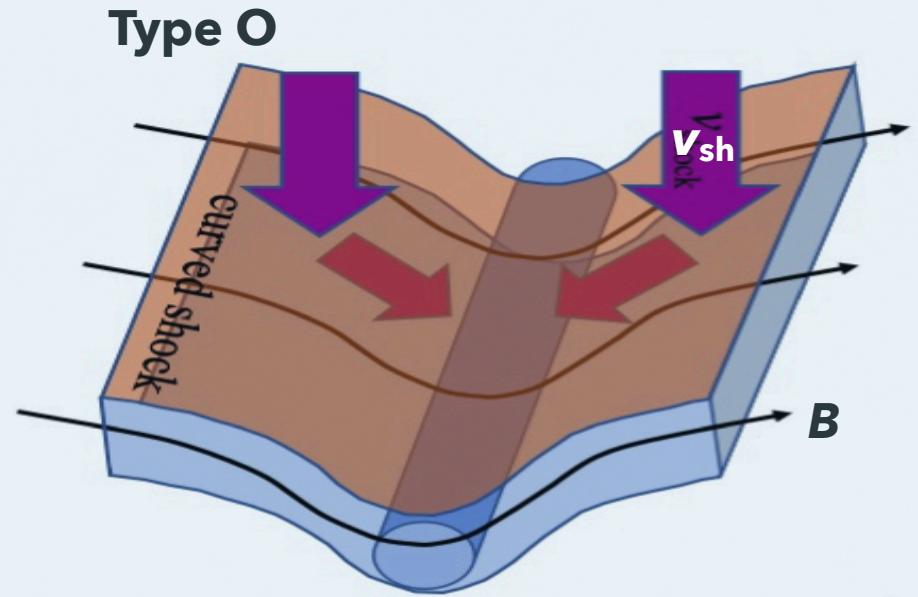


Shock waves trigger **filament formation**.

Formation Mechanisms of Filaments

Main filament formation mechanisms

- Type G: Gravitational fragmentation of shocked sheet
(e.g., Nagai+ 1998)
- Type C: Compressive component of turbulent flow
(e.g., Matsumoto+ 2015)
- **Type O**: Oblique MHD shock induced flow
(Inoue & Fukui 2013)

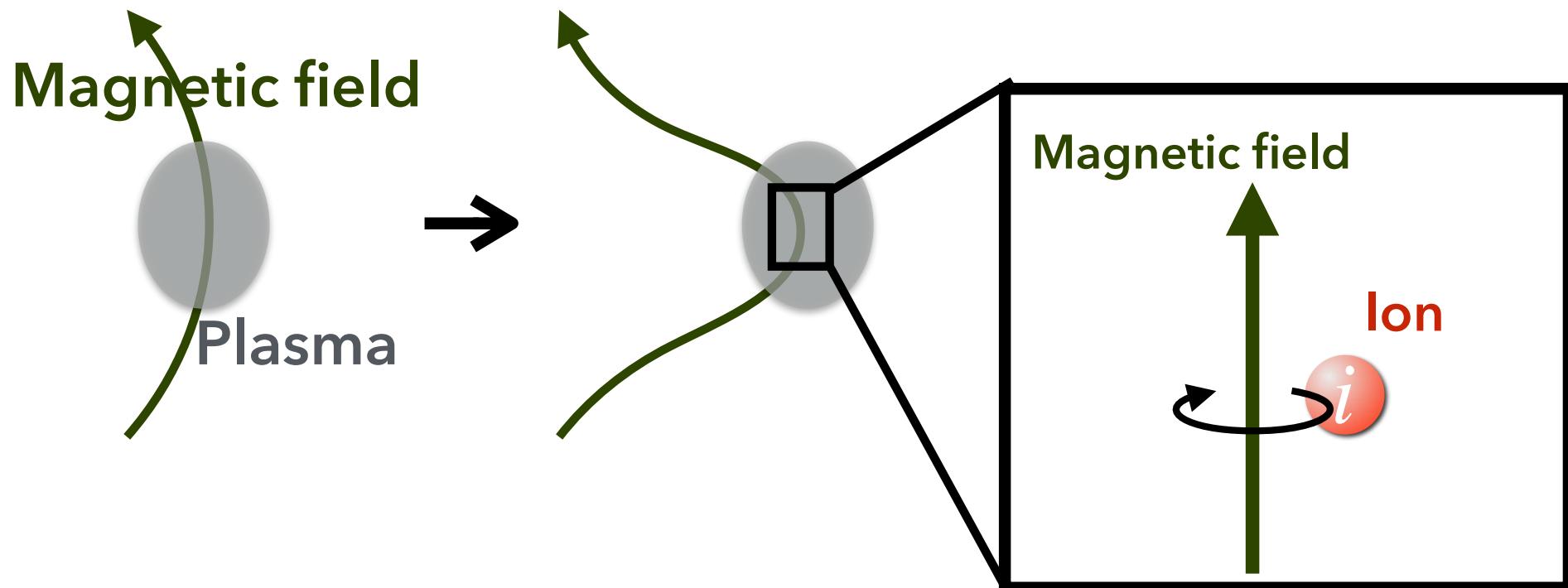


Abe+2021,
see also PPVII review (Pineda+ 2023)

Filament formation is well understood!

(Abe, Inoue, Inutsuka & Matsumoto 2021)

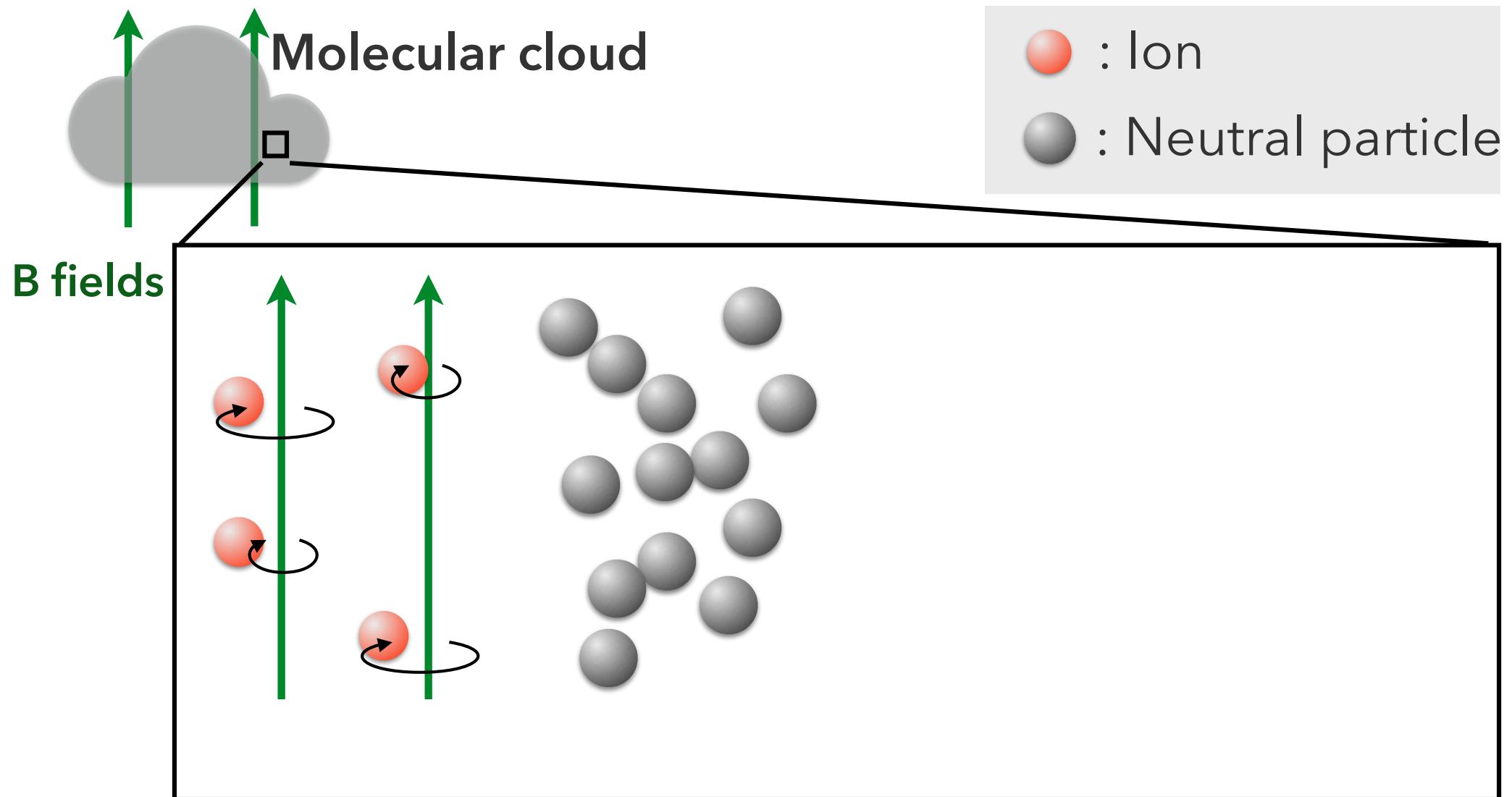
Properties of Magneto-Hydro Dynamics (MHD)



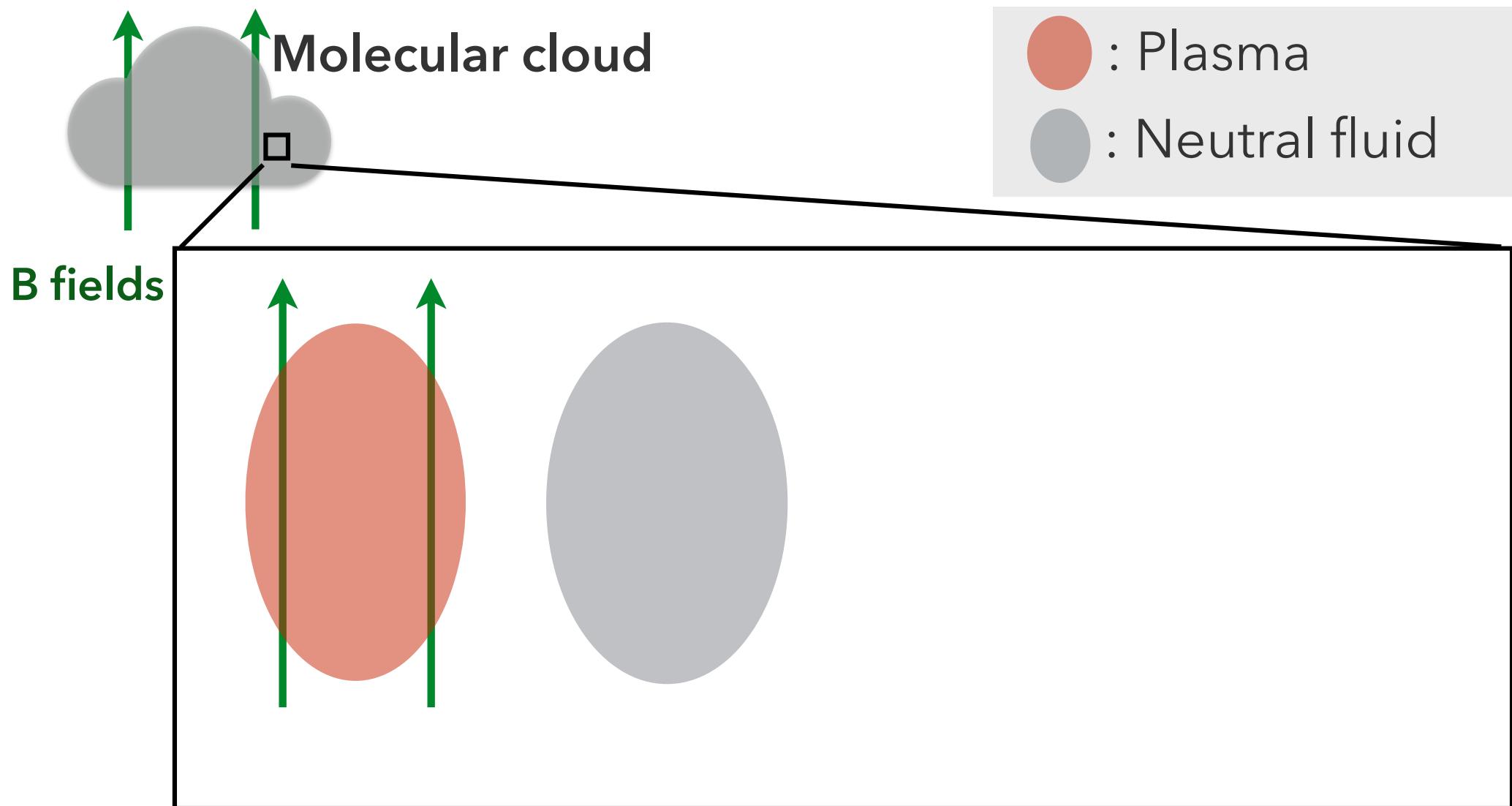
- Plasma can move easily along magnetic fields.
- Plasma drags magnetic field
- Plasma cannot move across magnetic fields.

“Magnetic frozen-in”

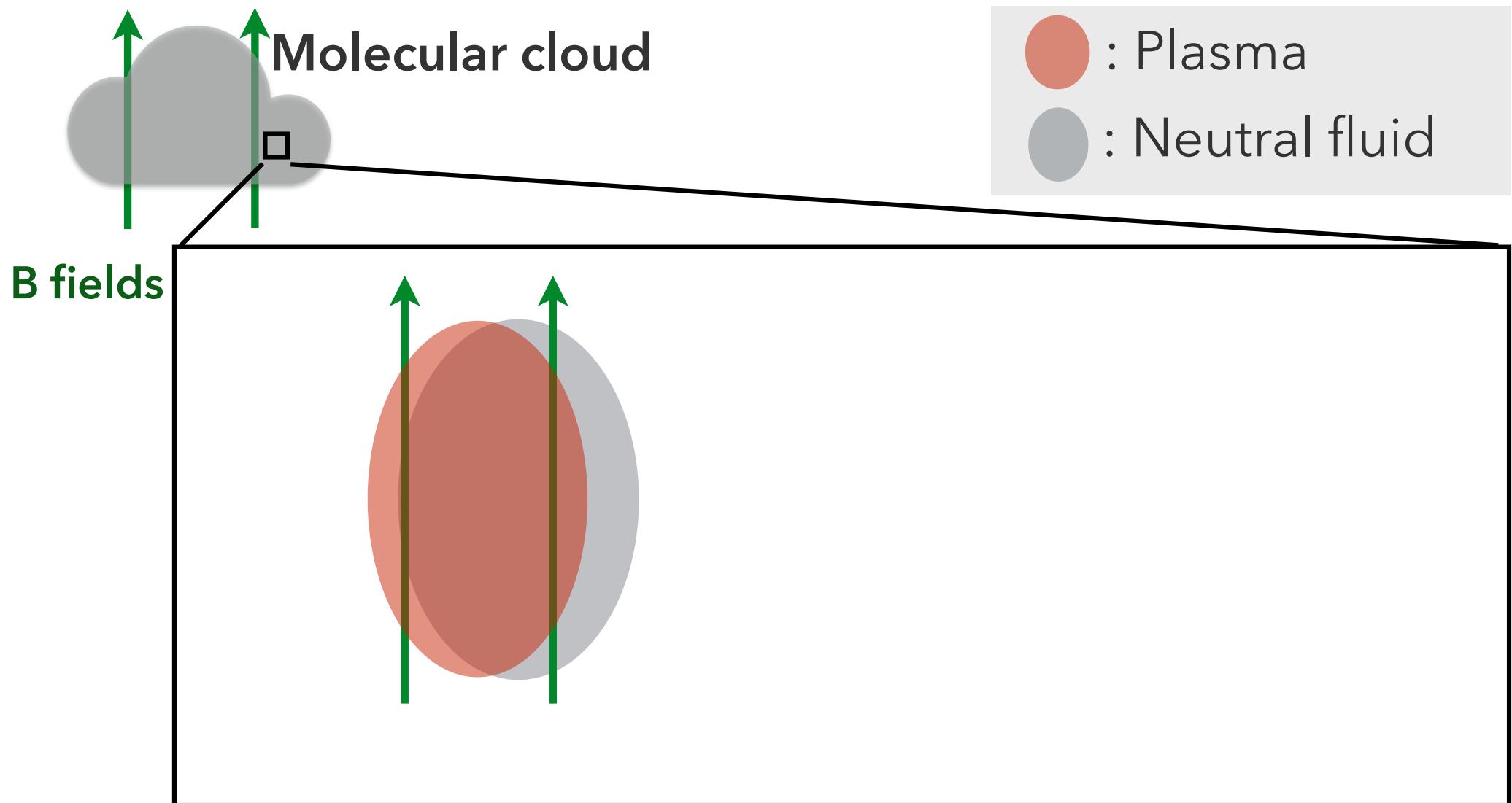
MHD Approximation for Molecular Clouds Dynamics



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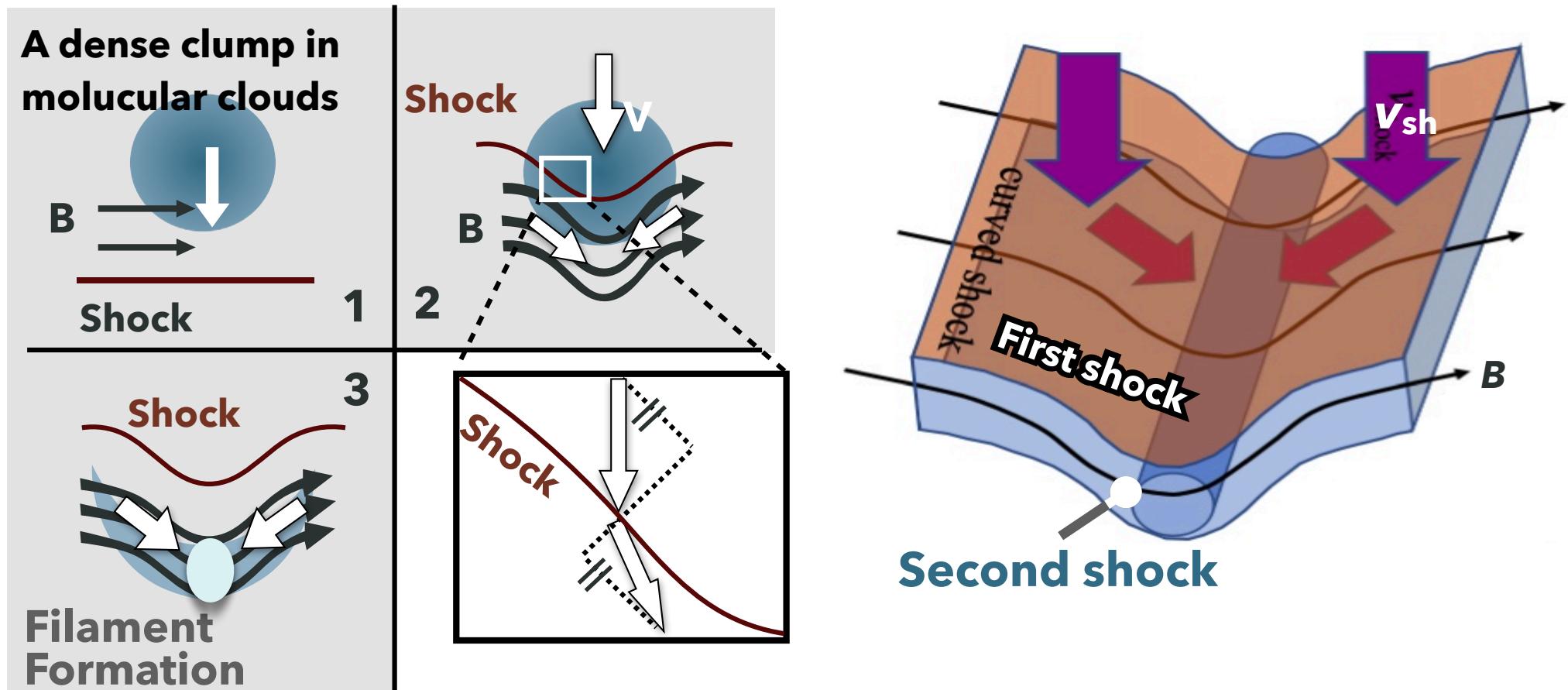
MHD Approximation for Molecular Clouds Dynamics



The dynamics of molecular clouds can be described as MHD.

An example of filament formation mechanism (Type O)

Interstellar shock wave with velocity of $\gtrsim 5$ km/s \rightarrow Type O
(Abe+ 2021)



Two stages of compression \rightarrow Filament formation

Importance of Filament Width

Condition of Star Formation

Statement in the Filament Paradigm

Stars form in filaments
above **critical line mass**
(e.g., André+2010)

Which is correct?

Statement by Lada et al.

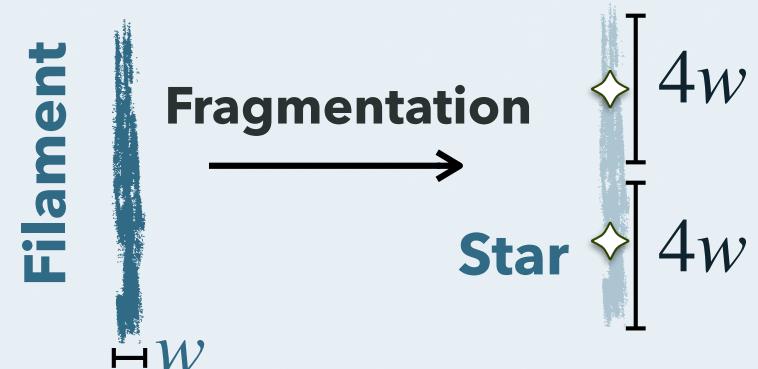
There is a threshold column density
($\sim 10^{22} \text{ cm}^{-2}$) for star formation to occur.
(e.g., Lada+2010)

Core Mass

(The most unstable mode of self-gravitational fragmentation)
 \sim (Filament width).

(Inutsuka & Miyama 1997)

→ The width determines
core mass



Physical origin of 0.1 pc width should be clarified

Importance of Filament Width

Condition of Star Formation

Statement in the Filament Paradigm

Stars form in filaments above **critical line mass**
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Critical line mass

$$2c_s^2/G$$

&

Width

0.1 pc

→ **Column density**

$$N \sim 10^{22} \text{ cm}^{-2}$$

Statement by Lada et al.

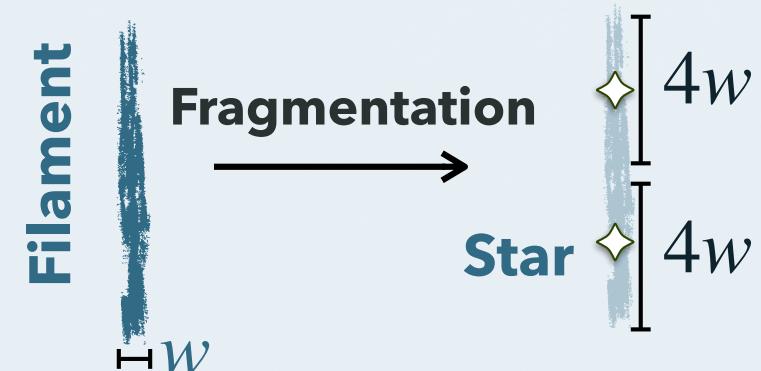
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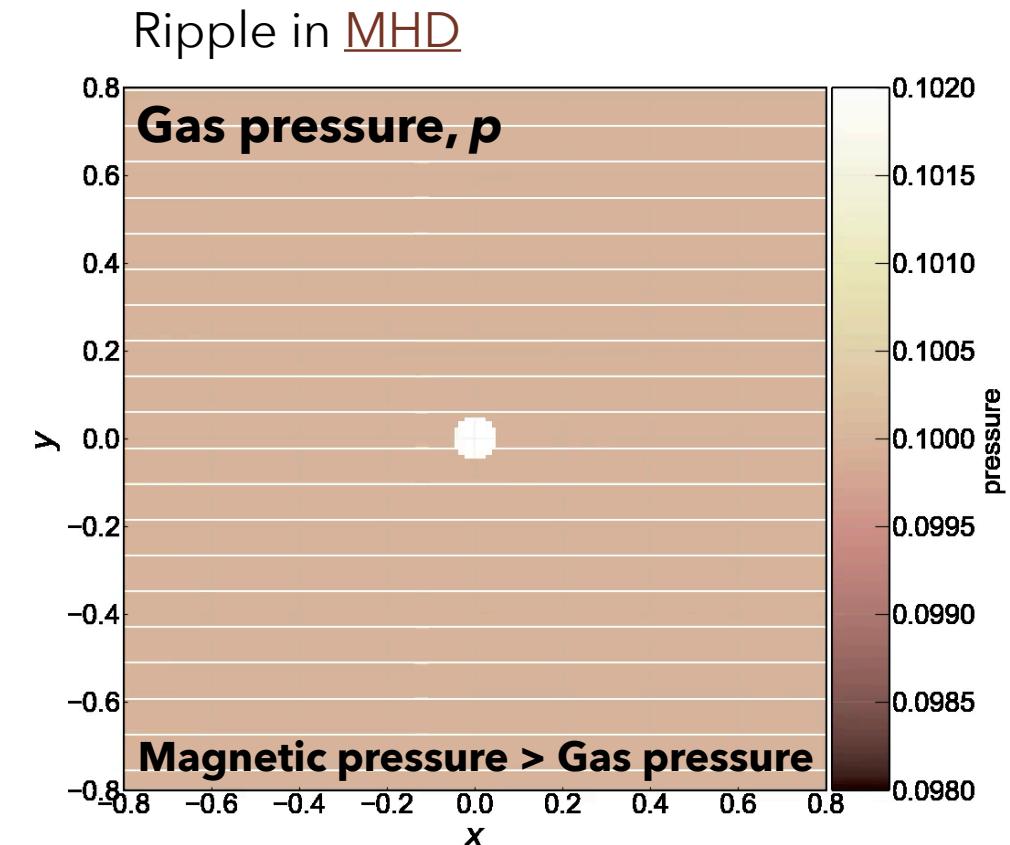
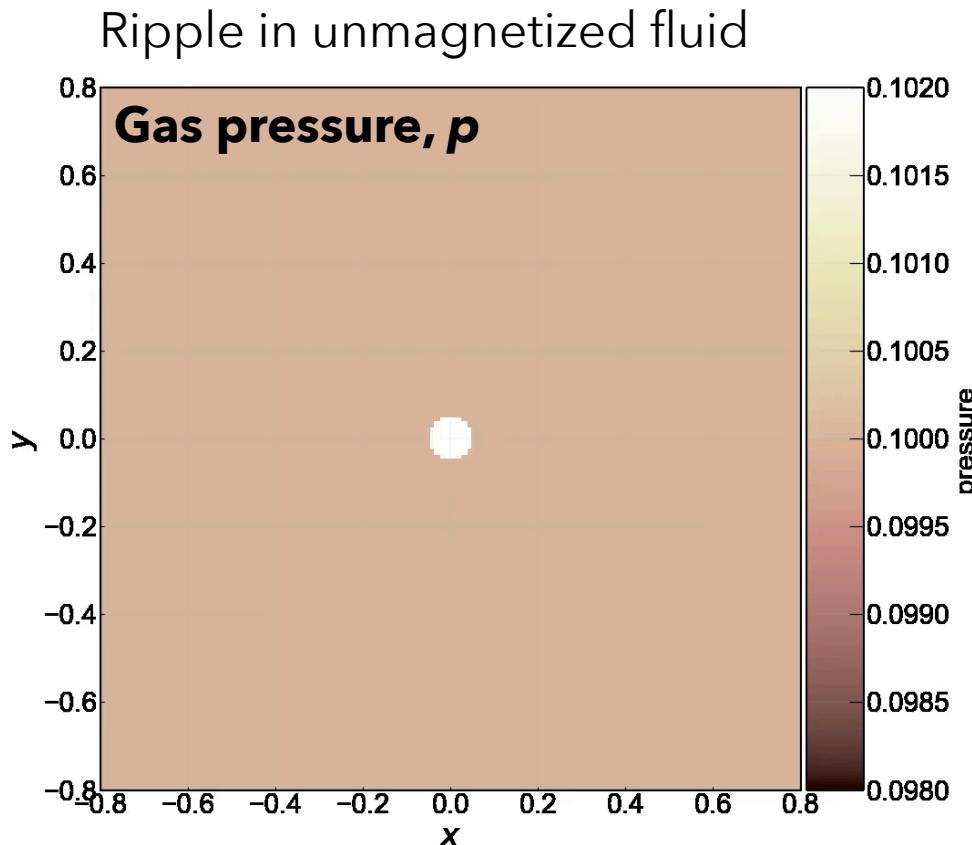
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Instability at Filament Boundaries

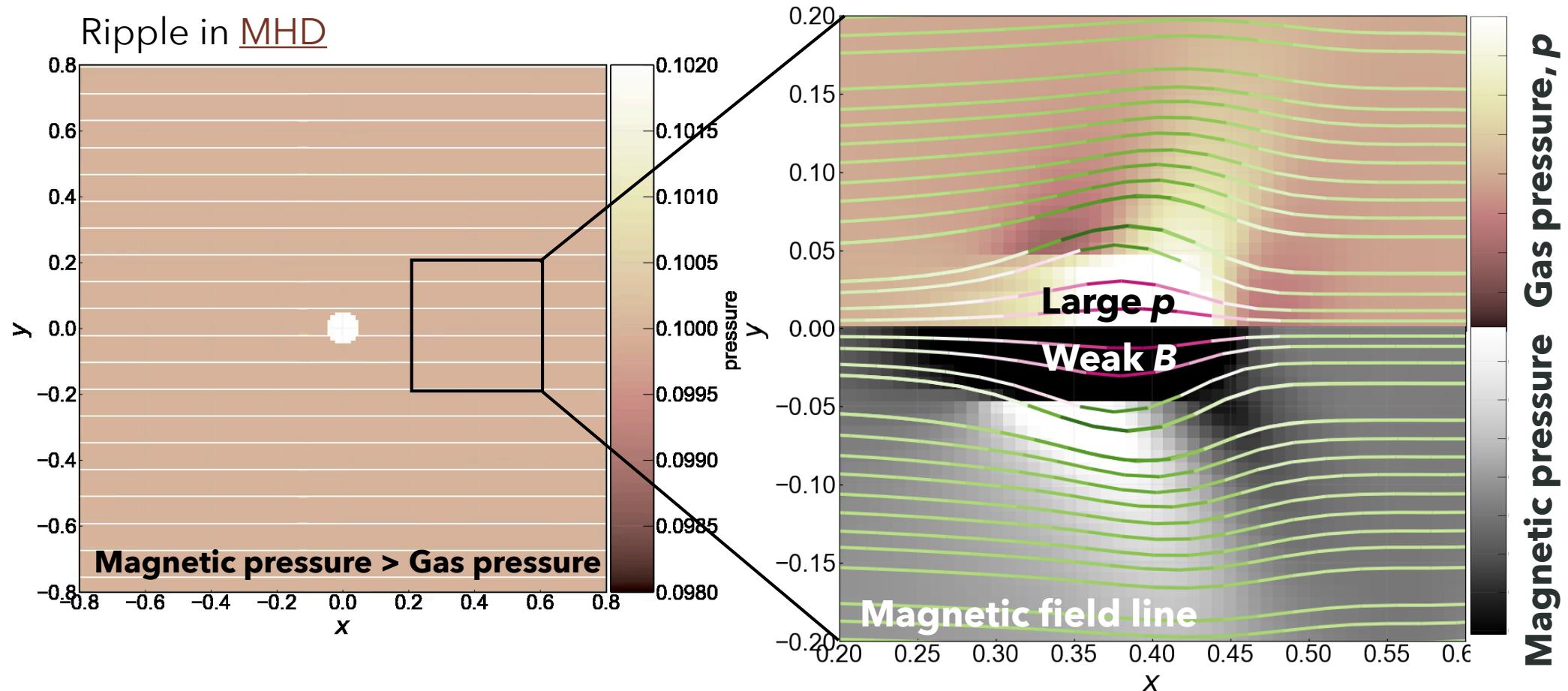
Remarks: Slow-mode (Slow Wave)



Slow-mode: Compressional wave that gas pressure and magnetic pressure have inverse phase of each other.

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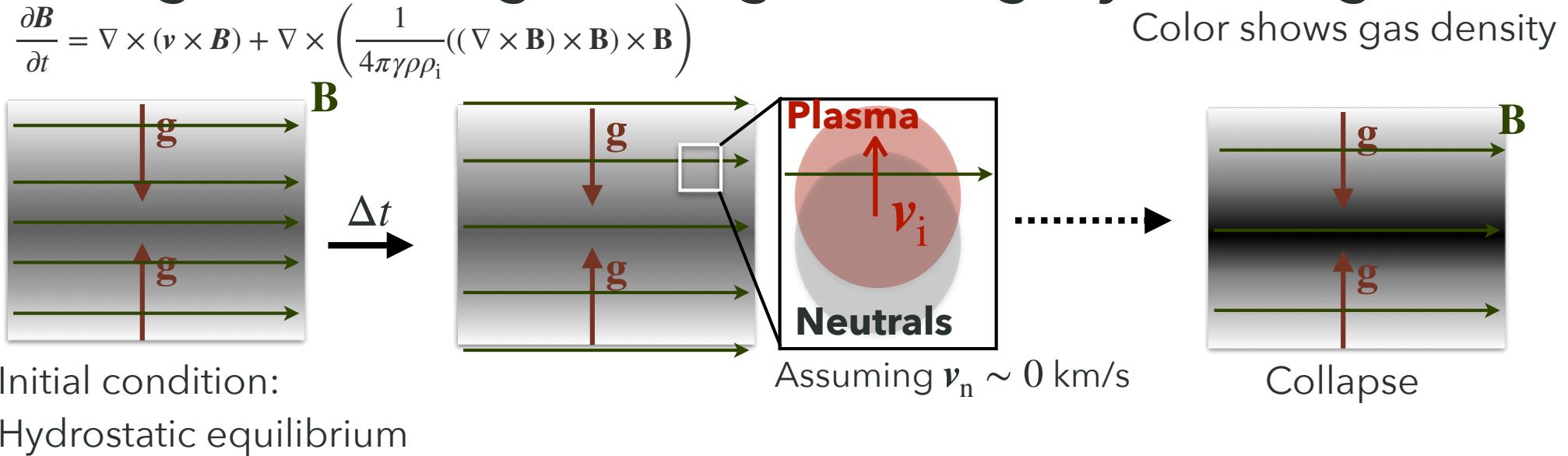


Slow-mode: Compressional wave that gas pressure and magnetic pressure have inverse phase of each other.

The Ion-Neutral Drift (Ambipolar Diffusion)

The neutral fluid and plasma are **not** perfectly coupled in molecular clouds

Ex. Magnetized self-gravitating slab of lightly ionized gas



Magnetic Reynolds
number
for ambipolar diffusion

$$R_{AD} = 1 \rightarrow \ell_{AD} = 0.09 \text{ pc} \left(\frac{B}{30 \mu\text{G}} \right)^2 \left(\frac{n}{10^3 \text{ cm}^{-3}} \right)^{-3/2} \left(\frac{v}{1 \text{ km/s}} \right)^{-1}$$

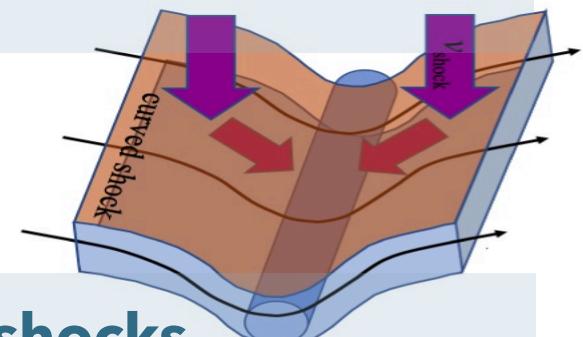
→ AD is effective in filaments ($\ell \sim 0.1$ pc)

Purpose

Aim

Discover the mechanism to keep upstream kinetic energy

→ Understand how massive filament width is maintained



This Study

In reality, there are two shocks.

Investigate whether turbulent ram pressure can be provided by **two slow-shocks simulations with ambipolar diffusion**.

Setup for Simulations

2D simulations using Athena++ code (Stone+ 2020)

Isothermal MHD including ambipolar diffusion (w/o self-gravity)

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot (\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + \mathbf{P} + \mathbf{B}^2/2) = 0$$

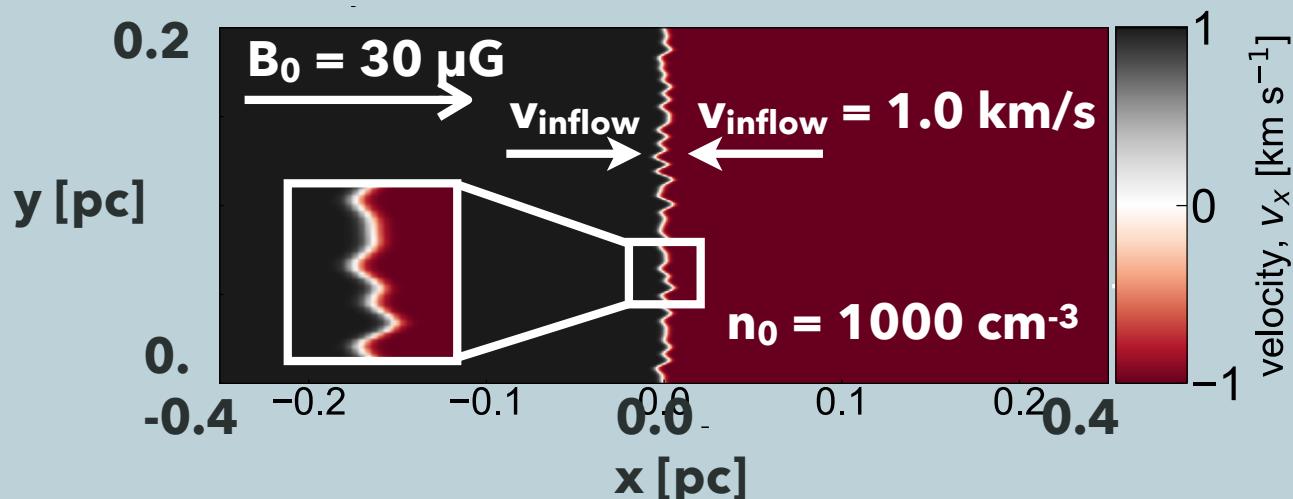
$$\frac{\partial E}{\partial t} + \nabla \cdot \left[(E + P + B^2/2) \mathbf{v} - \mathbf{B}(\mathbf{B} \cdot \mathbf{v}) + \frac{\eta_{\text{AD}}}{|B|^2} \{ \mathbf{B} \times (\mathbf{J} \times \mathbf{B}) \} \times \mathbf{B} \right] = 0$$

$$E = e + \frac{1}{2} \rho v^2 + \frac{B^2}{2} \quad P = (\gamma - 1)e \quad \gamma = 1.01$$

$$\frac{\partial \mathbf{B}}{\partial t} - \nabla \times \left[(\mathbf{v} \times \mathbf{B}) - \frac{\eta_{\text{AD}}}{|B|^2} \mathbf{B} \times ((\nabla \times \mathbf{B}) \times \mathbf{B}) \right] = 0$$

$$\eta_{\text{AD}} = \frac{B^2}{4\pi\gamma_{\text{in}}\rho_{\text{n}}\rho_{\text{i}}} \quad (\text{Shu 1992})$$

Initial Condition: Gas inflows along the B field \rightarrow filament formation



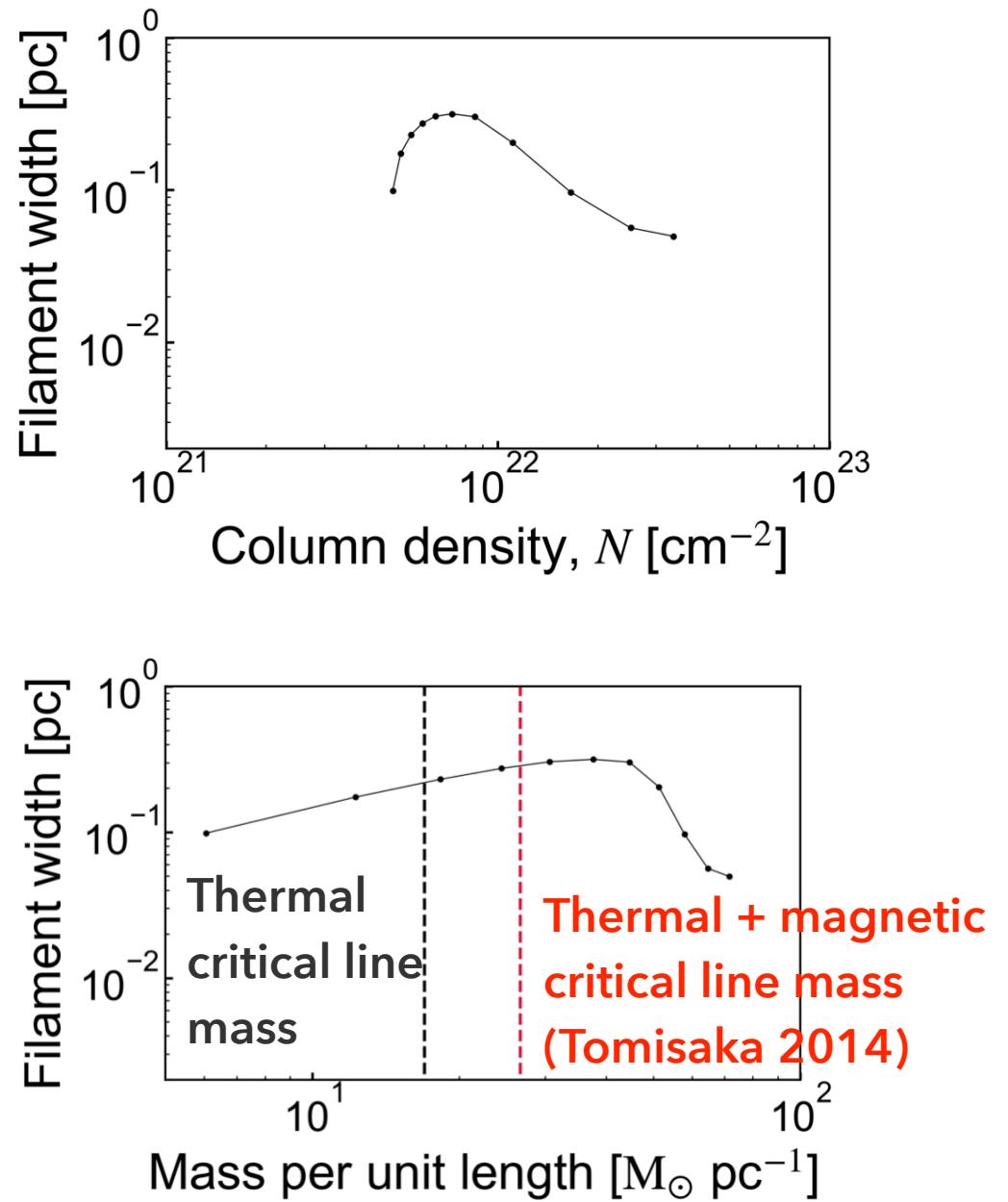
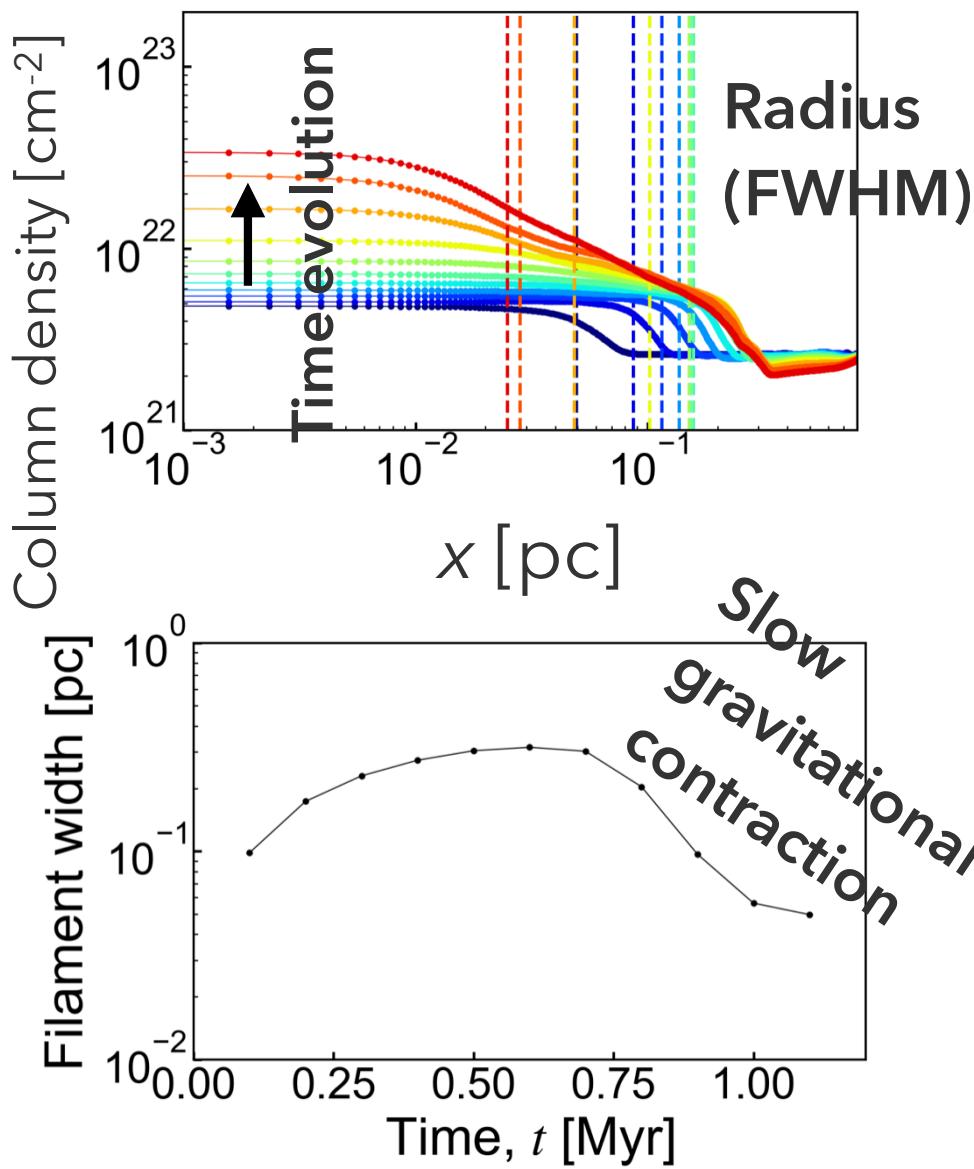
Boundary Condition

- $x \rightarrow \text{Free}$
- $y \rightarrow \text{Periodic}$

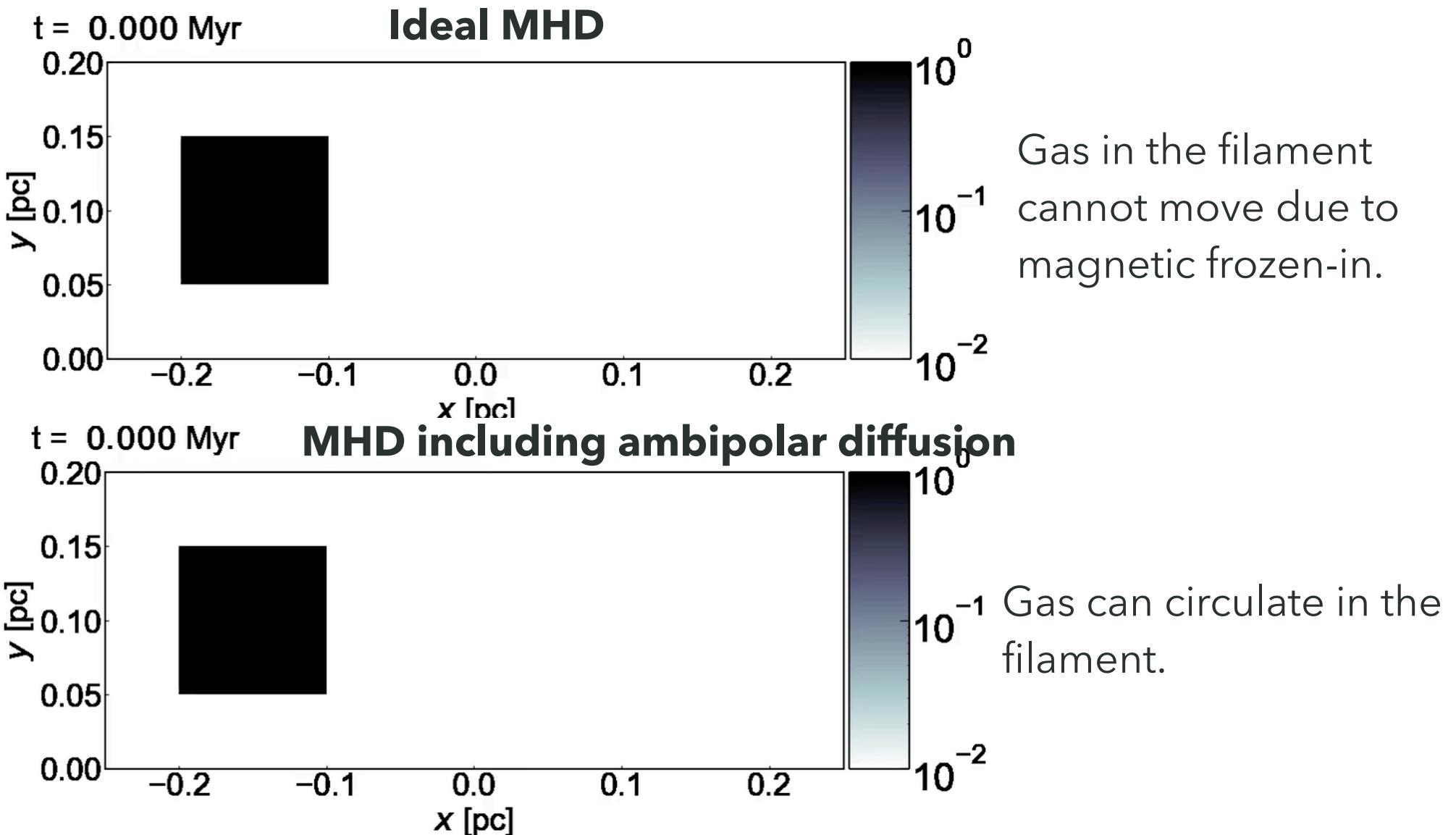
$$c_s = 0.2 \text{ km/s}$$

Resolution : $3.9 \times 10^{-4} \text{ pc}$

Time evolution of filament width

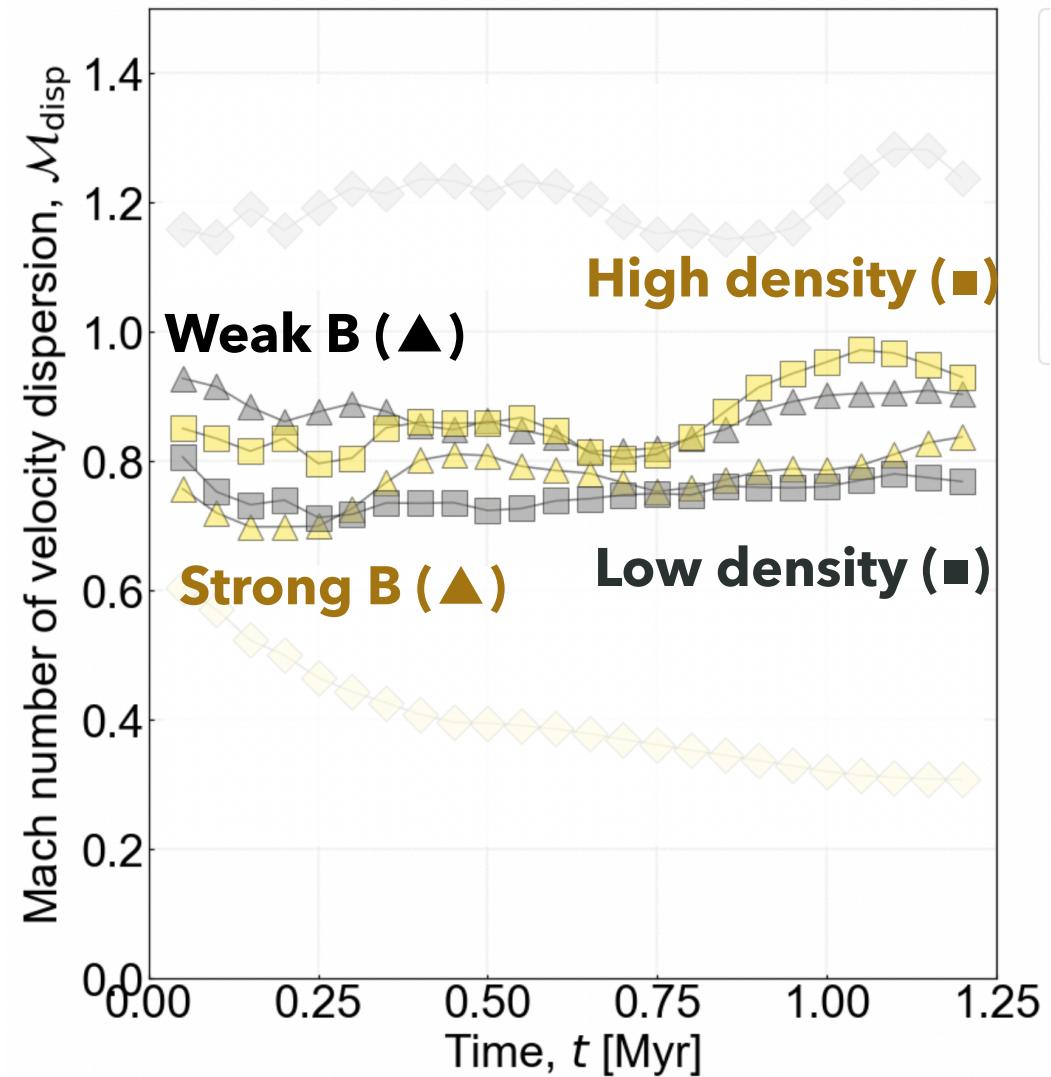
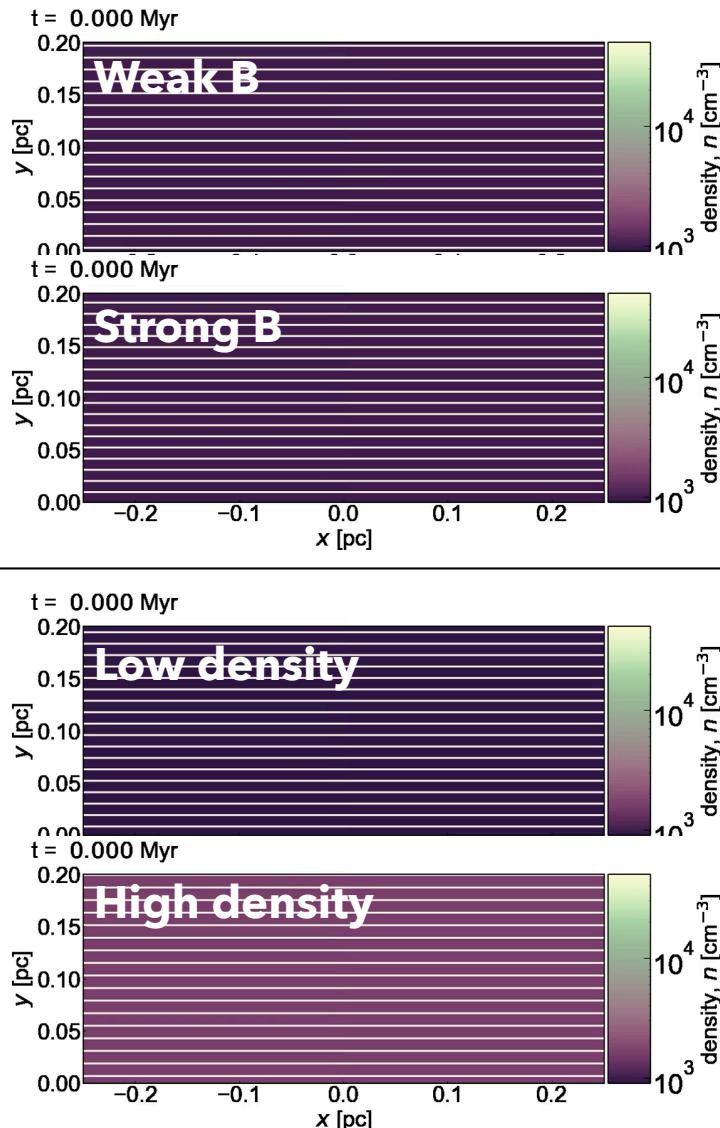


Additional Results: Passive Scalar



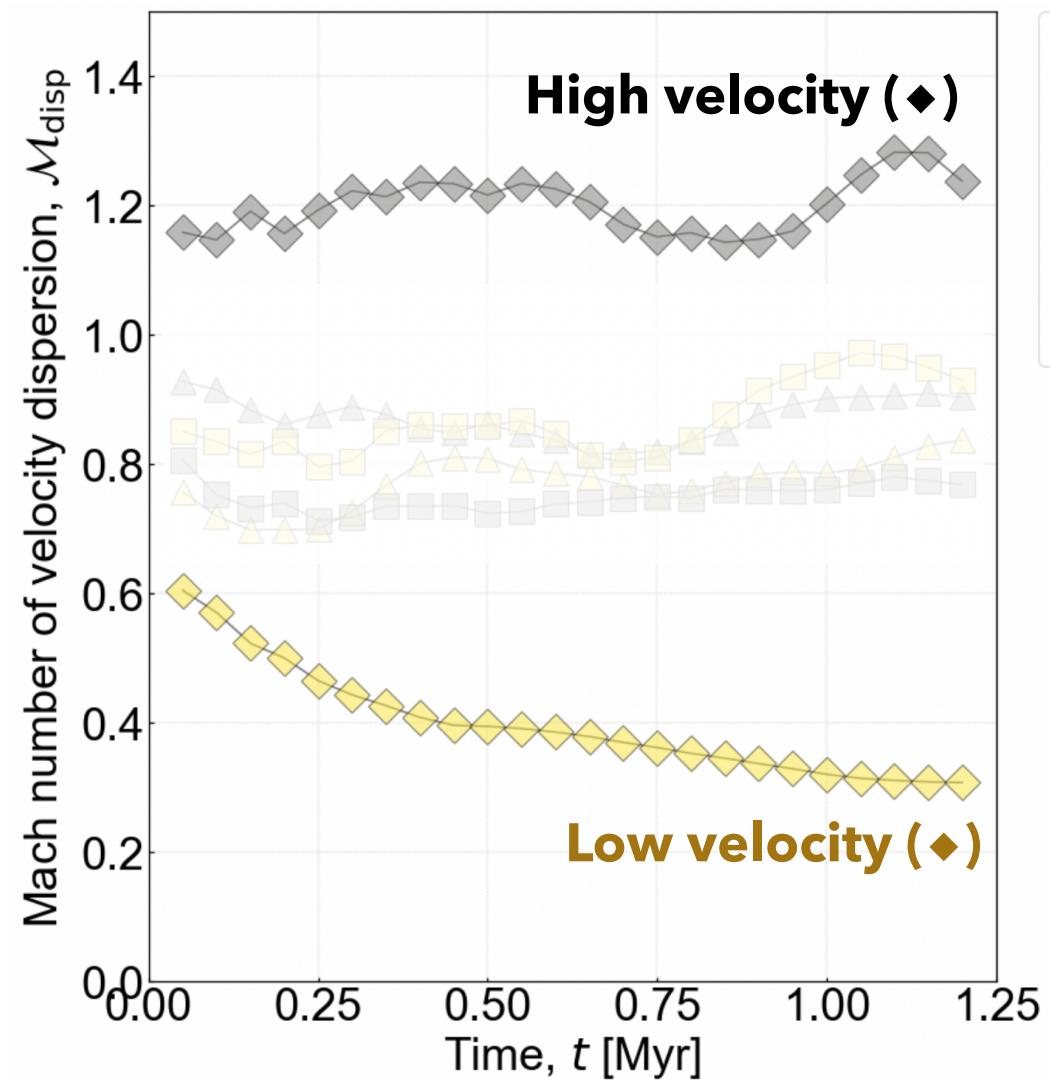
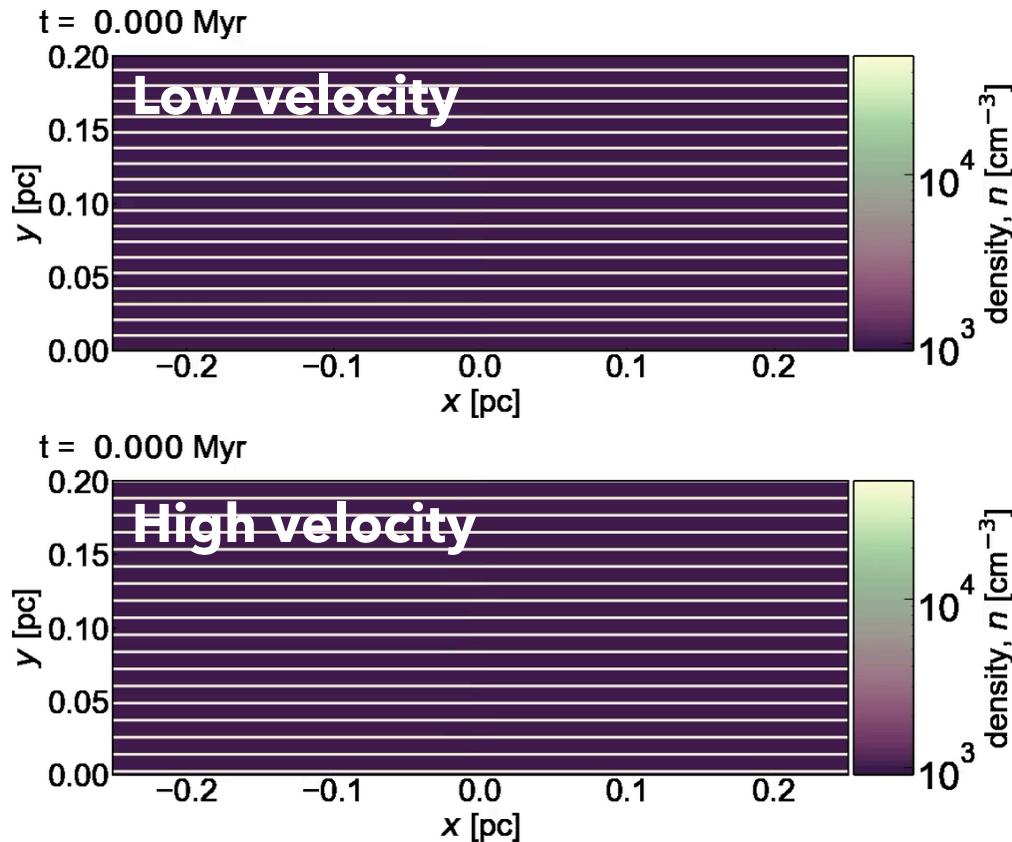
Confirm that the bullets are physical gas blobs.

Dependence on Density and Magnetic Field



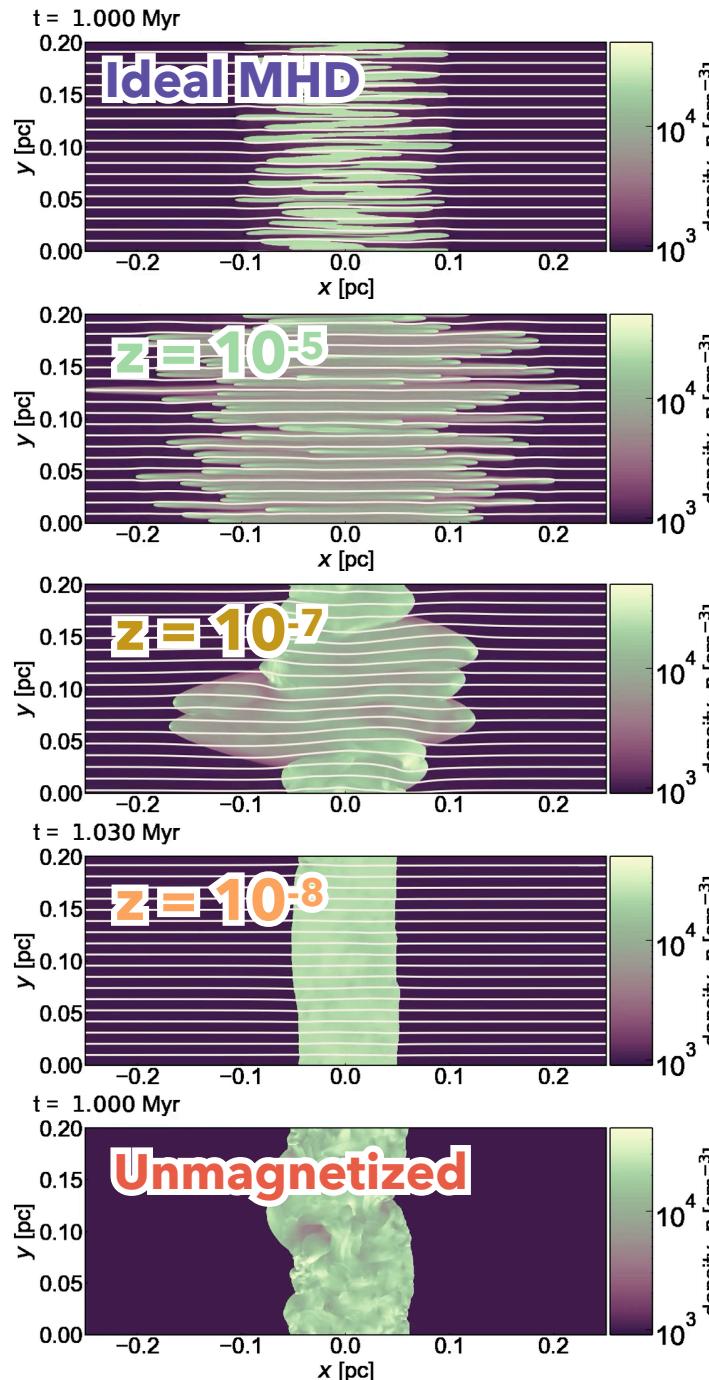
The strength of the generated anisotropic turbulence is independent of the magnetic field and density.

Dependence on Accretion Velocity



The strength of the generated anisotropic turbulence depends on the accretion velocity.

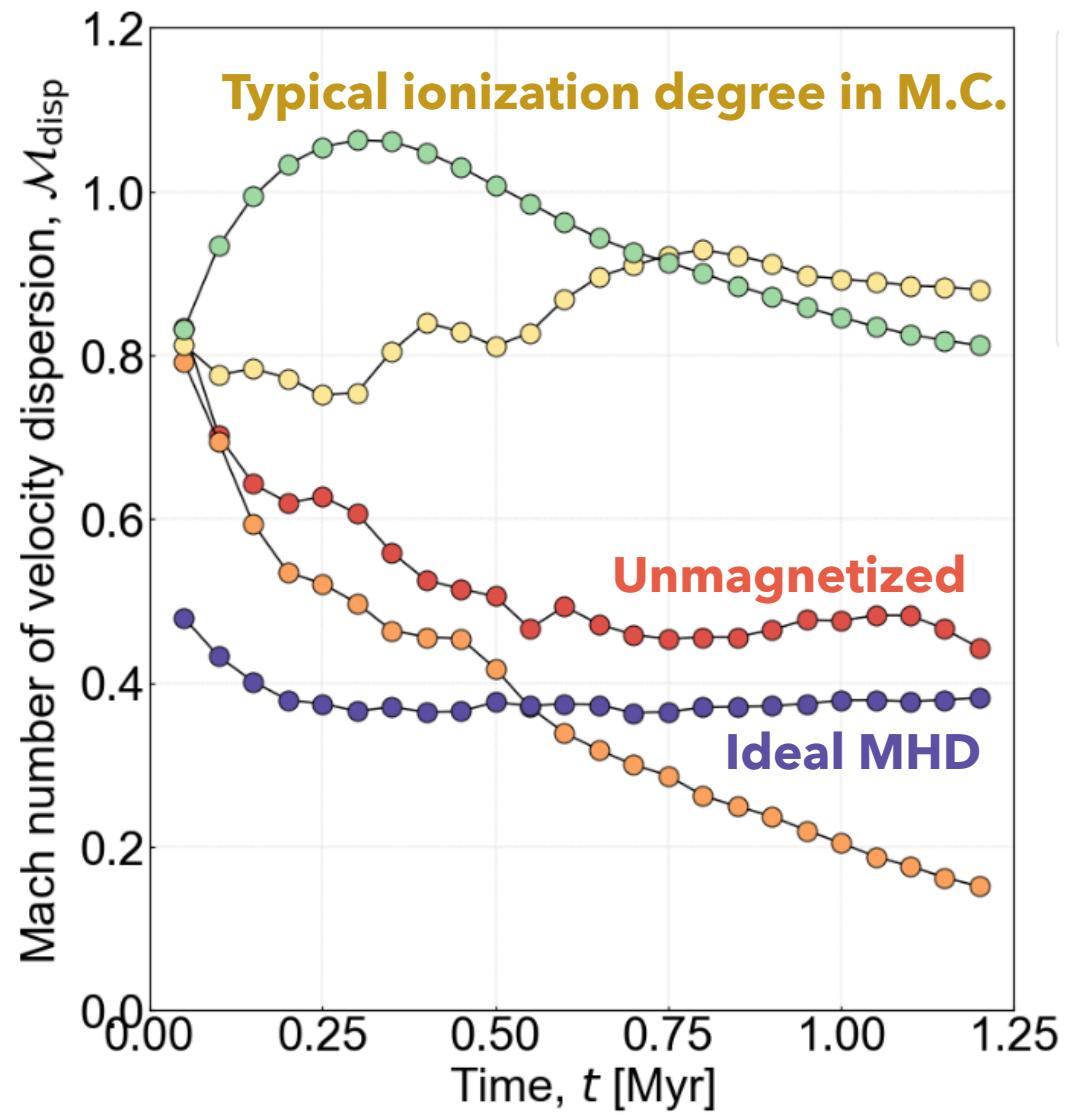
Dependence on Ionization Degree



Large

ionization degree

Small



Bullet mechanism exists even if the ionization degree changes by one or two orders of magnitude.

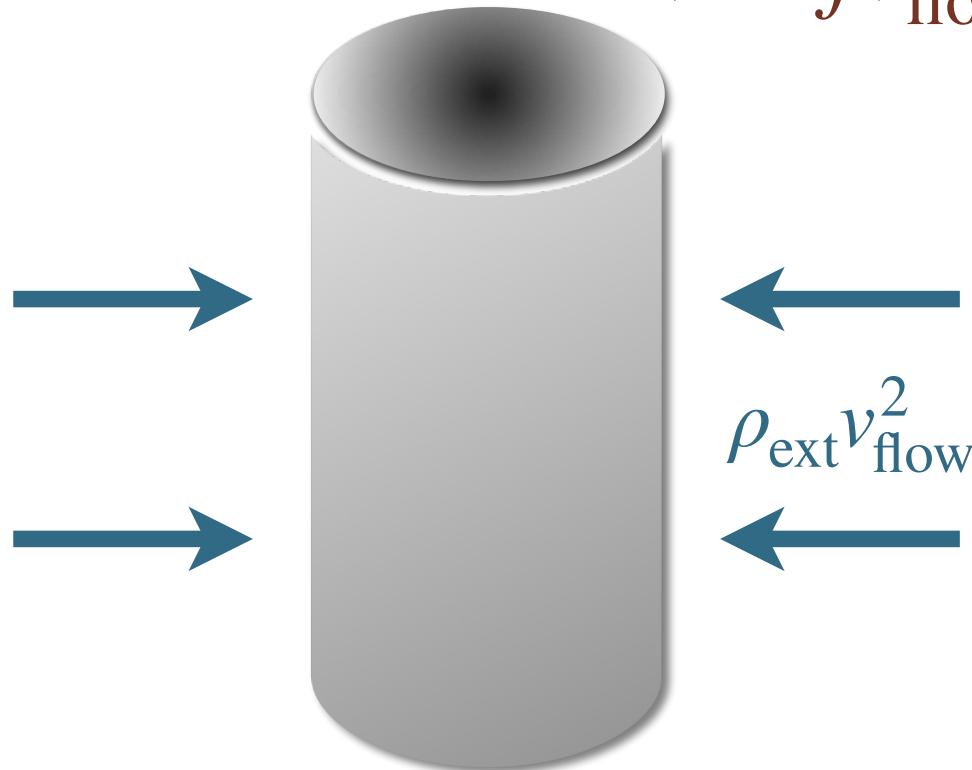
Discussion: Line mass v.s. FWHM (Theoretical Model)

Does flow-driven turbulence explain the observed line mass independent width?

Outline of our model

Velocity dispersion within the filament:

$$\Delta v = f v_{\text{flow}}, \quad (f \simeq 0.5)$$

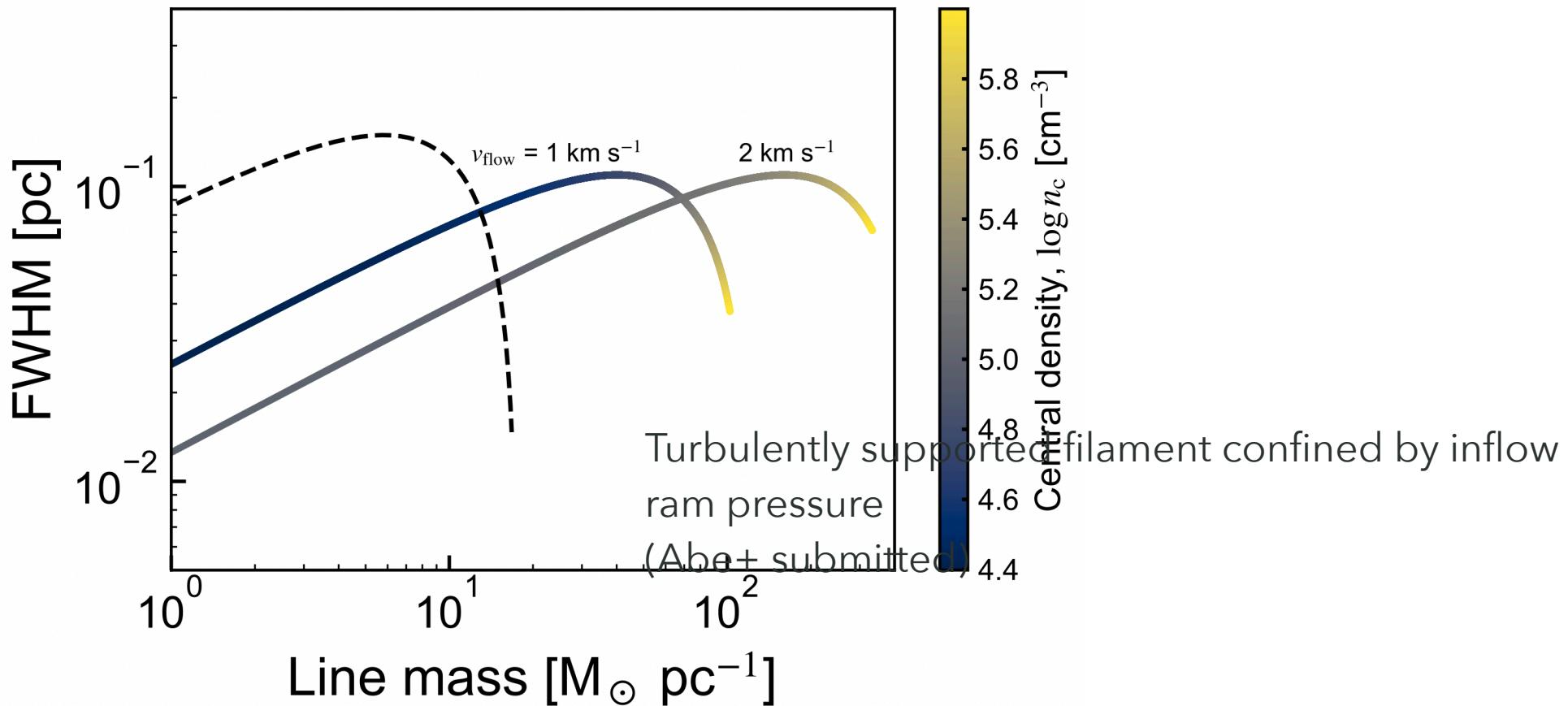


Discussion: Line mass v.s. FWHM (Theoretical Model)

Does flow-driven turbulence explain the observed line mass independent width?

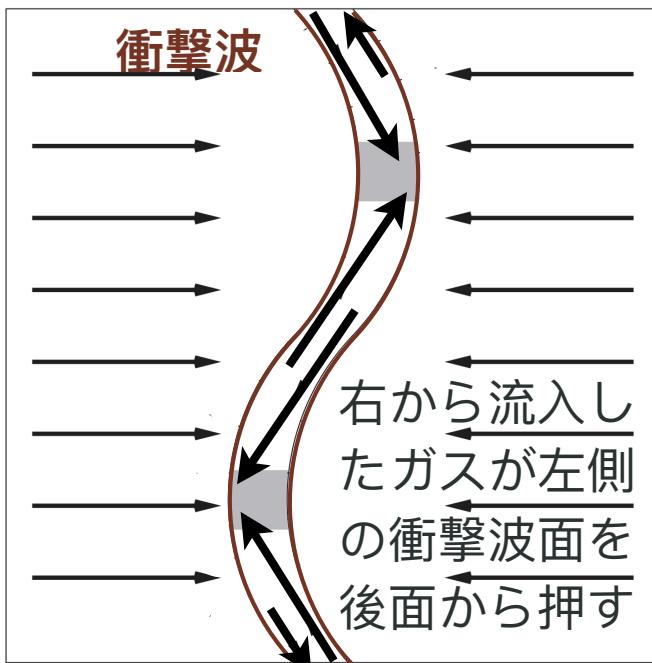
Thermally supported filament w/ external gas pressure

Fischera & Martin (2012)

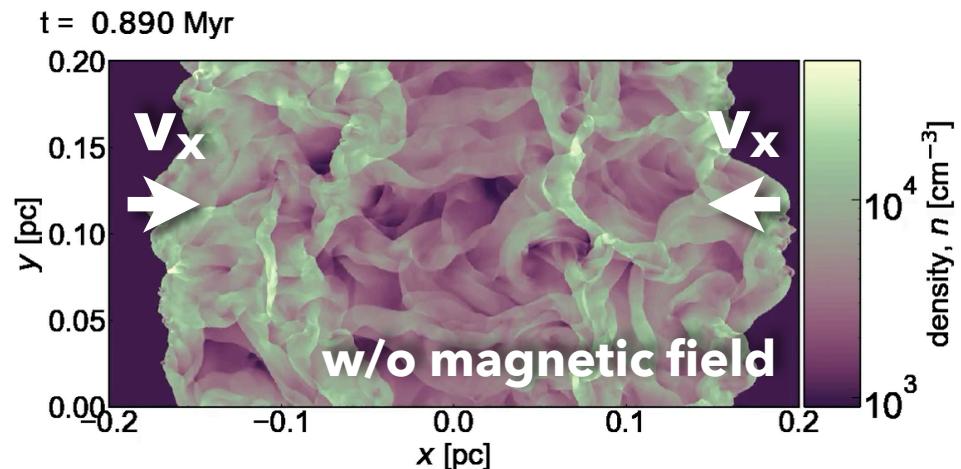


Nonlinear Thin-Shell Instability (NTSI)

Nonlinear Thin Shell Instability (NTSI): 磁場なし等温衝撃波2枚に挟まれた薄いシートにおける衝撃波面の揺らぎは不安定 (e.g., Vishniac 1994)



McLeod & Whitworth (2013)



反対側の衝撃波面への運動量輸送が本質

強磁場 (Slow shock) ではNTSIの成長を抑制する (Heitsh+ 2007)

※ NTSIには飽和があり乱流は維持できないのでフィラメントの幅の維持には向かない話

Future Works

Star Formation Rate Problem

Contradiction between the observation & the theoretical estimation

The observed
star formation rate (SFR)
 $\sim 1 \text{ M}_{\text{sun}}/\text{yr}$



Simple estimation of SFR

Molecular clouds mass in MW $\sim 10^9 \text{ M}_{\text{sun}}$.
Free-fall time scale of $\sim 10^6 \text{ yr}$,
→ The **SFR** $\sim 1,000 \text{ M}_{\text{sun}}/\text{yr}$.

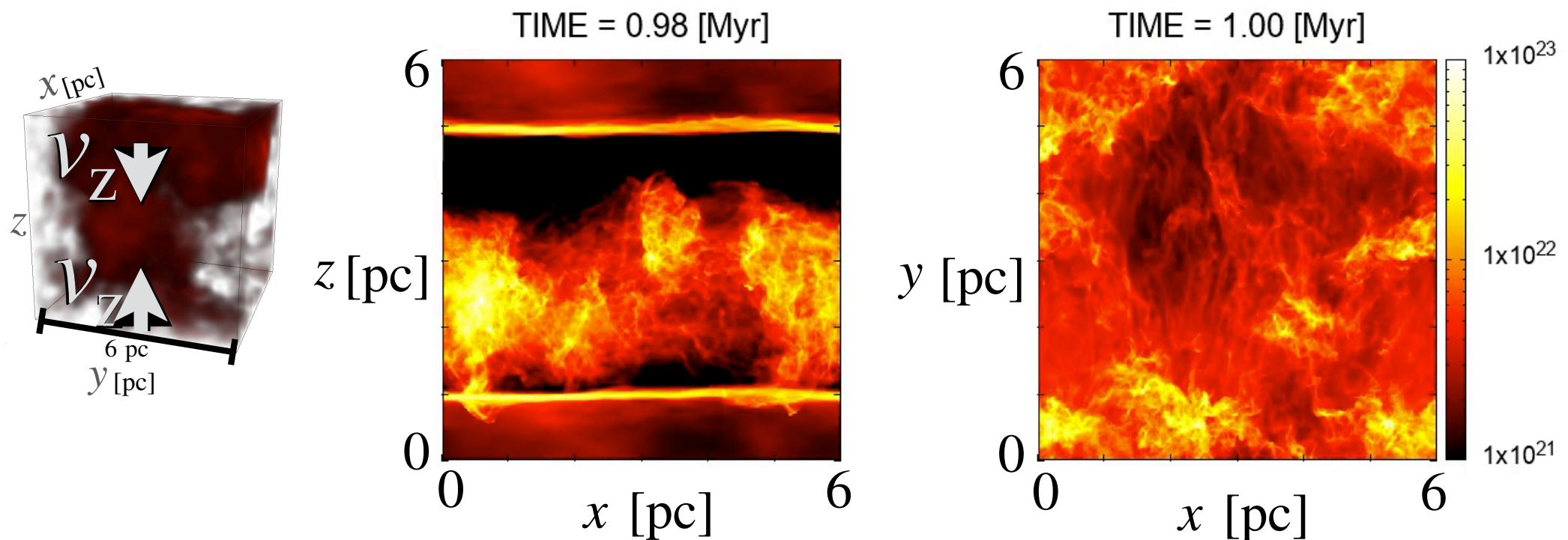
An example of solutions

→ Only a small fraction of the molecular cloud becomes stars.

How can such a situation to be realized?

Our Strategy: Filament "Resetting"

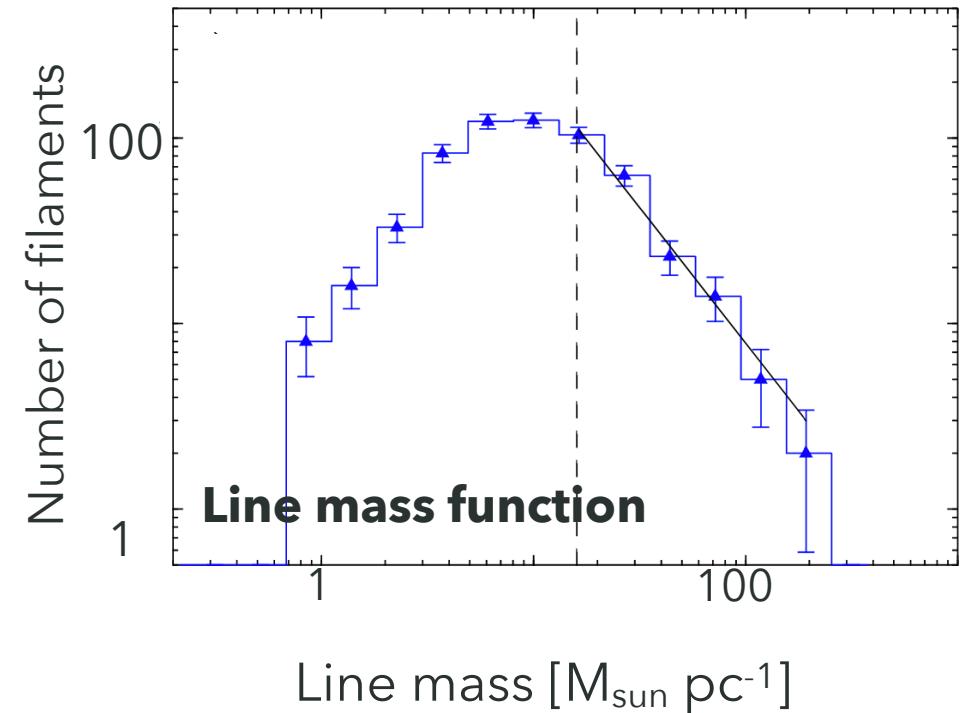
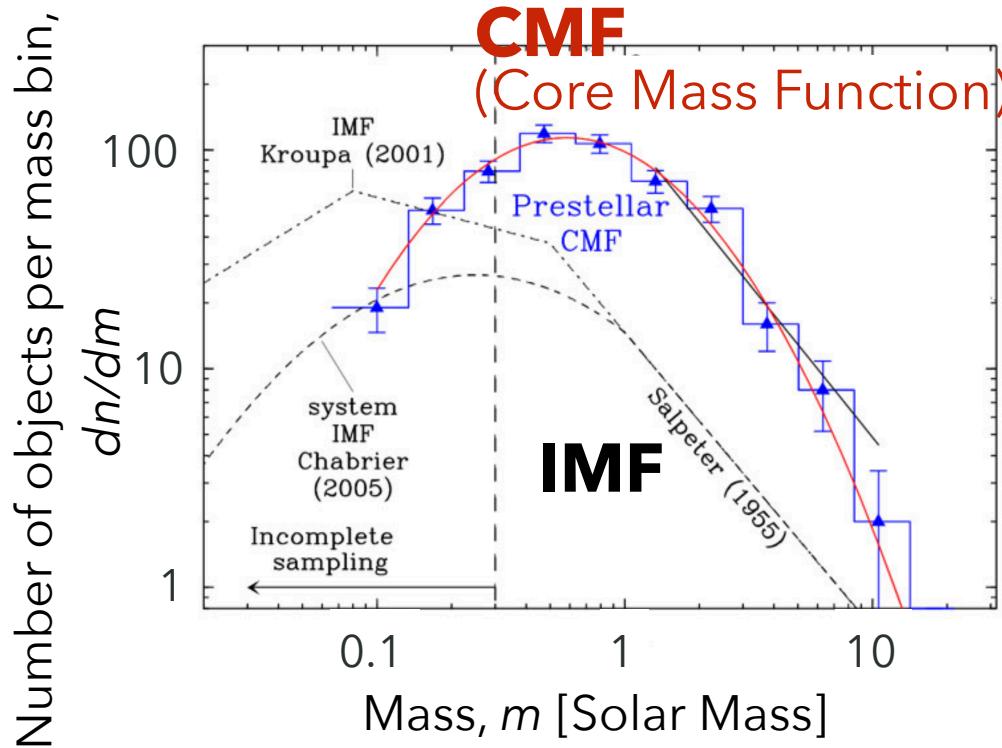
Idea: The effect of multi-compression causes filament **destruction** & formation
→ Filament mass fraction keeps low value?



First, I would like to reveal the conditions of the destruction.

Other Future Work

Link between CMF and Line Mass Function?



I would like to clarify the origins of the IMF through an understanding of filaments.

Summary & Future Works

We perform non-ideal MHD simulations and investigate filament evolution simulations.

“Anisotropic turbulence driven by the Bullet mechanism (SSI + ambipolar diffusion)” can maintain the width of massive filaments.

Future Works

- Solve the 0.1 pc problem
- Understanding the “filament resetting.” \rightarrow Solve the SFR problem
- Clarifying the origin of filament line mass function and relation with IMF